# An Invariance Principle for Isotropic Diffusions in Random Environment 

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Preliminary Draft


#### Abstract

We investigate in this work the asymptotic behavior of isotropic diffusions in random environment that are small perturbations of Brownian motion. When the space dimension is three or more, we prove an invariance principle as well as transience. Our methods also apply to questions of homogenization in random media.


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[^0]
## 0 Introduction

The mathematical investigation of transport in disordered media has been an active field of research over the last thirty years, rich in surprising effects and mathematical challenges. In a number of cases the method of the environment viewed from the particle has proven a powerful tool, cf. De Masi et al. [7], Kipnis-Varadhan [12], Kozlov [13], Molchanov [17], Olla [18], [19], Papanicolaou-Varadhan [21], [22]. However basic models such as random walk in random environment or Brownian motion perturbed by an environmentdependent drift, when the random drift is neither the gradient of a stationary function nor incompressible, have in essence not been amenable to this approach and remain to this day mathematical challenges. An intensive effort to understand these models has been launched in the last five years. Progress has been made, especially in the case of ballistic behavior, i.e. when the particle has a non-degenerate velocity, see for instance [27], [32] and the references therein. As for diffusive behavior, there has been some progress, cf. [4], but overall the topic has been little touched. The present work is precisely concerned with diffusive behavior, and investigates isotropic diffusions in random environment that are small perturbations of Brownian motion. When the space dimension is three or more, we prove transience and an invariance principle. The model we analyze is a continuous counterpart of the model studied by Bricmont-Kupiainen [5]. However our strategy of proof is different and we believe more transparent.

Let us first describe the setting in more details. The local characteristics, i.e. covariance and drift, of the diffusion in random environment are bounded stationary functions $a(x, \omega), b(x, \omega), x \in \mathbb{R}^{d}, \omega \in \Omega$, with respective values in the non-negative $d$-matrices and $\mathbb{R}^{d}, d \geq 3$; the set $\Omega$ is endowed with a group $\left(t_{x}\right)_{x \in \mathbb{R}^{d}}$ of jointly measurable transformations preserving the probability $\mathbb{P}$ on $\Omega$. We assume that for $\omega \in \Omega, a(\cdot, \omega)$ is uniformly elliptic, see (1.5), and that

$$
\begin{equation*}
a(\cdot, \omega) \text { and } b(\cdot, \omega) \text { satisfy a Lipschitz condition with constant } K \text {, cf. (1.4). } \tag{0.1}
\end{equation*}
$$

We denote with $P_{x, \omega}$ the law of the diffusion in the environment $\omega$, starting from $x$, i.e. the unique probability on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$ solution of the martingale problem attached to $x$ and

$$
\begin{equation*}
L=\frac{1}{2} \sum_{i, j=1}^{d} a_{i j}(y, \omega) \partial_{i j}^{2}+\sum_{i=1}^{d} b_{i}(y, \omega) \partial_{i}, \tag{0.2}
\end{equation*}
$$

cf. [26]. We let $\left(X_{t}\right)_{t \geq 0}$ stand for the canonical process on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$.
The random characteristics of the diffusion are assumed to have finite range dependence, namely for some $R>0$, under $\mathbb{P}$,

$$
\begin{equation*}
\sigma(a(x, \cdot), b(x, \cdot), x \in A) \text { and } \sigma(a(y, \cdot), b(y, \cdot), y \in B) \text { are independent } \tag{0.3}
\end{equation*}
$$ when $A, B \subseteq \mathbb{R}^{d}$ have mutual distance at least $R$.

Further they also fulfill a restricted isotropy condition, namely for any rotation matrix $r$ preserving the union of coordinate axes of $\mathbb{R}^{d}$,

$$
\begin{equation*}
(a(r x, \omega), b(r x, \omega))_{x \in \mathbb{R}^{d}} \text { has same law under } \mathbb{P} \text { as }\left(r a(x, \omega) r^{T}, r b(x, \omega)\right)_{x \in \mathbb{R}^{d}} \tag{0.4}
\end{equation*}
$$

we refer to Section 1 for details.

The main result of this article, of Theorem 6.3, states that
Theorem. ( $d \geq 3$ )
There is an $\eta_{0}(d, K, R)>0$, such that if

$$
\begin{equation*}
|a(x, \omega)-I| \leq \eta_{0},|b(x, \omega)| \leq \eta_{0}, \text { for all } x \in \mathbb{R}^{d}, \omega \in \Omega \tag{0.5}
\end{equation*}
$$

then for $\mathbb{P}$-a.e. $\omega$,

$$
\begin{align*}
& \frac{1}{\sqrt{t}} X_{\cdot t} \text { under } P_{0, \omega} \text { converges in law to Brownian motion on }  \tag{0.6}\\
& \mathbb{R}^{d} \text { with deterministic variance } \sigma^{2}>0 \text {, as } t \rightarrow \infty
\end{align*}
$$

and

$$
\begin{equation*}
\text { for all } x \in \mathbb{R}^{d} \text {, } P_{x, \omega} \text {-a.s., } \lim _{t \rightarrow \infty}\left|X_{t}\right|=\infty \tag{0.7}
\end{equation*}
$$

In other words for diffusions in random environment that are small perturbations of Brownian motion and satisfy the restricted isotropy condition (0.4), we prove transience and diffusive behavior. Our results also apply to questions of homogenization in random media, cf. Theorem 6.4, and show that

Theorem. ( $d \geq 3$ )
One can choose $\eta_{0}(d, K, R)>0$, so that when (0.5) holds, on a set of full $\mathbb{P}$-probability, for any bounded functions $f, g$ on $\mathbb{R}^{d}$, respectively continuous and Hölder continuous, the solution of the Cauchy problem:

$$
\left\{\begin{align*}
\partial_{t} u_{\epsilon} & =L_{\epsilon} u_{\epsilon}+g, \text { in }(0, \infty) \times \mathbb{R}^{d}  \tag{0.8}\\
\left.u_{\epsilon}\right|_{t=0} & =f
\end{align*}\right.
$$

where for $\epsilon>0$,

$$
\begin{equation*}
L_{\epsilon}=\frac{1}{2} \sum_{i, j=1}^{d} a_{i j}\left(\frac{x}{\epsilon}, \omega\right) \partial_{i j}^{2}+\sum_{i=1}^{d} \frac{1}{\epsilon} b_{i}\left(\frac{x}{\epsilon}, \omega\right) \partial_{i} \tag{0.9}
\end{equation*}
$$

converges uniformly on compact subsets of $\mathbb{R}_{+} \times \mathbb{R}^{d}$, as $\epsilon \rightarrow 0$, to the solution of the Cauchy problem

$$
\left\{\begin{align*}
\partial_{t} u_{0} & =\sigma^{2} \Delta u_{0}+g, \text { in }(0, \infty) \times \mathbb{R}^{d}  \tag{0.10}\\
\left.u_{0}\right|_{t=0} & =f
\end{align*}\right.
$$

When $b(\cdot, \omega) \equiv 0$, cf. [22], [31], or when $L$ is in divergence form, cf. [7], [13], [19], [20], [21], the method of the environment viewed from the particle applies successfully, and there is an extensive literature on invariance principles describing diffusive behavior and applications to homogenization. There is also ample literature on analogous discrete situations, cf. [2], [3], [12], [13], [14], [15]. On the other hand the case of general diffusions in random environment of type (0.2) remains poorly understood, reflecting the genuine non self-adjoint character of the problem and the absence of invariant measure at hand. We do not know of any work proving diffusive behavior, and in the context of random walks in random environment only of [4], [5].

We will now give some description of the proof of our results. The main point appears in Theorem 1.1. It states an induction step concerning the behavior of the diffusion in random environment along a sequence of length scales $L_{n} \simeq L_{0}^{(1+a)^{n}}$ and time scales $L_{n}^{2}$, where $a$ is a small positive number and $L_{0}$ in a large enough number, cf. (1.14), (1.15). Several assumptions are propagated from level $n$ to level $n+1$. A first assumption, cf. (1.47), states that up to a $\mathbb{P}$-probability decaying like a large negative power of $L_{n}$, the following holds. On the one hand, for starting points $x$ with distance const $L_{n}$ from the origin, the displacements of the path of the diffusion in the environment $\omega$ slightly beyond distances of order $L_{n}$ satisfy under $P_{x, \omega}$ a certain exponential control, cf. (1.39), and on the other hand the transition kernel at time $L_{n}^{2}$ of the diffusion:

$$
\begin{equation*}
R_{n}(x, d y)=P_{x, \omega}\left[X_{L_{n}^{2}} \in d y\right] \tag{0.11}
\end{equation*}
$$

is in a sense that we explain below "close" to the Gaussian kernel

$$
\begin{align*}
& R_{n}^{0}(x, d y)=\left(2 \pi \alpha_{n} L_{n}^{2}\right)^{-d / 2} \exp \left\{-\frac{(y-x)^{2}}{2 \alpha_{n} L_{n}^{2}}\right\} d y, \text { with }  \tag{0.12}\\
& \alpha_{n} \approx \mathbb{E} E_{0, \omega}\left[\left|X_{L_{n}^{2}}\right|^{2}\right] /\left(d L_{n}^{2}\right),
\end{align*}
$$

(cf. (1.22) for the true definition), after localization of $x$ in a box of size const $L_{n}$ around the origin. The way in which "close" is defined plays a pivotal role in this work. It refers to the operator norm $\|\cdot\|_{n}$, for linear transformations on the space of bounded Hölder continuous functions of order $\beta$ (some fixed number in ( $0, \frac{1}{2}$ ], cf. (1.13)), endowed with the norm $|\cdot|_{(n)}$, cf. (1.28), adapted to functions "living in scale $L_{n}$ ":

$$
\begin{equation*}
|f|_{(n)}=\sup _{x \in \mathbb{R}^{d}}|f(x)|+\sup _{x \neq y} \frac{|f(x)-f(y)|}{\left|\frac{x-y}{L_{n}}\right|^{\beta}} . \tag{0.13}
\end{equation*}
$$

In essence "close" means $\left\|\chi_{n, 0}\left(R_{n}-R_{n}^{0}\right)\right\|_{n} \leq$ const $L_{n}^{-\delta}$, where $\chi_{n, 0}$ is a cut-off function localizing $x$ in (0.11), (0.12), within distance const $L_{n}$ of the origin, cf. (1.38), and $\delta>0$ is a fraction of $\beta$, cf. (1.40).

A second assumption being propagated, cf. (1.48), states quantitatively the rarity of traps by describing the domination of the tails under $\mathbb{P}$ of certain variables measuring the strength of traps in boxes of size $L_{n}$, cf. (1.44), by the corresponding tails of i.i.d. variables equal to 0 with overwhelming probability. The word "traps" refers to the fact that in some pockets of the medium, the random $\operatorname{drifts} b(\cdot, \omega)$ may concur to capture the particle for a long time.

The third and last assumption entering the induction step, cf. (1.49), controls the behavior of $\alpha_{n}$.

Once Theorem 1.1 is proved, we show in Section 6 that when the local characteristics of the diffusion satisfy (0.5), we can start the induction stated in Theorem 1.1. So the induction assumptions propagate to all levels $n$, and with Borel-Cantelli's lemma we see that all boxes $L_{n}$ within distance const $L_{n+3}^{2}$ of the origin "behave well". With the Kantorovich-Rubinstein theorem, cf. [8], the Hölder-norm estimates and the controls on displacements of the diffusion, cf. (1.47), enable to construct "good couplings" between the diffusion in random environment and Brownian motion with variance $\alpha_{n}$, cf. Proposition 6.2. Since $\alpha_{n}$ converges to a positive limit, namely $\sigma^{2}$ of ( 0.6 ), the invariance principle easily follows. The transience of the diffusion, cf. (0.7), and the homogenization result (0.8), (0.10), also come as easy consequences of these coupling measures.

Let us explain how the article is organized and briefly comment on each section. Section 1 presents the setting and states Theorem 1.1. The proof of Theorem 1.1 occupies Sections 2 to 5 of the article.

Section 2 propagates from level $n$ to level $n+1$ the controls on the displacement of the path, cf. Proposition 2.2.

Section 3 propagates the controls on traps, cf. (1.48) and Proposition 3.3. Traps are a serious matter in our problem because a pocket of size $L$ has the potential, depending on the realization of the medium, to entrap the particle for times of exponential order in $L$. Hence pockets of relatively modest size may distort the diffusive behavior of the particle on many time scales $L_{n}^{2}$. This feature naturally affects the distribution of the variables in (1.44) that measure the strength of traps. We are in fact mainly interested in a small portion of the information contained in (1.48), namely ensuring that the variables in (1.44) vanish with "overwhelming probability", cf. (5.2), (5.3). But the inductive proof requires a control on the tails of the variables in (1.44). To carry the tail domination control (1.48) from level $n$ to level $n+1$, in essence we exhibit exit strategies for the particle from boxes of size $L_{n+1}$ before time $L_{n+1}^{2}$, which show that it is costly for the medium to produce a trap at level $n+1$ of a given strength. Depending on the strength in question, the exit strategy that is employed varies, and we distinguish four distinct regimes, (three regimes suffice when $d \geq 4$ ), cf. (3.20).

Sections 4 and 5 are devoted to the propagation from level $n$ to level $n+1$ of the Hölder-norm controls contained in (1.47).

In Section 4, we perform "surgery" in a large box of size const $L_{n+1}^{2}$ around the origin, which contains the relevant portion of the medium for our purpose. We investigate at a finite depth $n-m_{0}-1$, with $m_{0}$ a fixed number, cf. (1.17), this large box, remove all boxes of size $L_{n-m_{0}-1}$ where bad behavior in the sense of (1.47) occurs, and in essence replace them with good boxes. In this new artificial environment "after surgery", we analyze the diffusion at all the levels $n^{\prime}$ between $n-m_{0}-1$ and $n+1$. We show that with overwhelming $\mathbb{P}$-probability this environment not only does not develop in these intermediate levels bad Hölder-norm behavior with distance $L_{n+1}^{2}$ from the origin, but produces a decay of the relevant $\|\cdot\|_{n^{\prime}}$-norms faster than $L_{n^{\prime}}^{-\delta}$, cf. Proposition 4.11. Wavelets, cf. [6], [16], turn out to provide a powerful tool in the control of the $\|\cdot\|_{n^{\prime}}$-norms of certain random linear operators, cf. Lemma 4.5 and 4.6. Collecting lemmas 4.2 to 4.6 , one can read that the relevant $\|\cdot\|_{n^{\prime}}$ norms mentioned above "contract like $L_{n^{\prime}}^{-\beta / 3 \wedge(1-\beta) \wedge(d / 2-1)}$ ", see also Remark 4.7.

In Section 5, we compare at level $n+1$ the true environment with the environment after surgery constructed in Section 4. The difference between them resides in a few defects of size $L_{n-m_{0}-1}$. Thanks to the controls on traps in (1.48), we can assume that these defects have no trapping power. Then using a strategy close in spirit to Section 2 of [25], we show that the Hölder regularity of the kernels of the diffusion in the environment after surgery performed in Section 4, tends to repair the small defects of the true environment, cf. Proposition 5.1. One can then recover with large $\mathbb{P}-$ probability the bound $\| \chi_{n+1,0}\left(R_{n+1}-\right.$ $\left.R_{n+1}^{0}\right) \|_{n+1} \leq$ const $L_{n+1}^{-\delta}$, required to prove (1.47) at level $n+1$, and the discrepancy $\left|\alpha_{n+1}-\alpha_{n}\right|$ is controlled in Proposition 5.7.

Section 6 as indicated previously applies Theorem 1.1 to the proof of the main Theorem 6.3 , cf. also (0.6), (0.7), and to the derivation of an homogenization result, cf. Theorem 6.4 and (0.8), (0.10).

The Appendix collects some useful results on the norms $|\cdot|_{(n)}$ on the space of $\beta$-Hölder continuous functions, cf. (0.13), and on the control of the corresponding operator norms $\|\cdot\|_{n}$ with wavelets, cf. Proposition A.2.

The work by Bricmont-Kupiainen [5] was certainly a source of inspiration for the present work even if we had difficulty to follow some of their arguments. Our proof albeit using renormalization follows a different track. It may be helpful to highlight some of the differences beyond the fact that in [5] the setting is discrete and here it is continuous. In this article we introduce a family of Hölder-norms that play an important role both for their contraction properties and the couplings they enable to construct. They also motivate the use of wavelets. Further we directly compare the quenched transition kernels of the diffusion, cf. (0.11) to certain Gaussian kernels, cf. (0.12), and not to the $\mathbb{P}$-average of the kernels in (0.11). This simplifies the proof. Our bounds on traps are conducted in a different fashion, that is more in line with [29]. We do not carry in our induction a decomposition of the kernels into "small field" and "large field". The scales along which we perform renormalization here grow faster than geometrically, and we perform surgery at a finite depth, and compare what happens in true and "after surgery" environments. Our proof also enables to have, unlike [5], a concise induction step stated in Theorem 1.1. We believe this is a source of clarity.

Finally let us say a few words concerning the decision to work in a continuous rather than discrete setting. It entails some simplifications because a number of scaling arguments become natural and straightforward. But it also bears some technical intricacies related to regularity questions at small scales. Decisive was perhaps the fact that some of the calculations involving wavelets are more transparent and standard when one uses wavelets on $\mathbb{R}^{d}$, rather than wavelets on $\mathbb{Z}^{d}$, cf. [16], §7.3.3.

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## 1 Setting and main induction step

In this section we introduce notations for the main objects of interest and collect some of their elementary properties. We also present in Theorem 1.1 the induction assumption that will be propagated. The proof of Theorem 1.1 occupies the next four sections.

We let $\left(e_{i}\right)_{1 \leq i \leq d}$ stand for the canonical basis of $\mathbb{R}^{d}$, and $d \geq 3$ throughout the article. We respectively denote with $|\cdot|$ and $|\cdot|_{\infty}$ the Euclidean and supremum distances on $\mathbb{R}^{d}$. We let $B(x, r)$ and $\bar{B}(x, r)$ stand for the open and closed Euclidean balls with center $x \in \mathbb{R}^{d}$ and radius $r>0$, and write $B_{\infty}(x, r), \bar{B}_{\infty}(x, r)$ for the corresponding $|\cdot|_{\infty}$-balls. For $A, B$ subsets of $\mathbb{R}^{d}$ we denote with

$$
\begin{equation*}
d(A, B)=\inf \{|x-y| ; x \in A, y \in B\} \tag{1.1}
\end{equation*}
$$

their mutual $|\cdot|$-distance, and with $d_{\infty}(A, B)$ their analogously defined mutual $|\cdot|_{\infty^{-}}$ distance. When $\mathcal{U}$ is a finite subset, we write $|\mathcal{U}|$ for the cardinality of $\mathcal{U}$.

The random environment is described by $(\Omega, \mathcal{A}, \mathbb{P})$ a probability space endowed with $\left(t_{x}\right)_{x \in \mathbb{R}^{d}}$ a bi-measurable group of $\mathbb{P}$-preserving transformations. The diffusion matrix and
the drift of the diffusion in random environment are stationary functions $a(x, \omega), b(x, \omega)$, $x \in \mathbb{R}^{d}, \omega \in \Omega$, with respective values in the space $M_{d}^{+}$of non-negative $d$-matrices and $\mathbb{R}^{d}$ :

$$
\begin{equation*}
a\left(x, t_{y} \omega\right)=a(x+y, \omega), b\left(x, t_{y} \omega\right)=b(x+y, \omega), \text { for } x, y \in \mathbb{R}^{d}, \omega \in \Omega \tag{1.2}
\end{equation*}
$$

We assume that these functions are bounded and uniformly Lipschitz, i.e. there is $K>1$, such that for $x, y \in \mathbb{R}^{d}, \omega \in \Omega$,

$$
\begin{align*}
& |b(x, \omega)|+|a(x, \omega)| \leq K  \tag{1.3}\\
& |b(x, \omega)-b(y, \omega)|+|a(x, \omega)-a(y, \omega)| \leq K|x-y| \tag{1.4}
\end{align*}
$$

Further we assume that the diffusion matrix is uniformly elliptic, i.e. there is a $\nu>1$, such that for $x \in \mathbb{R}^{d}, \omega \in \Omega$ :

$$
\begin{equation*}
\frac{1}{\nu} I \leq a(x, \omega) \leq \nu I \tag{1.5}
\end{equation*}
$$

As mentioned in (0.3) the local characteristics of the diffusion satisfy a condition of finite range dependence. Namely for $A \subseteq \mathbb{R}^{d}$, we define

$$
\begin{equation*}
\mathcal{G}_{A}=\sigma(a(x, \cdot), b(x, \cdot) ; x \in A), \tag{1.6}
\end{equation*}
$$

and assume that for some $R>0$,

$$
\begin{equation*}
\mathcal{G}_{A} \text { and } \mathcal{G}_{B} \text { are independent under } \mathbb{P} \text { whenever } d(A, B) \geq R . \tag{1.7}
\end{equation*}
$$

Finally we assume that the local characteristics of the diffusion satisfy the restricted isotropy condition stated in (0.4).

We recall that $\left(X_{t}\right)_{t \geq 0}$ denotes the canonical process on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$. We write $\left(\mathcal{F}_{t}\right)_{t \geq 0}$ and $\left(\theta_{t}\right)_{t \geq 0}$ for the respective canonical right-continuous filtration and canonical shift on $C\left(\mathbb{R}_{+}, \mathbb{R}^{\bar{d}}\right)$. We also write $H_{B}$ and $T_{U}$ for the respective entrance time of $X$ in the closed set $B \subseteq \mathbb{R}^{d}$ and exit time of $X$ from the open set $U \subseteq \mathbb{R}^{d}$ :

$$
\begin{equation*}
H_{B}=\inf \left\{u \geq 0, X_{u} \in B\right\}, T_{U}=\inf \left\{u \geq 0, X_{u} \notin U\right\} \tag{1.8}
\end{equation*}
$$

In view of (1.2) - (1.5), for any $\omega \in \Omega, x \in \mathbb{R}^{d}$, the martingale problem attached to $(a(\cdot, \omega), b(\cdot, \omega), x)$, (or alternatively to $L$ in (0.2), and $x$ ) is well-posed, cf. [26]. The corresponding law $P_{x, \omega}$ on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$, unique solution of the above martingale problem, describes the diffusion in the environment $\omega$ and starting from $x$. We write $E_{x, \omega}$ for the expectation under $P_{x, \omega}$. Under $P_{x, \omega},(X$.$) satisfies the stochastic differential equation$

$$
\left\{\begin{align*}
d X_{t} & =\sigma\left(X_{t}, \omega\right) d \beta_{t}+b\left(X_{t}, \omega\right) d t  \tag{1.9}\\
X_{0} & =x, P_{x, \omega} \text {-a.s. }
\end{align*}\right.
$$

where $\sigma(\cdot, \omega)=a(\cdot, \omega)^{\frac{1}{2}}$ and $\beta$. is some $d$-dimensional $\left(\mathcal{F}_{t}\right)$-Brownian motion under $P_{x, \omega}$.
The laws $P_{x, \omega}$ are sometimes called "quenched laws" of the diffusion in random environment. We also need the "annealed laws", $P_{x}, x \in \mathbb{R}^{d}$, that are the semi-direct products on $\Omega \times C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$ :

$$
\begin{equation*}
P_{x}=\mathbb{P} \times P_{x, \omega} . \tag{1.10}
\end{equation*}
$$

We denote with $E_{x}$ the corresponding expectations. These laws typically destroy the Markovian property of ( $X$. ) but restore translation invariance and isotropy:
(1.11) the law of $(X .+y)$ under $P_{x}$ equals that of $(X$.$) under P_{x+y}$, for $x, y \in \mathbb{R}^{d}$,
and for $r$ a rotation matrix preserving the union of coordinate axes of $\mathbb{R}^{d}$, and $x \in \mathbb{R}^{d}$,

$$
\begin{equation*}
\text { the law of }\left(r X_{.}\right) \text {under } P_{x} \text { equals that of }\left(X_{.}\right) \text {under } P_{r x} . \tag{1.12}
\end{equation*}
$$

We now turn to the description of spatial scales. We first choose

$$
\begin{equation*}
\beta \in\left(0, \frac{1}{2}\right] \tag{1.13}
\end{equation*}
$$

that will later appear as an exponent of Hölder-continuous functions, as well as

$$
\begin{equation*}
a \in\left(0, \frac{\beta}{1000 d}\right], \text { and } c_{0}>1, \text { with } 2 c_{0} \log \left(1+\frac{a}{2}\right)>1 \tag{1.14}
\end{equation*}
$$

Then for $L_{0} \geq 10^{a^{-1}}$, an integer multiple of 5 , we define $L_{n}, n \geq 0$, by induction via:

$$
\begin{equation*}
L_{n+1}=\ell_{n} L_{n} \text { with } \ell_{n}=5\left[L_{n}^{a} / 5\right], n \geq 0 \tag{1.15}
\end{equation*}
$$

and by convention we set $L_{-1}=1$. We also need the auxiliary scales

$$
\begin{equation*}
D_{n}=L_{n} \exp \left\{c_{0}\left(\log \log L_{n}\right)^{2}\right\}, \widetilde{D}_{n}=L_{n} \exp \left\{2 c_{0}\left(\log \log L_{n}\right)^{2}\right\}, n \geq 0 \tag{1.16}
\end{equation*}
$$

The proof of Theorem 1.1, when deriving controls on certain Hölder-norms at scale $L_{n+1}$, requires one to work at depth $m_{0}+2$ in scale $L_{n-m_{0}-1}$, see Sections 4 and 5 , with $m_{0} \geq 2$ determined by

$$
\begin{equation*}
(1+a)^{m_{0}-2} \leq 100<(1+a)^{m_{0}-1} . \tag{1.17}
\end{equation*}
$$

We can now introduce the probability kernels that enter the renormalization scheme. To this end we first define

$$
\begin{equation*}
X_{u}^{*}=\sup _{s \leq u}\left|X_{s}-X_{0}\right|, u \geq 0 \tag{1.18}
\end{equation*}
$$

as well as the $\left(\mathcal{F}_{t}\right)$-stopping times describing the first time $X$. travels a distance $\widetilde{D}_{n}$ from its starting point:

$$
\begin{equation*}
T_{n}=\inf \left\{u \geq 0, X_{u}^{*} \geq \widetilde{D}_{n}\right\}, n \geq 0 \tag{1.19}
\end{equation*}
$$

We can then consider $n \geq 0, \omega \in \Omega$, the probability kernels on $\mathbb{R}^{d}$

$$
\begin{equation*}
R_{n}(x, d y)=P_{x, \omega}\left[X_{L_{n}^{2}} \in d y\right], \widetilde{R}_{n}(x, d y)=P_{x, \omega}\left[X_{L_{n}^{2} \wedge T_{n}} \in d y\right] . \tag{1.20}
\end{equation*}
$$

In the renormalization scheme we compare $R_{n}$ and $\widetilde{R}_{n}$ to a Gaussian probability kernel $R_{n}^{0}$ that we now define. To this end we denote with $W_{x}$ the $d$-dimensional Wiener measure starting from $x \in \mathbb{R}^{d}$. Then for $n \geq 0$, we set

$$
\begin{equation*}
R_{n}^{0}(x, d y)=W_{x}\left[X_{\alpha_{n} L_{n}^{2}} \in d y\right], \widetilde{R}_{n}^{0}(x, d y)=W_{x}\left[X_{\left(\alpha_{n} L_{n}^{2}\right) \wedge T_{n}} \in d y\right] \tag{1.21}
\end{equation*}
$$

where the positive constant $\alpha_{n}$ is such that:

$$
\begin{equation*}
E_{0}\left[\left|X_{L_{n}^{2} \wedge T_{n}}\right|^{2}\right]=E^{W_{0}}\left[\left|X_{\alpha_{n} L_{n}^{2}}\right|^{2}\right]=\alpha_{n} d L_{n}^{2}, n \geq 0 \tag{1.22}
\end{equation*}
$$

To compare $R_{n}$ and $\widetilde{R}_{n}$ to $R_{n}^{0}$, we will use the kernels

$$
\begin{equation*}
S_{n}=R_{n}-R_{n}^{0}, \widetilde{S}_{n}=\widetilde{R}_{n}-R_{n}^{0}, n \geq 0, \omega \in \Omega \tag{1.23}
\end{equation*}
$$

The local drift and the compensated second moments at level $n$ at site $x$ in the environment $\omega$ are defined via:

$$
\begin{align*}
\widetilde{d}_{n}(x, \omega) & =\int(y-x) \widetilde{R}_{n}(x, d y)=\int(y-x) \widetilde{S}_{n}(x, d y),  \tag{1.24}\\
\widetilde{\gamma}_{n}^{i, j}(x, \omega) & =\int(y-x)_{i}(y-x)_{j} \widetilde{S}_{n}(x, d y), 1 \leq i, j \leq d .
\end{align*}
$$

In view of the translation invariance and isotropy of $X$. under the annealed measure, cf. (1.11), (1.12), and of (1.22), we see that

$$
\begin{equation*}
\mathbb{E}\left[\widetilde{d}_{n}(x, \omega)\right]=0, \mathbb{E}\left[\widetilde{\gamma}_{n}(x, \omega)\right]=0, \text { for } n \geq 0, x \in \mathbb{R}^{d} \tag{1.25}
\end{equation*}
$$

Note also that for $x \in \mathbb{R}^{d}, n \geq 0$,

$$
\begin{equation*}
\widetilde{S}_{n}(x, d y) \text { depends in a } \mathcal{G}_{\bar{B}\left(x, \widetilde{D}_{n}\right)} \text {-fashion on } \omega \tag{1.26}
\end{equation*}
$$

(see (1.6) for the notation), and in particular

$$
\begin{equation*}
\widetilde{d}_{n}(x, \omega), \widetilde{\gamma}_{n}(x, \omega) \text { are } \mathcal{G}_{\bar{B}\left(x, \widetilde{D}_{n}\right)} \text {-measurable . } \tag{1.27}
\end{equation*}
$$

The finite range dependence property (1.7), together with stationarity and (1.25) yields the fact that $\left(\widetilde{d}_{n}(x, \omega), \widetilde{\gamma}_{n}(x, \omega)\right)_{x \in \mathcal{V}}$ are i.i.d. centered variables under $\mathbb{P}$, whenever $\mathcal{V}$ is a collection of points of $\mathbb{R}^{d}$ with mutual distance at least $2 \widetilde{D}_{n}+R$. This will be especially useful in Section 4.

In what follows we will use various norms. For $p \in[1, \infty]$, we denote with $|f|_{p}$ the $L^{p}$-norm of a measurable scalar function $f$ on $\mathbb{R}^{d}$. We also consider as already mentioned in (0.13) the Hölder-norm of order $\beta$, cf. (1.13), in scale $L_{n}$ :

$$
\begin{equation*}
|f|_{(n)}=\sup _{x \in \mathbf{R}^{d}}|f(x)|+L_{n}^{\beta} \sup _{x \neq y} \frac{|f(x)-f(y)|}{|x-y|^{\beta}}, n \geq 0 . \tag{1.28}
\end{equation*}
$$

Note that for $f, g$ scalar functions on $\mathbb{R}^{d}$ :

$$
\begin{equation*}
|f g|_{(n)} \leq|f|_{(n)}|g|_{(n)}, n \geq 0 \tag{1.29}
\end{equation*}
$$

The operator norm corresponding to $|\cdot|_{(n)}$ is denoted with $\|\cdot\|_{n}$ :

$$
\begin{equation*}
\|A\|_{n}=\sup _{|f|_{(n)}=1}|A f|_{(n)} \tag{1.30}
\end{equation*}
$$

for $A$ a linear operator mapping the space of Hölder-continuous functions of order $\beta$ into itself.

In Section 4 we need to compute in an efficient way the $\|\cdot\|_{n+1}$-norm of certain operators entering the linearization of $S_{n+1}$ expressed in terms of $n$, for $n_{0}-m_{0}-1 \leq n \leq n_{0}$, cf. Theorem 1.1 for the notation. This is done with the help of wavelets. Namely we choose a scaling function $\varphi$ and a mother wavelet $\psi$, which are compactly supported on $\mathbb{R}$, of class $C^{4}$, cf. Daubechies [6, Chapters 5,6], Mallat [16, Chapter 7]. In particular $\varphi, \psi$ have unit $L^{2}$-norms and $\int_{\mathbb{R}} \psi(t) d t=0$, cf. [6, p. 153], (intuitively one can think of the Haar wavelets $\varphi(t)=1_{[0,1)}(t), \psi(t)=1_{\left[0, \frac{1}{2}\right)}(t)-1_{\left[\frac{1}{2}, 1\right)}(t)$, which of course do not fulfill the smoothness assumption we require). Attached to this choice we have a multiresolution approximation of $L^{2}(\mathbb{R})$, namely a decreasing sequence of closed subspaces $V_{j}, j \in \mathbb{Z}$, of $L^{2}(\mathbb{R})$ :

$$
\begin{equation*}
\cdots \subset V_{2} \subset V_{1} \subset V_{0} \subset V_{-1} \subset V_{-2} \cdots, \tag{1.31}
\end{equation*}
$$

with dyadic scaling sending one space into the next, $V_{-\infty}=L^{2}(\mathbb{R}), V_{\infty}=\{0\}$, and $\varphi(\cdot-k), k \in \mathbb{Z}$, an orthonormal basis of $V_{0}, \psi(\cdot-k), k \in \mathbb{Z}$, an orthonormal basis of the complement of $V_{0}$ in $V_{-1}$. Since we are interested in functions on $\mathbb{R}^{d}$, we write

$$
\begin{equation*}
\theta_{0}=\varphi, \theta_{1}=\psi \tag{1.32}
\end{equation*}
$$

and for $\alpha \in\{0,1\}^{d}$ and $x=\left(x_{1}, \ldots, x_{d}\right) \in \mathbb{R}^{d}$, we define:

$$
\begin{equation*}
\theta_{\alpha}(x)=\theta_{\alpha_{1}}\left(x_{1}\right) \ldots \theta_{\alpha_{d}}\left(x_{d}\right), \tag{1.33}
\end{equation*}
$$

as well as for $\ell \in \mathbb{Z}, p \in \mathbb{Z}^{d}$ :

$$
\begin{equation*}
\theta_{\alpha, \ell, p}(x)=\theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right) . \tag{1.34}
\end{equation*}
$$

In this way given any "top scale" $2^{j_{0}}$, we have an orthogonal basis of $L^{2}\left(\mathbb{R}^{d}\right)$ made of $\theta_{\alpha, \ell, p}, \ell \leq j_{0}, p \in \mathbb{Z}^{d}$, with $\alpha \neq 0$ if $\ell<j_{0}$, and any $f \in L^{2}\left(\mathbb{R}^{d}\right)$ can be expanded as

$$
\begin{equation*}
f(x)=\sum_{\substack{\ell \leq j_{0}, p \in \mathbb{Z}^{d} \\ \alpha \neq 0, \text { for } \ell<j_{0}}} c_{\alpha, \ell, p}^{j_{0}} \theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right) . \tag{1.35}
\end{equation*}
$$

For our purpose the interest of this expansion stems from the fact that with an adequate choice of $j_{0}\left(\right.$ i.e. $\left.2^{j_{0}} \approx L_{n}\right)$ the norm $|f|_{(n)}$ is comparable to $\sup \left\{\left|c_{\alpha, \ell, p}^{j_{0}}\right| 2^{\beta\left(j_{0}-\ell\right)} ; \ell \leq j_{0}\right.$, $p \in \mathbb{Z}^{d}, \alpha \neq 0$ for $\left.\ell<j_{0}\right\}$. This leads to effective estimates on $\|\cdot\|_{n}$, cf. Proposition A. 2 from the Appendix. These controls will be very useful in the proof of Lemmas 4.5 and 4.6.

To formulate the Hölder-norm controls that enters the induction assumption of Theorem 1.1 we need certain cut-off functions which we now describe. We consider the [ 0,1$]$-valued radial function:

$$
\begin{equation*}
\chi(x)=1 \wedge(2-|x|)_{+}, x \in \mathbb{R}^{d} \tag{1.36}
\end{equation*}
$$

so that $\chi=1$ on $\bar{B}(0,1), \chi=0$ on $B(0,2)^{c}$. For $u \geq 1, x \in \mathbb{R}^{d}, n \geq 0$, we also consider

$$
\begin{align*}
\chi_{u}(\cdot) & =\chi\left(\frac{\cdot}{u}\right), \text { as well as }  \tag{1.37}\\
\chi_{n, x}(\cdot) & =\chi_{10 \sqrt{d} L_{n}}(\cdot-x)=\chi\left(\frac{-x}{10 \sqrt{d} L_{n}}\right) . \tag{1.38}
\end{align*}
$$

Of special importance for us will be the control of the norm $\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n}$ to measure the closeness of $\widetilde{R}_{n}$ to $R_{n}^{0}$, for starting points in a neighborhood of size const $L_{n}$ of $x$, (we incidentally mention that $\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n}$ is finite, cf. Remark 2.6.2)).

We are now ready to describe the induction assumption we will propagate. Part of the induction assumption, cf. (1.47), expresses the fact that with "high probability", $\left\|\chi_{n, 0} \widetilde{S}_{n}\right\|_{n}$ is "small" and for starting points $|y| \leq 30 \sqrt{d} L_{n}$, the tail of $X_{L_{n}^{2}}^{*}$ under $P_{y, \omega}$ has exponential decay. More precisely we introduce for $\omega \in \Omega, n \geq 0$, the set

$$
\begin{align*}
\mathcal{B}_{n}(\omega)= & \left\{x \in L_{n} \mathbb{Z}^{d} ; \text { for }|y-x| \leq 30 \sqrt{d} L_{n}, P_{y, \omega}\left[X_{L_{n}^{2}}^{*} \geq v\right] \leq e^{-\frac{v}{D_{n}}}\right. \\
& \text { for } \left.v \geq D_{n}, \text { and }\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n} \leq L_{n}^{-\delta}\right\} \tag{1.39}
\end{align*}
$$

with $\delta$ a number slightly larger than $\frac{\beta}{8}$, specifically:

$$
\begin{equation*}
\delta=\frac{5}{32} \beta \tag{1.40}
\end{equation*}
$$

We will in particular propagate an upper bound on $\mathbb{P}\left[0 \notin \mathcal{B}_{n}(\omega)\right]$, cf. (1.47).
Another part of the induction assumption involves the control of traps in the medium. For $n \geq 0, x \in L_{n} \mathbb{Z}^{d}$, we write

$$
\begin{equation*}
C_{n}(x)=x+L_{n}[0,1]^{d}, C_{n}^{\prime}(x)=x+L_{n}\left(-\frac{1}{4}, \frac{5}{4}\right)^{d} \tag{1.41}
\end{equation*}
$$



Fig. 1: The boxes $C_{n}(x), C_{n}^{\prime}(x), C_{n, \gamma}(x)$
We then chop each of the $2 d$ faces of $\partial C_{n}(x)$ into $5^{(d-1)}$ closed $(d-1)$-dimensional cubes of side-length $L_{n} / 5$, see (1.15), and denote with $C_{n, \gamma}(x), 1 \leq \gamma \leq 2 d 5^{(d-1)}$, the resulting closed $d$-dimensional cubes obtained by "expanding" in the outwards normal direction to $\partial C_{n}(x)$ the above mentioned ( $d-1$ )-dimensional cubes, (with some specific labelling of the collection of cubes expressed by the index $\gamma$ ). We clearly have

$$
\begin{equation*}
C_{n, \gamma}(x) \subseteq C_{n}^{\prime}(x), \text { for } 1 \leq \gamma \leq 2 d 5^{(d-1)}, n \geq 0, x \in L_{n} \mathbb{Z}^{d} \tag{1.42}
\end{equation*}
$$

To measure the possible presence of traps in $C_{n}(x)$, we want to control how well the diffusion starting in the smaller box $C_{n}(x)$ travels to the boundary boxes $C_{n, \gamma}(x)$ without leaving the larger box $C_{n}^{\prime}(x)$, within time $L_{n}^{2}$. To this end we pick a number

$$
\begin{equation*}
\zeta \in(0,2), \text { with } \zeta^{-1} \geq \frac{1}{2}+d 3^{d+1} \tag{1.43}
\end{equation*}
$$

see also (3.85), and introduce for $n \geq 0, x \in L_{n} \mathbb{Z}^{d}$, $A \subseteq C_{n}(x), 1 \leq \gamma \leq 2 d 5^{(d-1)}$, the random variables measuring the presence and strength of traps:

$$
\begin{equation*}
J_{n, x, A, \gamma}(\omega)=\inf \left\{u \geq 0 ; \inf _{y \in A} P_{y, \omega}\left[H_{C_{n, \gamma}(x)} \leq L_{n}^{2} \wedge T_{C_{n}^{\prime}(x)}\right] \geq c_{1} L_{n}^{-\zeta u}\right\} \tag{1.44}
\end{equation*}
$$

where $c_{1} \in(0,1)$ is the constant depending on $d$ and $\nu$, see also above (3.67):

$$
\begin{aligned}
c_{1}=\frac{1}{4} \inf \{ & W_{x}\left[X_{u} \in B, u<T_{\left.\left(-\frac{9}{40}, \frac{49}{40}\right)^{d}\right]}\right] ; u \in\left[\frac{1}{40 \nu}, \frac{4 \nu}{10}\right], x \in\left[-\frac{1}{10}, \frac{11}{10}\right]^{d}, \\
& \left.\quad \text { and } B \text { is a closed cube with side-length } \frac{1}{10}, \text { contained in }\left[-\frac{1}{5}, \frac{6}{5}\right]^{d}\right\}>0 .
\end{aligned}
$$

We call $n$-admissible family, for $n \geq 0$, an arbitrary collection
$\left(u_{x}, A_{x}, \gamma_{x}\right)_{x \in \mathcal{A}}$, where $\mathcal{A}$ is a finite subset of $L_{n} \mathbb{Z}^{d}$, and for $x \in \mathcal{A}$, $u_{x}>0, \gamma_{x} \in\left\{1, \ldots, 2 d 5^{(d-1)}\right\}$, and $A_{x} \subseteq C_{n}(x)$ is a union of boxes
$C_{n-1}(z)$ (with the convention $L_{-1}=1$, when $n=0$, cf. below (1.15)),
such that $d_{\infty}\left(A_{x}, A_{x^{\prime}}\right) \geq 10 d L_{n-1}$, when $x \neq x^{\prime}$.
In the induction step we will propagate an upper bound on $\mathbb{P}\left[\right.$ for $\left.x \in \mathcal{A}, J_{n, x, A_{x}, \gamma_{x}} \geq u_{x}\right]$ for $n$-admissible families that will show that with overwhelming probability the variables in (1.44) vanish. We are now almost ready to state the main Theorem 1.1. We just need to introduce two numbers $M_{0}$ and $M$ that will respectively govern the estimates on $\mathbb{P}\left[0 \notin \mathcal{B}_{n}(\omega)\right]$ and on the tail of the variables in (1.44).

$$
\begin{equation*}
M_{0} \geq 100 d(1+a)^{m_{0}+2}, M \geq 1000 M_{0} \tag{1.46}
\end{equation*}
$$

Throughout this article we denote with $c$ a positive constant varying from place to place that solely depends on $d, K, \nu, R, \beta, a, c_{0}, \varphi, \psi, \zeta, M_{0}, M$, cf. (1.3), (1.4), (1.5), (1.13), (1.7), (1.14), (1.32), (1.43), (1.46). Any additional dependence of the constant will appear in the notation. So for instance if $\mu$ is a parameter, $c(\mu)$ denotes a positive constant depending solely on $\mu, d, K, \nu, R, \beta, a, c_{0}, \varphi, \psi, \zeta, M_{0}, M$.

Theorem 1.1. (main induction step)
There are positive constants $c_{2}, c$, such that for $L_{0} \geq c$, for $n_{0} \geq m_{0}+1$, (cf. (1.17)), if for all $0 \leq n \leq n_{0}$,

$$
\begin{equation*}
\mathbb{P}\left[0 \notin \mathcal{B}_{n}(\omega)\right] \leq L_{n}^{-M_{0}}, \tag{1.47}
\end{equation*}
$$

and for all $n$-admissible families $\left(u_{x}, A_{x}, \gamma_{x}\right)_{x \in \mathcal{A}}$,

$$
\begin{align*}
& \mathbb{P}\left[\text { for all } x \in \mathcal{A}, J_{n, x, A_{x}, \gamma_{x}} \geq u_{x}\right] \leq L_{n}^{-\bar{M}_{n} \sum_{x \in \mathcal{A}}\left(u_{x}+1\right)} \text {, with } \\
& \bar{M}_{n}=M \prod_{0 \leq j<n}\left(1-\frac{c_{2}}{\log L_{j}}\right) \tag{1.48}
\end{align*}
$$

and if, with $\delta$ as in (1.40),
i) $\frac{1}{2 \nu} \leq \alpha_{n} \leq 2 \nu, 0 \leq n \leq n_{0}$,
ii) $\left|\alpha_{n+1}-\alpha_{n}\right| \leq L_{n}^{-\left(1+\frac{9}{10}\right) \delta}, 0 \leq n<n_{0}$,
then the estimates (1.47), (1.48) hold with $n_{0}+1$ in place of $n_{0}$, and

$$
\begin{equation*}
\left|\alpha_{n_{0}+1}-\alpha_{n_{0}}\right| \leq L_{n_{0}}^{-\left(1+\frac{9}{10}\right) \delta} \tag{1.50}
\end{equation*}
$$

The proof of Theorem 1.1 is the scope of the next four sections. The crucial control is (1.47). In Section 2 we propagate the localization estimate contained in (1.47), that pertains to the tail behavior of $X_{L_{n}^{2}}^{*}$. In Section 3 we propagate the control on traps that appears in (1.48). It is in fact used in a rather special case, at the beginning of Section 5, cf. (5.3). As mentioned in the Introduction, the more detailed (1.48) enables the induction proof to function. In Section 4 we perform surgery on the environment at scale $L_{n_{0}^{\prime}}$, with $n_{0}^{\prime}=n_{0}-m_{0}-1$, and remove possible defects within distance const $L_{n_{0}+1}^{2}$ from the origin, which (in essence) belong to $L_{n_{0}^{\prime}} \mathbb{Z}^{d} \backslash \mathcal{B}_{n_{0}^{\prime}}(\omega)$, and show that with high probability this modified environment behaves very well up to scale $L_{n_{0}+1}$. In Section 5 we compare the true and modified environment, and show with the help of the smoothness estimates of Section 4, and the control on traps from (1.48) and Section 3, that one can repair the defects possibly present in the true environment.

We have already discussed our convention concerning positive constants above Theorem 1.1. We will use in the sequel the expression "for large $L_{0}$ " in place of "when $L_{0} \geq c$ ". We will recurrently use the shorthand notation

$$
\begin{equation*}
\kappa_{n}=\exp \left\{c\left(\log \log L_{n}\right)^{2}\right\}, n \geq 0 \tag{1.51}
\end{equation*}
$$

From now on we assume $L_{0} \geq 10$, large enough so that

$$
\begin{equation*}
L_{n}<D_{n}<\widetilde{D}_{n}<L_{n+1}, \text { for } n \geq 0 \tag{1.52}
\end{equation*}
$$

We close this section with some bounds on the Brownian semigroup and on the semigroup of diffusion in random environment. We write $\left(P_{t}\right)_{t \geq 0}$ for the Brownian semigroup and $p_{t}(x, y)$ for its transition density so that

$$
\begin{align*}
p_{t}(x, y) & =(2 \pi t)^{-\frac{d}{2}} \exp \left\{-\frac{|y-x|^{2}}{2 t}\right\}, t>0, x, y \in \mathbb{R}^{d}, \text { and }  \tag{1.53}\\
P_{t} f(x) & =\int p_{t}(x, y) f(y) d y, t>0  \tag{1.54}\\
& =f(x), t=0, \text { with } x \in \mathbb{R}^{d}, f \text { bounded measurable }
\end{align*}
$$

Note that $P_{t}, t \geq 0$, contracts the $|\cdot|_{(n)}$-norm and

$$
\begin{equation*}
\left\|P_{t}\right\|_{n}=1, \text { for } t \geq 0 \tag{1.55}
\end{equation*}
$$

Also for $\gamma=\left(\gamma_{1}, \ldots, \gamma_{d}\right)$ a multi-index (i.e. $\gamma_{i} \geq 0$, integer), $f$ bounded measurable, $x \in \mathbb{R}^{d}, t>0$, one has

$$
\begin{equation*}
\left|D^{\gamma}\left(P_{t} f\right)(x)\right| \leq \frac{c(\gamma)}{t^{\frac{\gamma \mid}{2}}} \exp \left\{-\frac{d(x, \operatorname{Supp} f)^{2}}{4 t}\right\}\left[\left(\frac{|f|_{1}}{t^{\frac{d}{2}}}\right) \wedge|f|_{\infty}\right] \tag{1.56}
\end{equation*}
$$

with $|\gamma|=\gamma_{1}+\cdots+\gamma_{d}$, (the estimate readily follows from the identity: $D_{x}^{\gamma} p_{t}(x, y)=$ $(-1)^{|\gamma|} t^{-\frac{d+|\gamma|}{2}} D^{\gamma} q\left(\frac{y-x}{\sqrt{t}}\right)$, with $\left.q(z)=\frac{1}{\sqrt{2 \pi}} e^{-\frac{|z|^{2}}{2}}\right)$.

The semigroup of the diffusion in the environment $\omega$

$$
\begin{equation*}
\left(P_{t, \omega} f\right)(x)=E_{x, \omega}\left[f\left(X_{t}\right)\right], t \geq 0, x \in \mathbb{R}^{d}, f \text { as in (1.54), } \tag{1.57}
\end{equation*}
$$

thanks to (1.3) - (1.5), is known to admit a density $p_{t, \omega}(x, y)$, cf. Friedman [9], p. 24, which satisfies for $0<t \leq 1, x, y \in \mathbb{R}^{d}$ :

$$
\begin{align*}
p_{t, \omega}(x, y) & \leq \frac{c}{t^{\frac{d}{2}}} \exp \left\{-\frac{c|y-x|^{2}}{t}\right\}  \tag{1.58}\\
\left|D_{x} p_{t, \omega}(x, y)\right| & \leq \frac{c}{t^{\frac{d+1}{2}}} \exp \left\{-\frac{c|y-x|^{2}}{t}\right\} \tag{1.59}
\end{align*}
$$

As a consequence we can bound the norm $\left\|P_{t}\right\|_{L^{\infty} \rightarrow(n)}$ of $P_{t}$ between $L^{\infty}\left(\mathbb{R}^{d}\right)$ and the space of $\beta$-Hölder-continuous functions endowed with the norm $|\cdot|_{(n)}$.

## Lemma 1.2.

$$
\begin{equation*}
\left\|P_{t, \omega}\right\|_{L^{\infty} \rightarrow(n)} \leq c L_{n}^{\beta}, \text { for } t \geq 1, n \geq 0, \omega \in \Omega \tag{1.60}
\end{equation*}
$$

Proof. First note that for $s \geq 0, \omega \in \Omega$,

$$
\begin{equation*}
\left|P_{s, \omega} f\right|_{\infty} \leq|f|_{\infty} \tag{1.61}
\end{equation*}
$$

Then with (1.59) and the above we see that

$$
\begin{equation*}
\left|P_{1, \omega} f(x)-P_{1, \omega} f(y)\right| \leq c(|x-y| \wedge 1)|f|_{\infty} \leq c L_{n}^{\beta}\left(\left|\frac{x-y}{L_{n}}\right|^{\beta} \wedge 1\right)|f|_{\infty} \tag{1.62}
\end{equation*}
$$

We thus find

$$
\begin{equation*}
\left|P_{1, \omega} f\right|_{(n)} \leq c L_{n}^{\beta}|f|_{\infty}, \tag{1.63}
\end{equation*}
$$

and writing for $t \geq 1, P_{t, \omega}=P_{1, \omega} P_{t-1, \omega}$, the claim (1.60) now follows from (1.61), (1.63).

## 2 Localization estimates

We keep the notations of the previous section and in particular of Theorem 1.1. We begin here the proof of Theorem 1.1, the principal aim of this section is to propagate to level $n_{0}+1$ the tail estimates on $X^{*}$ implicit in (1.47), see also (1.39). This is achieved in Proposition 2.5. We also derive controls in Proposition 2.5 which in particular imply that $S_{n_{0}+1}$ and $\widetilde{S}_{n_{0}+1}$ are typically close in $\left\|\|_{n_{0}+1}\right.$-norm. We begin with some additional notations. With $K$ from (1.3), (1.4), and $n \geq 0$, we define:

$$
\begin{equation*}
\mathcal{T}_{n}=\left(-2 K L_{n}^{2}, 2 K L_{n}^{2}\right)^{d} \tag{2.1}
\end{equation*}
$$

and also introduce for $\omega \in \Omega$, the modification of $\mathcal{B}_{n}(\omega)$ in (1.39), see (1.16) for notations:

$$
\begin{align*}
\widetilde{\mathcal{B}}_{n}(\omega)= & \left\{x \in L_{n} \mathbb{Z}^{d} ; \text { for }|y-x| \leq 30 \sqrt{d} L_{n}, P_{y, \omega}\left[X_{L_{n}^{2}}^{*} \geq v\right] \leq \exp \left\{-\frac{v}{D_{n}}\right\},\right.  \tag{2.2}\\
& \text { for } \left.D_{n} \leq v \leq \widetilde{D}_{n}, \text { and }\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n} \leq L_{n}^{-\delta}\right\} .
\end{align*}
$$

Note that for $n \geq 0, x \in L_{n} \mathbb{Z}^{d}$, the event $\left\{x \in \widetilde{\mathcal{B}}_{n}(\omega)\right\}$ unlike $\left\{x \in \mathcal{B}_{n}(\omega)\right\}$ has a local dependence:

$$
\begin{equation*}
\left\{x \in \widetilde{\mathcal{B}}_{n}(\omega)\right\} \in \mathcal{G}_{\bar{B}\left(x, \widetilde{D}_{n}+30 \sqrt{d} L_{n}\right)}, \text { (see (1.6) for the notation). } \tag{2.3}
\end{equation*}
$$

In the terminology introduced above (1.51), and the notations of (1.22), (1.24), one has

Lemma 2.1. There is a constant $\bar{c}>0$, such that for large $L_{0}$, for any $\omega \in \Omega, n \geq 0$, with $\alpha_{n} \leq 2 \nu, x \in L_{n} \mathbb{Z}^{d}$ with $\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n} \leq L_{n}^{-\delta}$, and $|y-x| \leq 10 \sqrt{d} L_{n}$ :

$$
\begin{equation*}
\left|\widetilde{d}_{n}(y, \omega)\right| \leq \bar{\kappa}_{n} L_{n}^{1-\delta},\left|\widetilde{\gamma}_{n}(y, \omega)\right| \leq \bar{\kappa}_{n} L_{n}^{2-\delta}, \text { with } \bar{\kappa}_{n}=\exp \left\{\bar{c}\left(\log \log L_{n}\right)^{2}\right\} \tag{2.4}
\end{equation*}
$$

Proof. For $y$ as above and $1 \leq i, j \leq d$, we define, cf. (1.37),

$$
\begin{align*}
f_{i}(z) & =\chi_{\widetilde{D}_{n}}(z-y) \frac{(z-y)_{i}}{L_{n}}, \text { and }  \tag{2.5}\\
f_{i, j}(z) & =f_{i}(z) f_{j}(z) \tag{2.6}
\end{align*}
$$

Observe that

$$
\begin{equation*}
\left|f_{i}\right|_{(n)} \leq \kappa_{n} \text { and }\left|f_{i, j}\right|_{(n)} \stackrel{(1.29)}{\leq} \kappa_{n} \tag{2.7}
\end{equation*}
$$

Further using that $f_{i}(z)=\left(\frac{z-y}{L_{n}}\right)_{i}$ for $|z-y| \leq \widetilde{D}_{n}$, and Gaussian estimates, see (1.53), (here the control on $\alpha_{n}$ comes in play), one finds that

$$
\begin{equation*}
\left|\frac{\widetilde{d}_{n}(y, \omega)_{i}}{L_{n}}-\left(\widetilde{S}_{n} f_{i}\right)(y)\right| \leq e^{-\kappa_{n}}, \quad\left|\frac{\widetilde{\gamma}_{n}^{i, j}(y, \omega)}{L_{n}^{2}}-\left(\widetilde{S}_{n} f_{i, j}\right)(y)\right| \leq e^{-\kappa_{n}} . \tag{2.8}
\end{equation*}
$$

Since $\chi_{n, x}(y)=1$, cf. (1.38), and $\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n} \leq L_{n}^{-\delta}$, cf. (2.2), the claim now follows ( $L_{0}$ is large).

We now turn to the localization estimates.
Proposition 2.2. For large $L_{0}$, if for $n \geq 0$, (1.47) and $\frac{1}{2 \nu} \leq \alpha_{n} \leq 2 \nu$ hold, then

$$
\begin{align*}
\mathbb{P}\left[\text { for }|y| \leq 30 \sqrt{d} L_{n+1}, P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right]\right. & \left.\leq \exp \left\{-\frac{v}{D_{n+1}}\right\}, \text { for } v \geq D_{n+1}\right] \\
& \geq 1-\frac{1}{10} L_{n+1}^{-M_{0}} \tag{2.9}
\end{align*}
$$

Proof. Using the exponential inequality for martingales, cf. Revuz-Yor [23], p. 145, for large $L_{0}, n \geq 0, \omega \in \Omega, v \geq 2 K L_{n+1}^{2}$, and arbitrary $y$ we find

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right] \leq c \exp \left\{-\frac{c v^{2}}{L_{n+1}^{2}}\right\} \leq c \exp \{-c v\} \leq \exp \left\{-\frac{v}{D_{n+1}}\right\} \tag{2.10}
\end{equation*}
$$

Hence for proving (2.9) we can restrict $v$ to

$$
\begin{equation*}
D_{n+1} \leq v<2 K L_{n+1}^{2} \tag{2.11}
\end{equation*}
$$

For such $v$ and $\omega \in \Omega$, we define

$$
\begin{align*}
\mathcal{B}_{n, v}(\omega)= & \left\{x \in L_{n} \mathbb{Z}^{d}, P_{y, \omega}\left[X_{L_{n}^{2}}^{*} \geq u\right] \leq \exp \left\{-\frac{u}{D_{n}}\right\}, \text { for } D_{n} \leq u \leq \frac{v}{100}\right.  \tag{2.12}\\
& \text { and } \left.\left|\widetilde{d}_{n}(y, \omega)\right| \leq \bar{\kappa}_{n} L_{n}^{1-\delta}, \text { for }|y-x| \leq 10 \sqrt{d} L_{n}\right\},
\end{align*}
$$

where $\bar{\kappa}_{n}$ appears in (2.4). As in (2.3) the local dependence of the event $\left\{x \in \mathcal{B}_{n, v}(\omega)\right\}$, for $x \in L_{n} \mathbb{Z}^{d}$, is expressed by

$$
\begin{equation*}
\left\{x \in \mathcal{B}_{n, v}(\omega)\right\} \in \mathcal{G}_{\bar{B}\left(x,\left(\frac{v}{100} \vee \widetilde{D}_{n}\right)+10 \sqrt{d} L_{n}\right)} . \tag{2.13}
\end{equation*}
$$

In particular with (1.7) and (2.11), we see that when $L_{0}$ is large,

$$
\begin{align*}
& \text { for } x, x^{\prime} \in L_{n} \mathbb{Z}^{d} \text {, with }\left|x-x^{\prime}\right| \geq \frac{v}{40},\left\{x \in \mathcal{B}_{n, v}(\omega)\right\} \text { and }  \tag{2.14}\\
& \left\{x^{\prime} \in \mathcal{B}_{n, v}(\omega)\right\} \text { are independent } .
\end{align*}
$$

We then introduce, see (2.1):

$$
\begin{equation*}
\Omega_{n, v}=\left\{\omega \in \Omega, \mathcal{T}_{n+1} \cap L_{n} \mathbb{Z}^{d} \cap \mathcal{B}_{n, v}^{c}(\omega) \subset B\left(x_{0}, \frac{v}{70}\right), \text { for some } x_{0} \in L_{n} \mathbb{Z}^{d}\right\} \tag{2.15}
\end{equation*}
$$

Observe that when $L_{0}$ is large, $n \geq 0, v$ as in (2.11),

$$
\begin{align*}
& \mathbb{P}\left[\Omega_{n, v}^{c}\right] \leq \mathbb{P}\left[\mathcal{T}_{n+1} \cap L_{n} \mathbb{Z}^{d} \cap \mathcal{B}_{n, v}^{c}(\omega) \text { has diameter } \geq \frac{2 v}{70}-\sqrt{d} L_{n}\right] \leq \\
& \mathbb{P}\left[\text { for some } x, x^{\prime} \in \mathcal{T}_{n+1} \cap L_{n} \mathbb{Z}^{d}, \text { with }\left|x-x^{\prime}\right| \geq \frac{v}{40}, x \text { and } x^{\prime} \notin \mathcal{B}_{n, v}(\omega)\right] \leq  \tag{2.16}\\
& \left(c L_{n+1}^{2} / L_{n}\right)^{2 d} L_{n}^{-2 M_{0}} \leq L_{n+1}^{4 d} L_{n}^{-2 M_{0}},
\end{align*}
$$

where we have used $\mathcal{B}_{n}(\omega) \subset \widetilde{\mathcal{B}}_{n}(\omega)$, and hence with $(2.2),(2.4), \mathcal{B}_{n}(\omega) \subset \mathcal{B}_{n, v}(\omega)$, as well as (1.47) and (2.14) in the last step. We now pick some $\omega \in \Omega_{n, v}$. We can find some $x_{0}(\omega) \in \mathcal{T}_{n+1} \cap L_{n} \mathbb{Z}^{d}$, such that

$$
\begin{equation*}
\mathcal{T}_{n+1} \cap L_{n} \mathbb{Z}^{d} \cap \mathcal{B}_{n, v}^{c}(\omega) \subseteq B\left(x_{0}(\omega), \frac{v}{70}\right) \tag{2.17}
\end{equation*}
$$

We introduce the successive entrance times $R_{i}$ and exit times $D_{i}$ of $X$. in $B\left(x_{0}, \frac{v}{50}\right)$ and out of $B\left(x_{0}, \frac{v}{40}\right)$, (see (1.8) for the notation):

$$
\begin{align*}
R_{1} & =H_{B\left(x_{0}, \frac{v}{50}\right)}, D_{1}=T_{B\left(x_{0}, \frac{v}{40}\right)} \circ \theta_{R_{1}}+R_{1}, \text { and for } i \geq 1,  \tag{2.18}\\
R_{i+1} & =R_{1} \circ \theta_{D_{i}}+D_{i}, D_{i+1}=D_{1} \circ \theta_{D_{i}}+D_{i},
\end{align*}
$$

so that

$$
\begin{equation*}
R_{1} \leq D_{1} \leq R_{2} \leq \cdots \leq \infty \tag{2.19}
\end{equation*}
$$

We first discuss the more complicated case where

$$
\begin{equation*}
\left|x_{0}(\omega)\right| \leq \frac{v}{2} \tag{2.20}
\end{equation*}
$$

Then for $|y| \leq 30 \sqrt{d} L_{n+1}$, we write for large $L_{0}$,

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right] \leq P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v, R_{1} \leq L_{n+1}^{2}\right]+P_{y, \omega}\left[T_{B\left(0, \frac{3}{4} v\right)}<R_{1} \wedge L_{n+1}^{2}\right] \tag{2.21}
\end{equation*}
$$

where we have used that $P_{y, \omega}$-a.s., $T_{B\left(0, \frac{3}{4} v\right)}<L_{n+1}^{2}$, on $\left\{X_{L_{n+1}^{2}}^{*} \geq v\right\}$. To bound the first term on the right-hand side of (2.21), we consider on the event $\left\{X_{L_{n+1}^{2}}^{*} \geq v, R_{1} \leq L_{n+1}^{2}\right\}$ the last exit time of $B\left(x_{0}, \frac{v}{40}\right)$ before $T_{B\left(0, \frac{3}{4} v\right)}\left(<L_{n+1}^{2}, P_{y, \omega}\right.$-a.s. on this event), and the integer part of this time. We then find:

$$
\begin{align*}
& P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v, R_{1} \leq L_{n+1}^{2}\right] \leq \\
& P_{y, \omega}\left[\text { for some } k \leq L_{n+1}^{2}, \sup _{u \in[k, k+1]}\left|X_{u}-X_{k}\right| \geq \frac{v}{100}\right]+  \tag{2.22}\\
& P_{y, \omega}\left[\bigcup_{m \leq L_{n+1}^{2}}\left(\left\{X_{m} \in K\left(x_{0}\right)\right\} \cap \theta_{m}^{-1}\left\{T_{B\left(0, \frac{3}{4} v\right)}<R_{1} \wedge L_{n+1}^{2}\right\}\right)\right], \text { with } \\
& m \text { integer and } K\left(x_{0}\right)=\partial B\left(x_{0}, \frac{v}{40}\right)+\bar{B}\left(0, \frac{v}{100}\right)
\end{align*}
$$

Using an exponential inequality as in (2.10) to bound the first term on the right-hand side of (2.22), we find:

$$
\begin{align*}
& P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v, R_{1} \leq L_{n+1}^{2}\right] \leq \\
& c L_{n+1}^{2}\left(\exp \left\{-c v^{2}\right\}+\sup _{z \in K\left(x_{0}\right)} P_{z, \omega}\left[T_{B\left(0, \frac{3}{4} v\right)}<R_{1} \wedge L_{n+1}^{2}\right]\right) . \tag{2.23}
\end{align*}
$$

For convenience we write $\mathcal{K}_{n}=\left\{k \geq 0 ; k L_{n}^{2}<R_{1} \wedge L_{n+1}^{2} \wedge T_{\mathcal{T}_{n+1}}\right\}$. Keeping in mind the last term of (2.21), we write for $|z| \leq 30 \sqrt{d} L_{n+1}$, or $z \in K\left(x_{0}\right)$ :

$$
\begin{align*}
& P_{z, \omega}\left[T_{B\left(0, \frac{3}{4} v\right)}<R_{1} \wedge L_{n+1}^{2}\right] \leq  \tag{2.24}\\
& P_{z, \omega}\left[\text { for some } k \in \mathcal{K}_{n}, \sup _{u \in\left[k L_{n}^{2},(k+1) L_{n}^{2}\right]}\left|X_{u}-X_{k L_{n}^{2}}\right| \geq \frac{v}{100}\right]+ \\
& P_{z, \omega}\left[\text { for each } k \in \mathcal{K}_{n}, \sup _{u \in\left[k L_{n}^{2},(k+1) L_{n}^{2}\right]}\right] X_{u}-X_{k L_{n}^{2}} \left\lvert\,<\frac{v}{100}\right., \text { and } \\
& \left.\qquad T_{B\left(0, \frac{3}{4} v\right)}<R_{1} \wedge L_{n+1}^{2} \wedge T_{\left.\mathcal{T}_{n+1}\right]}\right] \leq \\
& c \ell_{n}^{2} \exp \left\{-\frac{v .12),(2.17)}{100 D_{n}}\right\}+P_{z, \omega}\left[\text { for each } k \in \mathcal{K}_{n}, \sup _{u \in\left[k L_{n}^{2},(k+1) L_{n}^{2}\right]}\left|X_{u}-X_{k L_{n}^{2}}\right|<\frac{v}{100},\right. \\
& \left.\quad \text { and } X_{R_{1} \wedge L_{n+1}^{2} \wedge T_{\mathcal{T}_{n+1}}}>\frac{v}{5}\right] \leq \\
& c \ell_{n}^{2} \exp \left\{-\frac{v}{100 D_{n}}\right\}+P_{z, \omega}\left[\text { for each } k \in \mathcal{K}_{n}, \sup _{u \in\left[k L_{n}^{2},(k+1) L_{n}^{2}\right]}\left|X_{u}-X_{k L_{n}^{2}}\right|<\frac{v}{100},\right. \\
& \left.\quad \text { and for some } m \in \mathcal{K}_{n},\left|X_{m L_{n}^{2}}-z\right|>\frac{v}{10}\right] .
\end{align*}
$$

We now have to bound the last term of (2.24). To this end we will use an exponential estimate. But we first need the following
Lemma 2.3. If $Z$ is a random variable on some probability space such that

$$
\begin{align*}
& E\left[e^{Z}\right] \leq 2, E\left[e^{-Z}\right] \leq 2 \text { and }  \tag{2.25}\\
& E[Z]=0, \tag{2.26}
\end{align*}
$$

then for $L \geq 1$,

$$
\begin{equation*}
E\left[\exp \left\{\sqrt{\frac{\log 2}{2}} \frac{Z}{L}\right\}\right] \leq 2^{1 / L^{2}} \tag{2.27}
\end{equation*}
$$

Proof. For $\alpha \in(0,1]$ and $u \in \mathbb{R}$, one has the inequality

$$
\begin{equation*}
\alpha^{-2}\left(e^{\alpha u}-1-\alpha u\right) \leq e^{u}+e^{-u}-2, \tag{2.28}
\end{equation*}
$$

that can be verified by expanding both sides in Taylor series and using that $\sum_{k \geq 2, \text { even }} \frac{u^{k}}{k!} \geq$ $\sum_{k \geq 2 \text { odd }} \frac{u^{k}}{k!}$. Hence we find

$$
\begin{equation*}
e^{\alpha Z} \leq 1+\alpha Z+\alpha^{2}\left[e^{Z}+e^{-Z}-2\right] \tag{2.29}
\end{equation*}
$$

Substituting $\alpha=\sqrt{\frac{\log 2}{2}} L^{-1}$, and taking expectations we find with (2.25), (2.26), that the left-hand side of (2.27) is smaller than

$$
1+\frac{\log 2}{L^{2}} \leq \exp \left\{\frac{\log 2}{L^{2}}\right\} \leq 2^{1 / L^{2}}
$$

This proves (2.27).

The desired exponential estimate comes in the next lemma where $y^{\prime}$ plays the role of $X_{m L_{n}^{2}}$ in the last term of (2.24). For $u \geq 0$, we write

$$
\begin{equation*}
\psi_{u}(\cdot)=[-u \vee \cdot] \wedge u \tag{2.30}
\end{equation*}
$$

Lemma 2.4. There is a constant $c$ such that for $L_{0}$ large, if $x \in \mathcal{B}_{n, v}(\omega)$ and $\left|y^{\prime}-x\right| \leq$ $10 \sqrt{d} L_{n}$, then for any $e \in \mathbb{Z}^{d}$, with $|e|=1$,

$$
\begin{equation*}
E_{y^{\prime}, \omega}\left[\exp \left\{\frac{c}{\ell_{n} D_{n}}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)-E_{y^{\prime}, \omega}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)\right]\right]\right\}\right] \leq 2^{\ell_{n}^{-2}} \tag{2.31}
\end{equation*}
$$

Proof. In view of Lemma 2.3, we only need to prove that for some $c$ and all $e$ as above:

$$
\begin{equation*}
E_{y^{\prime}, \omega}\left[\exp \left\{\frac{c}{D_{n}}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)-E_{y^{\prime}, \omega}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)\right]\right]\right\}\right] \leq 2 \tag{2.32}
\end{equation*}
$$

To this end note that with a small enough $c$ one has

$$
\begin{align*}
& E_{y^{\prime}, \omega}\left[\exp \left\{\frac{c}{D_{n}} \psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)\right\}\right] \leq 1+E_{y^{\prime}, \omega}\left[\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e>0\right. \\
& \left.\int_{0}^{\frac{v}{100} \wedge\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e} \frac{c}{D_{n}} \exp \left\{\frac{c}{D_{n}} u\right\} d u\right] \stackrel{(2.12)}{\leq} 1+\int_{0}^{\frac{v}{100}} \frac{c}{D_{n}} \exp \left\{(c-1) \frac{u}{D_{n}}+1\right\} d u  \tag{2.33}\\
& \leq 1+\frac{c}{1-c} e \leq \sqrt{2}
\end{align*}
$$

Then observe that when $L_{0}$ is large:

$$
\begin{align*}
& \left|E_{y^{\prime}, \omega}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)\right]-\widetilde{d}_{n}\left(y^{\prime}, \omega\right) \cdot e\right| \stackrel{(1.24)}{=}  \tag{2.34}\\
& \left|E_{y^{\prime}, \omega}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-y^{\prime}\right) \cdot e\right)-\left(X_{L_{n}^{2} \wedge T_{n}}-y^{\prime}\right) \cdot e\right]\right|
\end{align*}
$$

and since the integrand vanishes when $T_{n}>L_{n}^{2}$, (because $\left.\frac{v}{100}>\widetilde{D}_{n}\right)$,

$$
\begin{aligned}
& \leq \frac{2 v}{100} P_{y^{\prime}, \omega}\left[T_{n} \leq L_{n}^{2}\right] \stackrel{(2.11),(2.12)}{\leq} c L_{n+1}^{2} \exp \left\{-\frac{\widetilde{D}_{n}}{D_{n}}\right\} \\
& \begin{array}{c}
(1.15),(1.16) \\
\leq
\end{array} \exp \left\{2(1+a) \log L_{n}-\exp \left\{c_{0}\left(\log \log L_{n}\right)^{2}\right\}\right\}
\end{aligned}
$$

Moreover with (2.12) we find:

$$
\begin{equation*}
\left|\widetilde{d}_{n}\left(y^{\prime}, \omega\right) \cdot e\right| \leq \bar{\kappa}_{n} L_{n}^{1-\delta} \tag{2.35}
\end{equation*}
$$

Hence where $L_{0}$ is large, combining (2.33) - (2.35), we obtain (2.32). This concludes the proof of Lemma 2.4.

With the same $c$ as in (2.31), introducing for $e \in \mathbb{Z}^{d}$, with $|e|=1$, and $m \geq 0$, the notation

$$
\begin{align*}
& \mathcal{E}_{e, m}=  \tag{2.36}\\
& \exp \left\{\frac{c}{\ell_{n} D_{n}} \sum_{0 \leq j<m}\left(\psi_{\frac{v}{100}}\left(\left(X_{(j+1) L_{n}^{2}}-X_{j L_{n}^{2}}\right) \cdot e\right)-E_{X_{j L_{n}^{2}, \omega}}\left[\psi_{\frac{v}{100}}\left(\left(X_{L_{n}^{2}}-X_{0}\right) \cdot e\right)\right]\right)\right\},
\end{align*}
$$

we see as an application of (2.31) and the Markov property that for $m<\ell_{n}^{2},|z| \leq$ $30 \sqrt{d} L_{n+1}$ or $z \in K\left(x_{0}\right), e$ as above:

$$
\begin{equation*}
E_{z, \omega}\left[m L_{n}^{2}<R_{1} \wedge T_{\mathcal{I}_{n+1}}, \mathcal{E}_{e, m}\right] \leq 2 \tag{2.37}
\end{equation*}
$$

Note also that for large $L_{0}$, for $0 \leq m<\ell_{n}^{2}, P_{z, \omega}$-a.s. on the event $\left\{\left|X_{m L_{n}^{2}}-z\right|>\frac{v}{10}\right.$, $m L_{n}^{2}<R_{1} \wedge T_{\mathcal{T}_{n+1}}$, and for $\left.0 \leq k<m, \sup _{u \in\left[k L_{n}^{2},(k+1) L_{n}^{2}\right]}\left|X_{u}-X_{k L_{n}^{2}}\right|<\frac{v}{100}\right\}$, for some $e$ as above, with (2.12) and (2.34), (2.35):

$$
\begin{aligned}
\mathcal{E}_{e, m} & \geq \exp \left\{\frac{c}{\ell_{n} D_{n}} \sum_{0 \leq j<m}\left[\left(X_{(j+1) L_{n}^{2}}-X_{j L_{n}^{2}}\right) \cdot e-2 \kappa_{n} L_{n}^{1-\delta}\right]\right\} \\
& \geq \exp \left\{\frac{c}{\ell_{n} D_{n}}\left(\frac{v}{10 d}-\kappa_{n} \ell_{n}^{2} L_{n}^{1-\delta}\right)\right\} \geq \exp \left\{\frac{c}{\ell_{n} D_{n}} v\right\},
\end{aligned}
$$

using (1.14), (1.40) and $v \geq D_{n+1}$, in view of (2.11), in the last step. It now follows from (2.37) that the last term of (2.24) is smaller than $2 \ell_{n}^{2} \exp \left\{-\frac{c}{\ell_{n} D_{n}} v\right\}$. Hence we see that when $L_{0}$ is large the left-hand side of (2.24) is smaller than $c \ell_{n}^{2}\left(\exp \left\{-\frac{v}{100 D_{n}}\right\}+\right.$ $\exp \left\{-\frac{c}{\ell_{n} D_{n}} v\right\}$ ).

Using this bound in (2.23) and on the last term of (2.21), (recall that $|z| \leq 30 \sqrt{d} L_{n+1}$ or $z \in K\left(x_{0}(\omega)\right)$ in (2.24)), we obtain for large $L_{0}$ and $|y| \leq 30 \sqrt{d} L_{n+1}$ :

$$
\begin{align*}
P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right] & \leq c L_{n+1}^{2}\left(\exp \left\{-c v^{2}\right\}+\ell_{n}^{2} \exp \left\{-\frac{v}{100 D_{n}}\right\}+\ell_{n}^{2} \exp \left\{-\frac{c v}{\ell_{n} D_{n}}\right\}\right) \\
& \leq \exp \left\{-\frac{10 v}{D_{n+1}}\right\} \tag{2.38}
\end{align*}
$$

where we have used in the last step that for large $L_{0}$

$$
\begin{align*}
\frac{D_{n+1}}{\ell_{n} D_{n}} & \stackrel{(1.15),(1.16)}{\geq} \exp \left\{c_{0}\left[\left(\log \log L_{n}+\log \left(1+\frac{a}{2}\right)\right)^{2}-\left(\log \log L_{n}\right)^{2}\right]\right\}  \tag{2.39}\\
& \geq \exp \left\{2 c_{0} \log \left(1+\frac{a}{2}\right) \log \log L_{n}\right\}
\end{align*}
$$

with $2 c_{0} \log \left(1+\frac{a}{2}\right)>1$, by (1.14), as well as $v \geq D_{n+1}$, in view of (2.11).
We now turn to the simpler case where unlike (2.20)

$$
\begin{equation*}
\left|x_{0}(\omega)\right|>\frac{v}{2} . \tag{2.40}
\end{equation*}
$$

Then for $|y| \leq 30 \sqrt{d} L_{n+1}, L_{0}$ being large, we write:

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right] \leq P_{y, \omega}\left[T_{B\left(0, \frac{v}{3}\right)}<R_{1} \wedge L_{n+1}^{2}\right] \leq \exp \left\{-\frac{10 v}{D_{n+1}}\right\} \tag{2.41}
\end{equation*}
$$

repeating similar bounds as in (2.24), (leading to (2.38)). We now define, cf. (2.15),

$$
\begin{equation*}
\Omega_{n}=\bigcap_{m \geq 0 ; 10^{m}} \Omega_{n, 10^{m} D_{n+1}<2 K L_{n+1}^{2}}, \tag{2.42}
\end{equation*}
$$

and observe that for $\omega \in \Omega_{n}, v \in\left[D_{n+1}, 2 K L_{n+1}^{2}\right),|y| \leq 30 \sqrt{d} L_{n+1}$,

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v\right] \leq P_{y, \omega}\left[X_{L_{n+1}^{2}}^{*} \geq v_{m}\right] \stackrel{(2.38),(2.41)}{\leq} \exp \left\{-\frac{10 v_{m}}{D_{n+1}}\right\} \leq \exp \left\{-\frac{v}{D_{n+1}}\right\} \tag{2.43}
\end{equation*}
$$

with the notation $10^{m} D_{n+1}=v_{m} \leq v<10 v_{m}$.
In addition from (2.16) we deduce that when $L_{0}$ is large

$$
\begin{equation*}
\mathbb{P}\left[\Omega_{n}^{c}\right] \leq\left(\left[\log \left(\frac{2 K L_{n+1}^{2}}{D_{n+1}}\right) / \log 10\right]+1\right) c L_{n+1}^{4 d} L_{n}^{-2 M_{0}} \leq \frac{1}{10} L_{n+1}^{-M_{0}} \tag{2.44}
\end{equation*}
$$

since $2 M_{0}(1+a)^{-1}>M_{0}+4 d+1$, by (1.14), (1.46). Combining (2.10), (2.43), (2.44), we see that (2.9) is proved.

We will now conclude this section with an estimate on $\left\|\chi_{n, x}\left(S_{n}-\widetilde{S}_{n}\right)\right\|_{n}$ that will be repeatedly used in the sequel. We refer to (1.23), (1.30), (1.38) for the notations.

Proposition 2.5. Given $\kappa_{n}^{0}$ as in (1.51), for $L_{0}$ large, for any $n \geq 0$, $\omega \in \Omega$, if $x \in L_{n} \mathbb{Z}^{d}$ is such that for $|y-x| \leq 30 \sqrt{d} L_{n}$,

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n}^{2}}^{*} \geq \frac{\widetilde{D}_{n}}{2}\right] \leq e^{-\kappa_{n}^{0}} \tag{2.45}
\end{equation*}
$$

then there exists a $\kappa_{n}$ as in (1.51) such that

$$
\begin{equation*}
\left\|\chi_{n, x}\left(S_{n}-\widetilde{S}_{n}\right)\right\|_{n} \leq e^{-\kappa_{n}} \tag{2.46}
\end{equation*}
$$

Proof. We use the shorthand notation

$$
\begin{align*}
& \Delta_{n}=S_{n}-\widetilde{S}_{n}, \text { so that }  \tag{2.47}\\
& \Delta_{n} g(z) \stackrel{(1.23)}{=} E_{z, \omega}\left[g\left(X_{L_{n}^{2}}\right)-g\left(X_{L_{n}^{2} \wedge T_{n}}\right), T_{n}<L_{n}^{2}\right]
\end{align*}
$$

Note that for $f$ with $|f|_{(n)} \leq 1$, and $x, y$ as above (2.45),

$$
\begin{equation*}
\left|\Delta_{n} f(y)\right| \leq 2 P_{y, \omega}\left[T_{n}<L_{n}^{2}\right] \stackrel{(2.45)}{\leq} 2 e^{-\kappa_{n}^{0}} \tag{2.48}
\end{equation*}
$$

So when $L_{0}$ is large, we find that for $y, y^{\prime}$ in $B\left(x, 21 \sqrt{d} L_{n}\right)$, with $\left|y-y^{\prime}\right| \geq e^{-\kappa_{n}}$,

$$
\begin{equation*}
\left|\Delta_{n} f(y)-\Delta_{n} f\left(y^{\prime}\right)\right| \leq 2 e^{-\kappa_{n}^{0}} \leq L_{n}^{\beta}\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} e^{-\kappa_{n}} \leq\left|\frac{y-y^{\prime}}{L_{n}}\right| e^{-\kappa_{n}} \tag{2.49}
\end{equation*}
$$

(see above (1.51) for the convention we use, and we are only interested in $y, y^{\prime} \in B\left(x, 21 \sqrt{d} L_{n}\right)$ because $\chi_{n, x}$ is supported in $\bar{B}\left(x, 20 \sqrt{d} L_{n}\right)$, as follows from (1.38)).

We now consider for $\kappa_{n}$ as above (2.49),

$$
\begin{equation*}
\left|y-y^{\prime}\right| \leq e^{-\kappa_{n}} \tag{2.50}
\end{equation*}
$$

and write

$$
\begin{align*}
& \left|\Delta_{n} f(y)-\Delta_{n} f\left(y^{\prime}\right)\right| \leq a_{1}+a_{2}, \text { with } \\
& a_{1}=\left|E_{y^{\prime}, \omega}\left[f\left(X_{L_{n}^{2} \wedge T_{y^{\prime}}}\right)-f\left(X_{L_{n}^{2} \wedge T_{y}}\right)\right]\right|  \tag{2.51}\\
& a_{2}=\left|E_{y, \omega}\left[f\left(X_{L_{n}^{2}}\right)-f\left(X_{L_{n}^{2} \wedge T_{y}}\right)\right]-E_{y^{\prime}, \omega}\left[f\left(X_{L_{n}^{2}}\right)-f\left(X_{L_{n}^{2} \wedge T_{y}}\right)\right]\right|
\end{align*}
$$

and $T_{y}=T_{B\left(y, \widetilde{D}_{n}\right)}, T_{y^{\prime}}=T_{B\left(y^{\prime}, \widetilde{D}_{n}\right)}$ in the notations of (1.8). Writing

$$
\begin{equation*}
\tau=T_{y} \wedge T_{y^{\prime}} \tag{2.52}
\end{equation*}
$$

it follows from the strong Markov property at time $\tau$, with hopefully obvious notations, that

$$
\begin{align*}
& a_{1} \leq\left|E_{y^{\prime}, \omega}\left[T_{y^{\prime}}=\tau<L_{n}^{2} \wedge T_{y}, E_{X_{T_{y^{\prime}}, \omega}}\left[f\left(X_{T_{y} \wedge\left(L_{n}^{2}-\tau\right)}\right)-f\left(X_{0}\right)\right]\right]\right|+ \\
& \quad\left|E_{y^{\prime}, \omega}\left[T_{y}=\tau<L_{n}^{2} \wedge T_{y^{\prime}}, E_{X_{T_{y}, \omega}}\left[f\left(X_{T_{y^{\prime}} \wedge\left(L_{n}^{2}-\tau\right)}\right)-f\left(X_{0}\right)\right]\right]\right|  \tag{2.53}\\
&\left.\stackrel{\text { def }}{=} b_{1}+b_{2}, \text { (the inner expectations do not integrate } \tau\right) .
\end{align*}
$$

We will now bound $b_{1}, b_{2}$ being handled similarly. To this end we consider $z^{\prime} \in \partial B\left(y^{\prime}, \widetilde{D}_{n}\right) \cap$ $B\left(y, \widetilde{D}_{n}\right),\left(z^{\prime}\right.$ plays the role of $\left.X_{T_{y^{\prime}}}\right), 0 \leq u \leq\left(L_{n}^{2}-\tau\right)_{+}$, and $\mathcal{H}$ the half-space $\{z \in$ $\left.\mathbb{R}^{d} ; z \cdot \ell \geq v\right\}$, with $\ell$ the unit vector in the direction $z^{\prime}-y^{\prime}, v=z^{\prime} \cdot \ell+\left|y^{\prime}-y\right|$. So $d\left(\mathcal{H}, \bar{B}\left(y^{\prime}, \widetilde{D}_{n}\right)\right)=\left|y-y^{\prime}\right|$ in the notation (1.1), and $B\left(y, D_{n}\right) \subset \mathcal{H}^{c}$. We will use the shorthand notation, cf. (1.8), $H=H_{\mathcal{H}}$, and note that

$$
\begin{align*}
& E_{z^{\prime}, \omega}\left[\left|X_{T_{y} \wedge u}-X_{0}\right|^{\beta} \wedge 2\right] \leq  \tag{2.54}\\
& 2 P_{z^{\prime}, \omega}\left[H>\left|y^{\prime}-y\right|\right]+E_{z^{\prime}, \omega}\left[H \leq\left|y^{\prime}-y\right|,\left|X_{T_{y} \wedge u}-X_{0}\right|^{\beta} \wedge 2\right]
\end{align*}
$$

To bound the right-hand side of $(2.54)$, we first note that under $P_{z^{\prime}, \omega},\left(X_{s}-X_{0}\right) \cdot \ell$ admits the semimartingale decomposition

$$
\begin{equation*}
\left(X_{s}-X_{0}\right) \cdot \ell=M_{s}+A_{s}, s \geq 0 \tag{2.55}
\end{equation*}
$$

where in view of (1.3) - (1.5), for some $c>1$,

$$
\begin{equation*}
\frac{1}{c} s \leq\langle M\rangle_{s} \leq c s, \quad\left|A_{s}\right| \leq c s, s \geq 0 \tag{2.56}
\end{equation*}
$$

Observe also that with $c$ as above,

$$
\begin{equation*}
P_{z^{\prime}, \omega} \text {-a.s., } T_{y} \leq H \leq \widetilde{H} \stackrel{\text { def }}{=} \inf \left\{s \geq 0, M_{s} \geq c s+\left|y^{\prime}-y\right|\right\} \tag{2.57}
\end{equation*}
$$

As a result we find that

$$
\begin{align*}
P_{z^{\prime}, \omega}\left[H \leq\left|y^{\prime}-y\right|\right] & \geq P_{z^{\prime}, \omega}\left[\widetilde{H} \leq\left|y^{\prime}-y\right|\right] \geq P_{z^{\prime}, \omega}\left[\sup _{s \leq\left|y^{\prime}-y\right|} M_{s} \geq(c+1)\left|y^{\prime}-y\right|\right] \\
& \geq W\left[\sup _{s \leq c\left|y-y^{\prime}\right|} B_{s} \geq(c+1)\right] \stackrel{\text { scaling }}{=} W\left[\sup _{s \leq 1} B_{s} \geq c\left|y-y^{\prime}\right|^{\frac{1}{2}}\right]  \tag{2.58}\\
& =1-\int_{-c\left|y^{\prime}-y\right|^{\frac{1}{2}}}^{c\left|y^{\prime}-y\right|^{\frac{1}{2}}} e^{-\frac{v^{2}}{2}} \frac{d v}{\sqrt{v \pi}} \geq 1-c\left|y^{\prime}-y\right|^{\frac{1}{2}}
\end{align*}
$$

where $B$. denotes the canonical one-dimensional Brownian motion, $W$ the Wiener measure, and we have used time-change together with (2.56). This yields a bound on the first term in the right-hand side of (2.54). For the second term we note that with $c$ as in (2.56), we can define

$$
\begin{equation*}
\bar{H}=\inf \left\{s \geq 0, M_{s}=(c+1)\left|y^{\prime}-y\right|\right\} \tag{2.59}
\end{equation*}
$$

and $P_{z^{\prime}, \omega^{-}}$a.s. on the event $\left\{H \leq\left|y^{\prime}-y\right|\right\}$, one has $T_{y} \leq H \leq \bar{H}$, and hence

$$
\begin{align*}
E_{z^{\prime}, \omega}\left[H \leq\left|y^{\prime}-y\right|,\left|X_{T_{y} \wedge u}-X_{0}\right|^{\beta} \wedge 2\right] & \stackrel{(1.9)}{\leq}  \tag{2.60}\\
c\left|y^{\prime}-y\right|^{\beta}+E_{z^{\prime}, \omega}\left[H \leq\left|y^{\prime}-y\right|, \sup _{s \leq \bar{H}}\left|\int_{0}^{s} \sigma\left(X_{v}, \omega\right) d \beta_{v}\right|^{\beta}\right] & \leq \\
c\left|y^{\prime}-y\right|^{\beta}+c E_{z^{\prime}, \omega}\left[\bar{H}^{\frac{\beta}{2}}\right] & \leq c\left|y^{\prime}-y\right|^{\beta},
\end{align*}
$$

using Burkholder-Davis-Inequalities, cf. Karatzas-Shreve [11, p. 166], then once again a representation of $M$. as a time change of Brownian motion together with scaling, and the fact that moments of order less than $\frac{1}{2}$ of the hitting time of 1 by Brownian motion are finite, cf. [11, p. 96]. We can now collect (2.58), (2.60) to bound the left-hand side of (2.54). Coming back to the first line of (2.53), since $|f|_{(n)} \leq 1$, and $\beta \leq \frac{1}{2}$, cf. (1.13), we find (recall $\tau$ is not integrated in the inner expectation)

$$
\begin{align*}
b_{1} & \leq E_{y^{\prime}, \omega}\left[T_{y^{\prime}}=\tau<L_{n}^{2} \wedge T_{y}, E_{X_{T_{y^{\prime}}}, \omega}\left[\left|X_{T_{y} \wedge\left(L_{n}^{2}-\tau\right)}-X_{0}\right|^{\beta} \wedge 2\right]\right] \\
& \leq c\left|y-y^{\prime}\right|^{\beta} P_{y^{\prime}, \omega}\left[\tau=T_{y^{\prime}}<L_{n}^{2}\right] \stackrel{(2.45)}{\leq} c\left|y-y^{\prime}\right|^{\beta} e^{-\kappa_{n}^{0}} \leq\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} e^{-\kappa_{n}} . \tag{2.61}
\end{align*}
$$

A similar bound can be proved for $b_{2}$.
We then turn to the bound on $a_{2}$ in (2.51). We use the shorthand notation

$$
\begin{equation*}
t_{0}=\left(\log \left|y^{\prime}-y\right|\right)^{-2},(\operatorname{recall}(2.50)), \tag{2.62}
\end{equation*}
$$

and denote with $q_{t, \omega}\left(z, z^{\prime}\right)$ the sub-probability density of the diffusion in the environment $\omega$, killed when exiting the ball $B(y, 10)$, at time $t>0$, when starting in $z \in B(y, 10)$. We now find that

$$
\begin{align*}
& 1-\int q_{t_{0}, \omega}(y, z) \wedge q_{t_{0}, \omega}\left(y^{\prime}, z\right) d z \leq 1-\int q_{t_{0}, \omega}(y, z) d z+ \\
& \int\left|q_{t_{0}, \omega}(y, z)-q_{t_{0}, \omega}\left(y^{\prime}, z\right)\right| d z \leq 1-\int q_{t_{0}, \omega}(y, z) d z+ \\
& \int\left|p_{t_{0}, \omega}(y, z)-p_{t_{0}, \omega}\left(y^{\prime}, z\right)\right| d z+1-\int q_{t_{0}, \omega}(y, z) d z+  \tag{2.63}\\
& 1-\int q_{t_{0}, \omega}\left(y^{\prime}, z\right) d z \leq c e^{-\frac{c}{t_{0}}}+c\left|\frac{y-y^{\prime}}{\sqrt{t_{0}}}\right| \stackrel{(2.62)}{\leq} c\left|\frac{y-y^{\prime}}{\sqrt{t_{0}}}\right|
\end{align*}
$$

for large $L_{0}$, using (1.59) and standard estimates.
With the help of (2.63), we can construct on some auxiliary probability space two processes $Y$. and $Y^{\prime}$. with same laws as $X$. under $P_{y, \omega}$ and $P_{y^{\prime}, \omega}$ such that

$$
\begin{align*}
& P[G] \geq 1-c\left|\frac{y-y^{\prime}}{\sqrt{t_{0}}}\right|, \text { with }  \tag{2.64}\\
& G=\left\{Y_{u}=Y_{u}^{\prime} \text { for } u \geq t_{0}, \text { and } Y \text { and } Y^{\prime} \text { do not exit } B(y, 10) \text { up to time } t_{0}\right\} .
\end{align*}
$$

We now see that with a slight abuse of notations, when $L_{0}$ is large:

$$
\begin{align*}
& a_{2} \leq\left|E\left[G^{c}, f\left(Y_{L_{n}^{2}}\right)-f\left(Y_{L_{n}^{2} \wedge T_{B\left(y, \widetilde{D}_{n}\right)}(Y)}\right)-\left(f\left(Y_{L_{n}^{2}}^{\prime}\right)-f\left(Y_{L_{n}^{2} \wedge T_{B\left(y, \widetilde{D}_{n}\right)}\left(Y^{\prime}\right)}^{\prime}\right)\right)\right]\right| \\
& \leq 4 P\left[G^{c}, T_{B\left(y, \widetilde{D}_{n}\right)}(Y)<L_{n}^{2} \text { or } T_{B\left(y, \widetilde{D}_{n}\right)}\left(Y^{\prime}\right)<L_{n}^{2}\right]  \tag{2.65}\\
& \text { Hölder,(2.45)} \\
& \quad \leq\left[G^{c}\right]^{\frac{1+\beta}{2}} e^{-\kappa_{n}} \stackrel{(2.62),(2.64)}{\leq}\left|y-y^{\prime}\right|^{\beta} e^{-\kappa_{n}} \leq\left|\frac{y^{\prime}-y}{L_{n}}\right|^{\beta} e^{-\kappa_{n}} .
\end{align*}
$$

Collecting the bounds (2.51), (2.53), (2.61), (2.65), together with (2.49), we see that when $L_{0}$ is large, for $y, y^{\prime}$ in $B\left(x, 21 \sqrt{d} L_{n}\right)$,

$$
\begin{equation*}
\left|\Delta_{n} f(y)-\Delta_{n} f\left(y^{\prime}\right)\right| \leq\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} e^{-\kappa_{n}} \tag{2.66}
\end{equation*}
$$

This together with (2.48) and (1.38) readily implies (2.46), (see also (A.4) - (A.6) of the Appendix).

## Remark 2.6.

1) We have used the assumption $\beta \leq \frac{1}{2}$, cf. (1.12), in the estimate (2.58).
2) Note that the estimates on (2.51), and (1.60) can also be used to show that for $\omega \in \Omega, n \geq 0, x \in L_{n} \mathbb{Z}^{d}$,

$$
\begin{equation*}
\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n} \leq c L_{n}^{\beta} \tag{2.67}
\end{equation*}
$$

## 3 Controlling traps

We continue the proof of Theorem 1.1. The main objective in this section is to propagate "at level $n_{0}+1$ " the estimate (1.48), and this comes in Proposition 3.3. As mentioned in the Introduction and in Section I below Theorem 1.1, the main purpose of the control (1.48) on the tails of the variables in (1.44) measuring the strength of traps, is to later obtain the estimate (5.3), when "repairing defects". This only involves a small portion of (1.48), but (1.48) is there to let the induction proof function. As a preparation for our main task we first construct certain couplings of the diffusion in random environment with Brownian motion of variance $\alpha_{n}$, cf. (1.22), at times $k L_{n}^{2}, k \geq 0$. These couplings will be very handy later in this section as well as in Section 6. We begin with some notations. We denote with $d_{n, \beta}(\cdot, \cdot)$ the distance function on $\mathbb{R}^{d}$ :

$$
\begin{equation*}
d_{n, \beta}\left(y, y^{\prime}\right)=\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta}, y, y^{\prime} \in \mathbb{R}^{d}, n \geq 0 \tag{3.1}
\end{equation*}
$$

We define for $\nu, \nu^{\prime}$ probabilities on $\mathbb{R}^{d}$, for which

$$
\begin{gather*}
\int|y|^{\beta} \nu(d y)<\infty, \int|y|^{\beta} \nu^{\prime}(d y)<\infty,  \tag{3.2}\\
D_{n, \beta}\left(\nu, \nu^{\prime}\right)=\sup \left\{\left|\int f d \nu-\int f d \nu^{\prime}\right| ; \text { where } f \text { on } \mathbb{R}^{d}\right. \text { is such that } \\
\left.\left|f(y)-f\left(y^{\prime}\right)\right| \leq d_{n, \beta}\left(y, y^{\prime}\right), \text { for } y, y^{\prime} \in \mathbb{R}^{d}\right\} \\
=\inf \left\{\int_{\mathbb{R}^{d} \times \mathbb{R}^{d}} d_{n, \beta}\left(y, y^{\prime}\right) \rho\left(d y, d y^{\prime}\right) ; \text { with } \rho\right. \text { probability }  \tag{3.3}\\
\text { having } \left.\nu, \nu^{\prime} \text { as first and second marginals }\right\},
\end{gather*}
$$

where the last equality results from the Kantorovich-Rubinstein Theorem, cf. Dudley [8, Theorem 20.1]. The function $D_{n, \beta}$ is sometimes called Kantorovich-Rubinstein or

Vasserstein distance. We now consider a continuous function $h$ with values in $[0,1]$, and for $\omega \in \Omega, n \geq 0$, define the probability kernel on $\mathbb{R}^{d}$

$$
\begin{equation*}
\widetilde{R}_{n, h}(x, d y)=R_{n}^{0}(x, d y)+h(x) \widetilde{S}_{n}(x, d y), \text { cf. }(1.21),(1.23), \tag{3.4}
\end{equation*}
$$

(so when $h \equiv 0, \widetilde{R}_{n, h}=R_{n}^{0}$, and when $h \equiv 1, \widetilde{R}_{n, h}=\widetilde{R}_{n}$ ).
We are now ready to state and prove the above mentioned result concerning coupling measures.

Proposition 3.1. Let $h$ be a continuous $[0,1]$-valued function on $\mathbb{R}^{d}, \omega \in \Omega$, and $n \geq 0$ such that $\frac{1}{2 \nu} \leq \alpha_{n} \leq 2 \nu$. Then for $y \in \mathbb{R}^{d}$, there is a measure $Q_{n, y}$ on the canonical space $\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)^{\mathbf{N}}$ endowed with the canonical $\sigma$-algebra and the canonical processes $\bar{X}_{k}, k \geq 0$, $\bar{X}_{k}^{0}, k \geq 0$, such that
under $Q_{n, y}, \bar{X}_{k}, k \geq 0$, (resp. $\bar{X}_{k}^{0}, k \geq 0$ ) has the law of the Markov chain on $\mathbb{R}^{d}$, starting at $y$ with transition kernel $\widetilde{R}_{n, h}$ (resp. $R_{n}^{0}$ )
and for $k_{0} \geq 1, \gamma>0$,

$$
\begin{equation*}
Q_{n, y}\left[\left|\bar{X}_{k}-\bar{X}_{k}^{0}\right| \geq \gamma, \text { for some } k \leq k_{0}\right] \leq k_{0}^{2}\left(\frac{\gamma}{L_{n}}\right)^{-\beta}\left(\kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}}\right) \tag{3.6}
\end{equation*}
$$

with $\Gamma_{n, h}=\sup _{x \in L_{n} \mathbb{Z}^{d}: \chi_{n, x} h \neq 0}\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n}$.
Remark 3.2. Note that under $Q_{n, y}$ above, $\left(\bar{X}_{k}^{0}\right)_{k \geq 0}$ has same law as $\left(X_{\alpha_{n} k L_{n}^{2}}\right)_{k \geq 0}$ under $W_{y}$, the Wiener measure starting from $y$, cf. above (1.21). The inequality (3.6) highlights one of the interests in controlling the norms $\|\cdot\|_{n}$.

Proof of Proposition 3.1: For $z \in \mathbb{R}^{d}$, denote with $K_{z}$ the non-empty compact subset of $M_{1}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)$, the set of probability measures on $\mathbb{R}^{d} \times \mathbb{R}^{d}$, endowed with the topology of weak convergence

$$
\begin{align*}
K_{z}= & \left\{\rho \in M_{1}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right) ; \rho \text { has marginals } R_{n, h}(z, \cdot) \text { and } R_{n}^{0}(z, \cdot),\right. \\
& \text { and } \left.D_{n, \beta}\left(\widetilde{R}_{n, h}(z, \cdot), R_{n}^{0}(z, \cdot)\right)=\int d_{n, \beta}\left(z_{1}, z_{2}\right) \rho\left(d z_{1}, d z_{2}\right)\right\} . \tag{3.7}
\end{align*}
$$

Observe that for any sequences $z_{i}, \rho_{i}, i \geq 1$, with $\rho_{i} \in K_{z_{i}}$, for $i \geq 1$, and $z_{i}$ converging to $z_{\infty}, \rho_{i}$ is tight and has a limit point $\rho_{\infty}$ such that:

$$
\begin{align*}
\int d_{n, \beta}\left(z_{1}, z_{2}\right) \rho_{\infty}\left(d z_{1}, d z_{2}\right) & \leq \liminf _{i} D_{n, \beta}\left(\widetilde{R}_{n, h}\left(z_{i}, \cdot\right), R_{n}^{0}\left(z_{i}, \cdot\right)\right)  \tag{3.8}\\
& =D_{n, \beta}\left(\widetilde{R}_{n, h}\left(z_{\infty}, \cdot\right), R_{n}^{0}\left(z_{\infty}, \cdot\right)\right)
\end{align*}
$$

as follows straightforwardly by applying the triangle inequality satisfied by $D_{n, \beta}$, as well as (2.67) and (3.3). This shows that $\rho_{\infty} \in K_{z_{\infty}}$. Then with Stroock-Varadhan [26, Lemma 12.1.8 and Theorem 12.1 .10 , p. 289], we can find a probability kernel $\widetilde{\rho}_{z}\left(d z_{1}, d z_{2}\right), z \in \mathbb{R}^{d}$, such that

$$
\begin{equation*}
\text { for } z \in \mathbb{R}^{d}, \widetilde{\rho}_{z}(\cdot) \in K_{z} \tag{3.9}
\end{equation*}
$$

and define the transition probability $\bar{\rho}_{z, z_{0}}\left(d z^{\prime}, d z_{0}^{\prime}\right)$ on $\mathbb{R}^{d} \times \mathbb{R}^{d}$ :

$$
\begin{equation*}
\int g\left(z^{\prime}, z_{0}^{\prime}\right) \bar{\rho}_{z, z_{0}}\left(d z^{\prime}, d z_{0}^{\prime}\right)=\int g\left(z_{1}, z_{2}-z+z_{0}\right) \widetilde{\rho}_{z}\left(d z_{1}, d z_{2}\right), \tag{3.10}
\end{equation*}
$$

for $g$ bounded measurable on $\mathbb{R}^{d} \times \mathbb{R}^{d}$, and $z, z_{0} \in \mathbb{R}^{d}$.
We then define $Q_{n, y}$ as the canonical law of the Markov chain with transition kernel $\bar{\rho}$ and initial distribution concentrated on $(y, y)$. With (3.7), (3.9), it is straightforward to check that (3.5) holds. To prove (3.6), observe that for $k \geq 1$ :

$$
\begin{align*}
& E^{Q_{n, y}}\left[d_{n, \beta}\left(\bar{X}_{k}, \bar{X}_{k}^{0}\right)\right] \leq E^{Q_{n, y}}\left[d_{n, \beta}\left(\bar{X}_{k-1}, \bar{X}_{k-1}^{0}\right)\right]+ \\
& E^{Q_{n, y}}\left[d_{n, \beta}\left(\bar{X}_{k}, \bar{X}_{k}^{0}-\bar{X}_{k-1}^{0}+X_{k-1}\right)\right] \stackrel{(3.9),(3.10)}{=}  \tag{3.11}\\
& E^{Q_{n, y}}\left[d_{n, \beta}\left(\bar{X}_{k-1}, \bar{X}_{k-1}^{0}\right)\right]+E^{Q_{n, y}}\left[D_{n, \beta}\left(\widetilde{R}_{n, h}\left(\bar{X}_{k-1}, \cdot\right), R_{n}^{0}\left(\bar{X}_{k-1}, \cdot\right)\right)\right]
\end{align*}
$$

To bound the rightmost term, note that for $z \in \mathbb{R}^{d}$, when $x \in L_{n} \mathbb{Z}^{d}$ is such that $|z-x| \leq$ $\sqrt{d} L_{n}$, and $f$ has Lipschitz constant at most 1 with respect to $d_{n, \beta}(\cdot, \cdot)$, one finds

$$
\begin{equation*}
\left|\widetilde{R}_{n, h} f(z)-R_{n}^{0} f(z)\right| \stackrel{(3.4)}{=} h(z)\left|\widetilde{S}_{n} f(z)\right|=h(z)\left|\widetilde{S}_{n}(f(\cdot)-f(x))(z)\right| \tag{3.12}
\end{equation*}
$$

and since $\widetilde{R}_{n}(z, \cdot)$ is supported in $\bar{B}\left(z, \widetilde{D}_{n}\right)$ with (1.23), (1.37)

$$
\begin{aligned}
& \leq h(z)\left|\widetilde{S}_{n} F(z)\right|+h(z)\left|\widetilde{S}_{n}\left[\left(1-\chi_{2 \sqrt{d} \widetilde{D}_{n}}(\cdot-x)\right)(f(\cdot)-f(x))\right](z)\right| \\
& \leq h(z)\left|\left(\chi_{n, x} \widetilde{S}_{n} F\right)(z)\right|+h(z) R_{n}^{0}\left[1_{B\left(x, 2 \sqrt{d} \widetilde{D}_{n}\right)^{c}}(\cdot)\left|\frac{-x}{L_{n}}\right|^{\beta}\right](z),
\end{aligned}
$$

with $F(\cdot)=\chi_{2 \sqrt{d} \widetilde{D}_{n}}(\cdot-x)(f(\cdot)-f(x))$. Note that

$$
\begin{equation*}
|F|_{(n)} \leq \kappa_{n}, \tag{3.13}
\end{equation*}
$$

and we now see that the left-hand side of (3.12) is smaller than

$$
h(z)\left(\kappa_{n}\left\|\chi_{n, x} \widetilde{S}_{n}\right\|_{n}+W_{0}\left[X_{\alpha_{n} L_{n}^{2}} \notin B\left(0,2 \sqrt{d} \widetilde{D}_{n}\right)\right]^{\frac{1}{2}} E^{W_{0}}\left[\left|\frac{X_{\alpha_{n} L_{n}^{2}}}{L_{n}}\right|^{2 \beta}\right]^{\frac{1}{2}}\right) \leq \kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}}
$$

With (3.3), we see that we have shown that

$$
\begin{equation*}
\sup _{z \in \mathbb{R}^{d}} D_{n, \beta}\left(\widetilde{R}_{n, h}(z, \cdot), R_{n}^{0}(z, \cdot)\right) \leq \kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}} \tag{3.14}
\end{equation*}
$$

Coming back to (3.11), using induction over $k$, and the fact that $\bar{X}_{0}=\bar{X}_{0}, Q_{n, y-\text { a.s., we }}$ find for $k \geq 0$,

$$
\begin{equation*}
E^{Q_{n, y}}\left[d_{n, \beta}\left(\bar{X}_{k}, \bar{X}_{k}^{0}\right)\right] \leq k\left(\kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}}\right) . \tag{3.15}
\end{equation*}
$$

The application of Chebyshev's inequality now yields for $\gamma>0, k_{0} \geq 1$ :
$\left(\frac{\gamma}{L_{n}}\right)^{\beta} Q_{n, y}\left[\left|\bar{X}_{k}-\bar{X}_{k}^{0}\right| \geq \gamma\right.$, for some $\left.k \leq k_{0}\right] \leq \sum_{k=1}^{k_{0}} k\left(\kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}}\right) \leq k_{0}^{2}\left(\kappa_{n} \Gamma_{n, h}+e^{-\kappa_{n}}\right)$, which proves (3.6).

We can now return to the main object of this section, namely propagating (1.48) "at level $n_{0}+1$ ". The idea is to devise exit strategies from $C_{n_{0}+1}(x)$ for the path, that show that it is costly for the environment to produce $J_{n_{0}+1, x, \cdot, \text {. }}$ variables above level $u_{x}$, for $x$ in a finite collection $\mathcal{A}$. The nature of the exit strategies depends on the level $u_{x}$, and there are four regimes, (only three when $d \geq 4$ ), cf. (3.20). The higher the $u_{x}$, the more the exit strategy relies on the control (1.48) at level $n_{0}$. The lower the $u_{x}$, the more the exit strategy relies on "good-behavior" of the environment around $C_{n_{0}+1}^{\prime}(x)$ at the micro-level $n_{0}-1$, in the sense of (2.2), so that good couplings with Brownian motion resulting from Proposition 3.1 can be employed. Good behavior is precisely expressed by the events $\mathcal{C}_{x}$, cf. (3.24), (3.32), and below (3.33). As one of the first steps, we reduce ourselves to a situation of "only good behavior", cf. (3.36). This involves a certain thinning procedure of $\mathcal{A}$ singling out local high values of $u_{x}$ and showing that bad behavior of the environment at these sites is costly, cf. (3.36). We then have to control the probability that the variables $J_{n_{0}+1, x, r, \text {, }}$ are bigger than $u_{x}$, for $x$ in a thinning of $\mathcal{A}$, in the presence of good-behavior of the environment, cf. Lemma 3.4. This is done with the help of the exit strategies that enable to bound the variables, $J_{n_{0}+1, x, \cdot, \text {, }}$ from above, in terms of $J_{n_{0}, \cdot,,,}$ variables, cf. (3.50), (3.58), (3.71), (3.76), and then use the induction assumption, cf. (3.78). The constant $\zeta$, cf. (1.43), (1.44), is important in the treatment of the lower values of $u_{x}$, cf. (3.85). We then go back from the estimates on the thinned collection with good-behavior of the environment to the general upper bound in (3.86).

Proposition 3.3. One can choose a (large enough) positive constant $c_{2}$ in (1.48), such that for large $L_{0}$ and $n_{0} \geq m_{0}+1$, if (1.49) holds for $n_{0}$ and (1.47), (1.48) hold for $0 \leq n \leq n_{0}$, then (1.48) holds for $n_{0}+1$.
Proof. We consider $\left(u_{x}, A_{x}, \gamma_{x}\right)_{x \in \mathcal{A}}$, with $\mathcal{A}$ a finite subset of $L_{n_{0}+1} \mathbb{Z}^{d}$, an $\left(n_{0}+1\right)$ admissible family, cf. (1.45). From the definition (1.44), we see that

$$
\begin{align*}
J_{n, x, A \cup B, \gamma}=J_{n, x, A, \gamma} \vee J_{n, x, B, \gamma}, & \text { for } n \geq 0, x \in L_{n} \mathbb{Z}^{d}, A, B \subset C_{n}(x), \\
& 1 \leq \gamma \leq 2 d 5^{d-1} \tag{3.16}
\end{align*}
$$

As a result we see that

$$
\begin{align*}
& \mathbb{P}\left[\forall x \in \mathcal{A}, J_{n_{0}+1, x, A_{x}, \gamma_{x}} \geq u_{x}\right] \leq  \tag{3.17}\\
& \left(c \ell_{n_{0}-1}^{d} \ell_{n_{0}}^{d}\right)^{|\mathcal{A}|} \widetilde{\sup } \mathbb{P}\left[\forall x \in \mathcal{A}, J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}\right]
\end{align*}
$$

where $\widetilde{\sup }$ stands for the supremum over families $z_{x} \in L_{n_{0}-1} \mathbb{Z}^{d}, x \in \mathcal{A}$, with $C_{n_{0}-1}\left(z_{x}\right) \subseteq$ $C_{n_{0}+1}(x)$, and $d_{\infty}\left(C_{n_{0}-1}\left(z_{x}\right), C_{n_{0}-1}\left(z_{x^{\prime}}\right)\right) \geq 10 d L_{n_{0}}$, for $x \neq x^{\prime}$, in $\mathcal{A}$.

We will now work on the rightmost term of (3.17). To this end we introduce a thinning $\widetilde{\mathcal{A}}$ of $\mathcal{A}$ as follows. We pick some $x_{1} \in \mathcal{A}$ such that $u_{x_{1}}=\max _{x} u_{x}$, and define $\mathcal{N}_{1}=\{x \in$ $\mathcal{A},\left|x-x_{1}\right|_{\infty} \leq L_{n_{0}+1}$, and $\left.\left(u_{x}+1\right) \log L_{n_{0}}<\left(u_{x_{1}}+1\right)\right\}$, where we recall that $|\cdot|_{\infty}$ denotes the sup-norm on $\mathbb{R}^{d}$. So $\mathcal{N}_{1}$ corresponds to the boxes $C_{n_{0}+1}(x), x \in \mathcal{A}$, adjacent to $C_{n_{0}+1}\left(x_{1}\right)$, with value $\left(u_{x}+1\right)$ smaller than $\left(u_{x_{1}}+1\right) / \log L_{n_{0}}$. We define

$$
\mathcal{A}_{1}=\mathcal{A} \backslash\left(\mathcal{N}_{1} \cup\left\{x_{1}\right\}\right) .
$$

Either $\mathcal{A}_{1}=\emptyset$, in which case the process stops, or $\mathcal{A}_{1} \neq \emptyset$, and we repeat the same procedure to $\mathcal{A}_{1}$, and define $x_{2}, \mathcal{N}_{2}$ as above, and set $\mathcal{A}_{2}=\mathcal{A}_{1} \backslash\left(\mathcal{N}_{2} \cup\left\{x_{1}\right\}\right)$, and so on. After $p$ steps, with $p \leq|\mathcal{A}|$, one has $\mathcal{A}_{p}=\emptyset$, and the process stops. We then write

$$
\begin{equation*}
\tilde{\mathcal{A}}=\left\{x_{1}, \ldots, x_{p}\right\}=\mathcal{A} \backslash \bigcup_{1 \leq i \leq p} \mathcal{N}_{i} \tag{3.18}
\end{equation*}
$$

and observe that

$$
\begin{align*}
& x, x^{\prime} \in \widetilde{\mathcal{A}} \text { and }\left|x-x^{\prime}\right|_{\infty} \leq L_{n_{0}+1} \text { implies }\left(\log L_{n_{0}}\right)^{-1} \leq \frac{u_{x^{\prime}}+1}{u_{x}+1} \leq \log L_{n_{0}}, \text { and } \\
& \sum_{x \in \mathcal{A}}\left(u_{x}+1\right) \leq\left(1+\frac{3^{d}}{\log L_{n_{0}}}\right) \sum_{x \in \widetilde{\mathcal{A}}}\left(u_{x}+1\right) \tag{3.19}
\end{align*}
$$

We introduce the notation $a_{d}=\frac{3}{4}(d-2) a$, and partition $\widetilde{\mathcal{A}}$ into four subsets:

$$
\begin{array}{ll}
\widetilde{\mathcal{A}}_{1}=\left\{x \in \widetilde{\mathcal{A}} ; u_{x} \geq L_{n_{0}}^{a}\right\}, & \widetilde{\mathcal{A}}_{2}=\left\{x \in \widetilde{\mathcal{A}} ; L_{n_{0}}^{a_{d}} \leq u_{x}<L_{n_{0}}^{a}\right\}  \tag{3.20}\\
\widetilde{\mathcal{A}}_{3}=\left\{x \in \widetilde{\mathcal{A}} ; \log L_{n_{0}} \leq u_{x}<L_{n_{0}}^{a_{d} \wedge a}\right\}, & \widetilde{\mathcal{A}}_{4}=\left\{x \in \widetilde{\mathcal{A}} ; 0<u_{x}<\log L_{n_{0}}\right\} .
\end{array}
$$

Note that $\widetilde{\mathcal{A}}_{2}=\emptyset$, whenever $d \geq 4$.
Our aim is to produce an upper bound on quantities of the type $\mathbb{P}[\forall x \in \widetilde{\mathcal{A}}$, $\left.J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{z}} \geq u_{x}\right]$. We will in essence show that $\left\{J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}\right\}$ is unlikely by producing an exit strategy for the process that leads before time $L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}(x)}^{\prime}$ from $y \in C_{n_{0}-1}\left(z_{x}\right) \subseteq C_{n_{0}+1}(x)$ to the box $C_{n_{0}+1, \gamma_{x}}(x)$ with side-length $L_{n_{0}+1} / 5$ that borders $\partial C_{n_{0}+1}(x)$, cf. below (1.41). The nature of this strategy depends on which $\widetilde{\mathcal{A}}_{i}, 1 \leq i \leq 4$, $x$ belongs to. In particular when $x \in \widetilde{\mathcal{A}}_{2}$, or $x \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$, the exit strategy involves certain events describing a "good behavior" of the environment "at level $n_{0}-1$ ". We first specify these events.

We introduce for $x \in \widetilde{\mathcal{A}}_{2}$, (recall this only concerns the case of dimension $d=3$ ), the numbers $\alpha_{x}, \nu_{x}, \nu_{x}^{\prime}$ such that:

$$
\begin{align*}
& u_{x}=L_{n_{0}}^{\alpha_{x}},\left(\text { so that by }(3.20), \alpha_{x} \in\left[\frac{3}{4} a, a\right)\right), \text { and } \\
& 0<\nu_{x} \stackrel{\text { def }}{=} \frac{1}{2}\left(a-\frac{\alpha_{x}}{2}\right)<\nu_{x}^{\prime} \stackrel{\text { def }}{=} \frac{5}{8} \alpha_{x}+\frac{a}{4}<\frac{7}{8} a \tag{3.21}
\end{align*}
$$

We will now define for $x \in \widetilde{\mathcal{A}}_{2}$ the event $\mathcal{C}_{x}$ which in essence specifies the presence in $C_{n_{0}+1}(x)$ of channels of width $L_{n_{0}}^{1+\nu_{x}}$ within distance $\sim L_{n_{0}}^{1+\nu_{x}^{\prime}}$ of any point of $C_{n_{0}+1}(x)$ where the process easily travels. More precisely call a box $B=z+\left[0, L_{n_{0}}^{1+\nu_{x}}\right]^{d}$, $z \in$ $L_{n_{0}}^{1+\nu_{x}} \mathbb{Z}^{d}$, of side-length $L_{n_{0}}^{1+\nu_{x}}, n_{0}$-good for $\omega$, if all $y \in L_{n_{0}-1} \mathbb{Z}^{d}$ within $|\cdot|$-distance $30 \sqrt{d} L_{n_{0}-1}$ of $B$ belong to $\widetilde{\mathcal{B}}_{n_{0}-1}(\omega)$, cf. (2.2). Then set

$$
\begin{align*}
C_{n_{0}+1}^{0}(x)= & \left\{z \in C_{n_{0}+1}(x) ; d\left(z, C_{n_{0}+1}(x)^{c}\right)>L_{n_{0}}^{1+\nu_{x}^{\prime}}\right\}, \\
& \text { and for } e \in \mathbb{Z}^{d},|e|=1  \tag{3.22}\\
C_{n_{0}+1}^{e}(x)= & \left(C_{n_{0}+1}^{0}(x)+2 e L_{n_{0}}^{1+\nu_{x}^{\prime}}\right) \backslash C_{n_{0}+1}(x)
\end{align*}
$$

We now define for $z \in C_{n_{0}+1}^{0}(x), z^{\prime} \in C_{n_{0}+1}^{e}(x)$, ( $e$ as above), and $s>0$ :
$\mathcal{C}_{x}^{z, z^{\prime}, s}=\left\{\omega \in \Omega\right.$; there is a nearest-neighbor path of $n_{0}$-good boxes

$$
B_{1}=z_{1}+\left[0, L_{n_{0}}^{1+\nu_{x}}\right], \ldots, B_{k}=z_{k}+\left[0, L_{n_{0}}^{1+\nu_{x}}\right], k \leq 4 L_{n_{0}}^{a-\nu_{x}}
$$ moving in the $e$-direction after the first $i \in[1, k]$, for which

$$
\left.d_{\infty}\left(z_{i}, C_{n_{0}+1}(x)^{c}\right) \leq \frac{1}{2} L_{n_{0}}^{1+\nu_{x}^{\prime}} \text {, with } d_{\infty}\left(z, B_{1}\right) \vee d_{\infty}\left(z^{\prime}, B_{k}\right) \leq s L_{n_{0}}^{1+\nu_{x}^{\prime}}\right\}
$$

as well as the event

$$
\begin{equation*}
\mathcal{C}_{x}=\bigcap_{z, z^{\prime}} \mathcal{C}_{x}^{z, z^{\prime}, 1}, \text { where } z \text { runs over } C_{n_{0}+1}^{0}(x), z^{\prime} \text { over } \bigcup_{|e|=1} C_{n_{0}+1}^{e}(x) \tag{3.24}
\end{equation*}
$$

(note that requiring $z, z^{\prime}$ to have rational coordinates does not change (3.24), and makes clear that $\mathcal{C}_{x}$ is an event). We now bound $\mathbb{P}\left[\mathcal{C}_{x}^{c}\right]$. We observe that

$$
\begin{equation*}
\mathbb{P}\left[\mathcal{C}_{x}^{c}\right] \leq c L_{n_{0}}^{2 d\left(a-\nu_{x}\right)} \sup _{z, z^{\prime}} \mathbb{P}\left[\left(\mathcal{C}_{x}^{z, z^{\prime}, \frac{1}{2}}\right)^{c}\right], L_{0} \text { large }, \tag{3.25}
\end{equation*}
$$

where $z, z^{\prime}$ respectively run over $\left(L_{n_{0}}^{1+\nu_{x}} \mathbb{Z}^{d}\right) \cap C_{n_{0}+1}^{0}(x)$, and $\bigcup_{|e|=1}\left(L_{n_{0}}^{1+\nu_{x}} \mathbb{Z}^{d}\right) \cap C_{n_{0}+1}^{e}(x)$.
We now set $w=L_{n_{0}}^{-\left(1+\nu_{x}\right)}\left(z^{\prime}-z\right) \in \mathbb{Z}^{d}$, and for convenience assume that $z^{\prime} \in C_{n_{0}+1}^{e_{3}}(x)$ and $w_{i} \geq 0,1 \leq i \leq d(=3)$; the other cases being handled in a similar fashion. For $\theta=\left(\theta_{1}, \theta_{2}\right) \in 2 \mathbb{Z}^{2}$, with $\theta_{1}, \theta_{2} \leq 0$, we define $k_{\theta}=w_{1}+w_{2}+w_{3}+\left|\theta_{1}\right|+\left|\theta_{2}\right|$, and for $0 \leq i<k_{\theta}$,

$$
p_{0}^{\theta}=\left(0, \theta_{1}, \theta_{2}\right), p_{i+1}^{\theta}-p_{i}^{\theta}=\left\{\begin{array}{l}
(1,0,0), 0 \leq i<w_{1}+\left|\theta_{1}\right|  \tag{3.26}\\
(0,1,0), w_{1}+\left|\theta_{1}\right| \leq i<w_{1}+w_{2}+\left|\theta_{1}\right|+\left|\theta_{2}\right| \\
(0,0,1), w_{1}+w_{2}+\left|\theta_{1}\right|+\left|\theta_{2}\right| \leq i<k_{\theta}
\end{array}\right.
$$

as well as

$$
z_{i+1}^{\theta}=z+L_{n_{0}}^{1+\nu_{x}} p_{i}^{\theta}, B_{i+1}^{\theta}=z_{i+1}^{\theta}+\left[0, L_{n_{0}}^{1+\nu_{x}}\right]^{d} .
$$

Note that for $\theta \neq \theta^{\prime}$,

$$
\begin{equation*}
d_{\infty}\left(B_{i}^{\theta}, B_{i^{\prime}}^{\theta^{\prime}}\right) \geq L_{n_{0}}^{1+\nu_{x}}, 1 \leq i \leq k_{\theta}, 1 \leq i^{\prime} \leq k_{\theta^{\prime}} \tag{3.27}
\end{equation*}
$$

and for $\left|\theta_{1}\right|,\left|\theta_{2}\right| \leq \frac{1}{100} L_{n_{0}}^{\nu_{x}^{\prime}-\nu_{x}}, L_{0}$ large,

$$
\left\{\begin{array}{l}
k_{\theta} \leq\left(3 L_{n_{0}+1}+L_{n_{0}}^{1+\nu_{x}^{\prime}}\right) L_{n_{0}}^{-\left(1+\nu_{x}\right)}+\frac{2}{100} L_{n_{0}}^{\nu_{x}^{\prime}-\nu_{x}} \stackrel{(3.21)}{\leq} 4 L_{n_{0}}^{a-\nu_{x}},  \tag{3.28}\\
d_{\infty}\left(z, B_{1}^{\theta}\right) \vee d_{\infty}\left(z^{\prime}, B_{k_{\theta}}^{\theta}\right)<\frac{1}{2} L_{n_{0}}^{1+\nu_{x}^{\prime}}, \text { and for } \\
1 \leq i<k_{\theta}, d_{\infty}\left(z_{i}^{\theta}, C_{n_{0}+1}(x)^{c}\right) \leq \frac{1}{2} L_{n_{0}}^{1+\nu_{x}^{\prime}}, \text { implies } z_{i+1}^{\theta}-z_{i}^{\theta}=L_{n_{0}}^{1+\nu_{x}} e_{3} .
\end{array}\right.
$$



Fig. 2: Candidates for paths of good boxes corresponding to the exit strategy for $\widetilde{\mathcal{A}}_{2}$. Solid lines are made of boxes of side-length $L_{n_{0}}^{1+\nu_{x}}$ and distance between paths of boxes are at least $L_{n_{0}}^{1+\nu_{x}}$

So the paths $B_{i}^{\theta}, 1 \leq i \leq k_{\theta}$, satisfy the requirements set forth in the definition of $\mathcal{C}_{x}^{z, z^{\prime}, \frac{1}{2}}$.
Then for any such given path $B_{i}^{\theta}, 1 \leq i \leq k_{\theta}$,

$$
\begin{equation*}
\mathbb{P}\left[\text { one of the } B_{i}^{\theta} \text { is not } n_{0} \text {-good }\right] \leq\left(\frac{c L_{n_{0}}^{1+\nu_{x}}}{L_{n_{0}-1}}\right)^{d} 4 L_{n_{0}}^{a-\nu_{x}} L_{n_{0}-1}^{-M_{0}} \leq \frac{1}{2} \tag{3.29}
\end{equation*}
$$

when $L_{0}$ is large, cf. (1.14), (1.46), (1.47), (3.21).
Then using independence, cf. (1.7), (2.3), (3.27), we see that

$$
\mathbb{P}\left[\left(\mathcal{C}_{x}^{z, z^{\prime}, \frac{1}{2}}\right)^{c}\right] \leq\left(\frac{1}{2}\right)^{\left(c L_{n_{0}}^{\nu_{x}^{\prime}-\nu_{x}}\right)^{2}}
$$

and using (3.25), we find when $L_{0}$ is large, for $x \in \widetilde{\mathcal{A}}_{2}$,

$$
\begin{align*}
\mathbb{P}\left[\mathcal{C}_{x}^{c}\right] \leq & c L_{n_{0}}^{2 d\left(a-\nu_{x}\right)} \exp \left\{-c L_{n_{0}}^{2\left(\nu_{x}^{\prime}-\nu_{x}\right)}\right\} \stackrel{(3.21)}{\leq} \exp \left\{-c L_{n_{0}}^{\frac{a}{16}+\alpha_{x}}\right\}<  \tag{3.30}\\
& L_{n_{0}-1}^{-6 d 9^{d} M\left(u_{x}+1\right) \log L_{n_{0}}} .
\end{align*}
$$

When $x \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$, (we are back in the case of a general $d \geq 3$ ), the event $\mathcal{C}_{x}$ will in place of (3.24) require that there are "few" boxes $C_{n_{0}-1}(z) \subseteq C_{n_{0}+1}^{\prime}(x)$, cf. (1.41), with $z \notin \widetilde{\mathcal{B}}_{n_{0}-1}(\omega)$. Just as in (3.24), the good behavior of the environment is specified at level $n_{0}-1$. More precisely for $z \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$ and $\omega \in \Omega$, we introduce the compact sets

$$
\begin{equation*}
K_{x, \omega}=\bigcup_{z} \bar{B}\left(z, 30 \sqrt{d} L_{n_{0}-1}\right) \supset \widetilde{K}_{x, \omega}=\bigcup_{z} \bar{B}\left(z, 29 \sqrt{d} L_{n_{0}-1}\right), \tag{3.31}
\end{equation*}
$$

where the unions run over the set of $z \in L_{n_{0}-1} \mathbb{Z}^{d}$, with $d\left(z, C_{n_{0}+1}^{\prime}(x)\right) \leq 30 \sqrt{d} L_{n_{0}-1}$, such that $z \notin \widetilde{\mathcal{B}}_{n_{0}-1}(\omega)$. We then define for $x \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$,

$$
\begin{align*}
& \mathcal{C}_{x}=\left\{\omega \in \Omega ; K_{x, \omega} \text { is contained in the union of } N_{x}\right. \text { open balls } \\
&\text { with radius } \left.4 \widetilde{D}_{n_{0}-1} \text { and centers in } L_{n_{0}-1} \mathbb{Z}^{d}\right\}, \tag{3.32}
\end{align*}
$$

with $N_{x}=\left[12 d 9^{d}(1+a)^{2} \frac{M}{M_{0}}\left(u_{x}+1\right) \log L_{n_{0}}\right]+1$.
For $x \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$, on $\mathcal{C}_{x}^{c}$, arguing by contradiction we can find $N_{x}$ disjoint open balls with radius $\frac{3}{2} \widetilde{D}_{n_{0}-1}$, and centers in $L_{n_{0}-1} \mathbb{Z}^{d} \cap\left(x+L_{n_{0}+1}\left[-\frac{1}{2}, \frac{3}{2}\right]^{d}\right) \cap \widetilde{\mathcal{B}}_{n_{0}-1}^{c}(\omega)$. As a result with (1.7), (1.47), (2.3), we find that for large $L_{0}$, for $x \in \widetilde{\mathcal{A}}_{3} \cup \widetilde{\mathcal{A}}_{4}$ :

$$
\begin{align*}
& \mathbb{P}\left[\mathcal{C}_{x}^{c}\right] \leq\left(c\left(\ell_{n_{0}-1} \ell_{n_{0}}\right)^{d} L_{n_{0}-1}^{-M_{0}}\right)^{N_{x}} \leq\left(c L_{n_{0}-1}^{d a(2+a)-M_{0}}\right)^{N_{x}} \\
& \stackrel{(1.46)}{\leq} L_{n_{0}-1}^{-M_{0} N_{x} / 2} \leq L_{n_{0}+1}^{-6 d 9^{d} M\left(u_{x}+1\right)\left(\log L_{n_{0}}\right)} \tag{3.33}
\end{align*}
$$

For convenience, we set $\mathcal{C}_{x}=\Omega$, for $x \in \widetilde{\mathcal{A}}_{1}$. We now come back to the rightmost term of (3.17), and observe that

$$
\begin{align*}
& \mathbb{P}[\forall x \in \mathcal{A},\left.J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}\right] \leq \\
& 2^{|\mathcal{A}|} \sup _{\mathcal{G} \subseteq \widetilde{\mathcal{A}} \backslash \widetilde{\mathcal{A}}_{1}} \mathbb{P}\left[\text { for } x \in \widetilde{\mathcal{A}}, J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}, \mathcal{C}_{x} \text { for } x \in \mathcal{G}\right.  \tag{3.34}\\
& \mathcal{C}_{x}^{c} \text { for } x \in \widetilde{\left.\mathcal{A} \backslash\left(\widetilde{\mathcal{A}_{1}} \cup \mathcal{G}\right)\right]}
\end{align*}
$$

For $\mathcal{G}$ as above we chose $\mathcal{M}=\mathcal{M}(\mathcal{G})$ a maximal set of non-adjacent $x$ in $\widetilde{\mathcal{A} \backslash(\widetilde{\mathcal{A}}} \widetilde{\mathcal{A}} \cup \mathcal{G})$, (i.e. with mutual $|\cdot|_{\infty}$-distance at least $\left.L_{n_{0}+1}\right)$, and denote by $\overline{\mathcal{M}}$ the set of $x \in \widetilde{\mathcal{A}}$ adjacent to $\mathcal{M}$. Coming back to the definitions of the events $\mathcal{C}_{x}$ in (3.24), and the definition of the variables $J_{n_{0}+1, x, A, \gamma}$, with $A \subseteq C_{n_{0}+1}(x)$, cf. (1.44), we see with the help of (1.7) that when $L_{0}$ is large the collection of events

$$
\begin{equation*}
\mathcal{C}_{x}^{c}, x \in \mathcal{M},\left\{\forall x \in \widetilde{\mathcal{A}} \backslash \overline{\mathcal{M}}, J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}, \mathcal{C}_{x}\right\} \text { are independent. } \tag{3.35}
\end{equation*}
$$

This fact together with (3.30), (3.33), yields that for large $L_{0}$

$$
\mathbb{P}\left[\forall x \in \mathcal{A}, \quad J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}\right] \leq
$$

$$
\begin{equation*}
2^{|\mathcal{A}|} \sup _{\mathcal{G} \subseteq \widetilde{\mathcal{A}} \backslash \widetilde{\mathcal{A}}_{1}}\left\{L _ { n _ { 0 } + 1 } ^ { - 6 d 9 ^ { d } M \sum _ { x \in \mathcal { M } } ( u _ { x } + 1 ) ( \operatorname { l o g } L _ { n _ { 0 } } ) } \mathbb { P } \left[\forall x \in \widetilde{\left.\left.\mathcal{A} \backslash \overline{\mathcal{M}}, J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}, \text { and } \mathcal{C}_{x}\right]\right\} . . . . . .}\right.\right. \tag{3.36}
\end{equation*}
$$

With the help of (3.19) we also have a lower bound on the exponent in the first term in the right-hand side of (3.36), that we will later use in (3.86):

$$
\begin{equation*}
6 d 9^{d} M \sum_{x \in \mathcal{M}}\left(u_{x}+1\right) \log L_{n_{0}} \geq 6 d 3^{d} M \sum_{x \in \overline{\mathcal{M}}}\left(u_{x}+1\right) \tag{3.37}
\end{equation*}
$$

We will now bound the last term in the right-hand side of (3.36):

$$
\begin{equation*}
I \stackrel{\text { def }}{=} \mathbb{P}\left[\forall x \in \mathcal{D}, J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \geq u_{x}, \text { and } \mathcal{C}_{x}\right], \text { with } \mathcal{D}=\widetilde{\mathcal{A}} \backslash \overline{\mathcal{M}} \tag{3.38}
\end{equation*}
$$

Our main control comes in the next
Lemma 3.4. For any positive number $c_{2}$ there are $c^{\prime}, c\left(c_{2}\right)>0$, (see above Theorem 1.1 for the convention concerning constants, and $c_{2}$ is not yet a constant), such that for $L_{0} \geq c\left(c_{2}\right), \prod_{n \geq 0}\left(1-c_{2}\left(\log L_{n}\right)^{-1}\right) \geq \frac{1}{2}$, and with the notation (1.48),

$$
\begin{equation*}
I \leq L_{n_{0}+1}^{-\sum_{x \in \tilde{\mathcal{A}} \mid \overline{\mathcal{M}}} \bar{M}_{n_{0}}\left(1-c^{\prime}\left(\log L_{n_{0}}\right)^{-1}\right)\left(u_{x}+1\right)} L_{n_{0}}^{-\bar{M}_{n_{0}} d 3^{d+1} a\left|\widetilde{\mathcal{A}}_{4} \backslash \overline{\mathcal{M}}\right|} \tag{3.39}
\end{equation*}
$$

Proof. We define for $1 \leq i \leq 4$, in the notations of (3.20), (3.38),

$$
\begin{equation*}
\mathcal{D}_{i}=\mathcal{D} \cap \widetilde{\mathcal{A}}_{i} \tag{3.40}
\end{equation*}
$$

The proof involves the construction of "exit strategies" for the process somewhat in the spirit of what was done in [29]. The nature of these exit strategies from $C_{n_{0}+1}(x)$, leading to $C_{n_{0}+1, \gamma_{x}}(x)$ before time $L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}$, when starting in $C_{n_{0}-1}\left(z_{x}\right)$, depends on which $\mathcal{D}_{i}, 1 \leq i \leq 4, x$ belongs to.

The exit strategy first uses an "exit path" based on a sequence of nearest-neighbor boxes (of size $L_{n_{0}}$ ), $C_{n_{0}}\left(y_{j, x}\right), 0 \leq j \leq j_{x}$, starting at $C_{n_{0}}\left(y_{0, x}\right)$, containing or close to $C_{n_{0}-1}\left(z_{x}\right)$, leading to a final location, the nature of which depends on which $\mathcal{D}_{i}, 1 \leq i \leq 4$, $x$ belongs to.

More precisely we consider a family $\pi_{x}, x \in \mathcal{D}$, of finite sequences $\pi_{x}=\left(y_{j, x}, \gamma_{j, x}\right)_{0 \leq j \leq j_{x}}$ in $L_{n_{0}} \mathbb{Z}^{d} \times\left\{1, \ldots, 2 d 5^{(d-1)}\right\}$, so that writing for simplicity $\left(0 \leq j \leq j_{x}\right)$ :

$$
\begin{align*}
C^{j, x} & =C_{n_{0}}\left(y_{j, x}\right), \Delta^{j, x}=C_{n_{0}, \gamma_{j, x}}\left(y_{j, x}\right)  \tag{3.41}\\
\Delta^{-1, x} & =C_{n_{0}-1}\left(z_{x}\right), \text { for } x \in \mathcal{D} \backslash \mathcal{D}_{4}, \Delta^{-1, x}=C_{n_{0}}\left(y_{0, x}\right), \text { for } x \in \mathcal{D}_{4},
\end{align*}
$$

we have:

$$
\left\{\begin{array}{l}
C_{n_{0}-1}\left(z_{x}\right) \subseteq C^{0, x}, C^{j, x} \subseteq C_{n_{0}+1}(x), 0 \leq j \leq j_{x}, \text { and }  \tag{3.42}\\
\Delta^{j, x} \subseteq C^{j+1, x}, 0 \leq j<j_{x}, \text { when } x \in \mathcal{D} \backslash \mathcal{D}_{4}, \\
\left|y_{0, x}-y_{x}\right|_{\infty} \leq L_{n_{0}}, \text { if } C_{n_{0}}\left(y_{x}\right) \supseteq C_{n_{0}-1}\left(z_{x}\right), \text { when } x \in \mathcal{D}_{4}, \\
\text { (i.e. } \left.C^{0, x} \text { is adjacent to the } n_{0} \text {-box containing } C_{n_{0}-1}\left(z_{x}\right)\right)
\end{array}\right.
$$

and moreover the $\Delta^{j, x}$ are spread apart:

$$
\begin{equation*}
\min \left\{d_{\infty}\left(\Delta^{j, x}, \Delta^{j^{\prime}, x^{\prime}}\right) ;(j, x) \neq\left(j^{\prime}, x^{\prime}\right),-1 \leq j \leq j_{x},-1 \leq j^{\prime} \leq j_{x^{\prime}}\right\} \geq 10 d L_{n_{0}-1} \tag{3.43}
\end{equation*}
$$



Fig. 3: An example where $\mathcal{D}_{1}=\left\{x_{1}, x_{2}\right\}, \mathcal{D}_{2}=\left\{x_{3}\right\}, \mathcal{D}_{3}=\left\{x_{4}\right\}, \mathcal{D}_{4}=\left\{x_{5}, x_{6}\right\}$. In black the boxes $C_{n_{0}-1}\left(z_{x}\right), x \in \mathcal{D}$, and in grey the boxes $C^{j, x}$. The black boxes are at least at mutual $|\cdot|_{\infty}$-distance $10 d L_{n_{0}}$.

We now describe the additional requirements on the $\pi_{x}$ involving which $\mathcal{D}_{i}, 1 \leq i \leq 4, x$ belongs to. So in addition to the above requirements, $\pi_{x}$ are such that:

- when $x \in \mathcal{D}_{4}$ :

$$
\begin{align*}
& j_{x}=0, \text { and in addition to the last line of }(3.42), \gamma_{0, x} \in\left\{1, \ldots, 2 d 5^{(d-1)}\right\}  \tag{3.44}\\
& \text { is arbitrary. }
\end{align*}
$$

- When $x \in \mathcal{D}_{3}$ :

$$
\begin{equation*}
j_{x}=n_{x}+3 d \stackrel{\text { def }}{=}\left[\frac{\left(u_{x}+1\right)}{3}\right]+3 d \tag{3.45}
\end{equation*}
$$

and the nearest-neighbor path $\left(y_{j, x}\right)$ after at most $2 d$ steps is such that $C^{j, x}$ remains inside $C_{n_{0}+1}(x)$ at $|\cdot|_{\infty}$-distance at least $2 L_{n_{0}}$ from $\partial C_{n_{0}+1}(x)$, and moves "along some coordinate direction".

- When $x \in \mathcal{D}_{2}$ :

$$
\begin{equation*}
j_{x} \leq c L_{n_{0}}^{\nu_{x}^{\prime}} \tag{3.46}
\end{equation*}
$$

and the finite sequence $\left(y_{j, x}, \gamma_{y, x}\right)_{j \leq j_{x}}$ is now such that after at-most $2 d$ steps $C^{j, x}$ remains inside $C_{n_{0}+1}(x)$ at $|\cdot|_{\infty}$-distance at least $2 L_{n_{0}}$ from $\partial C_{n_{0}+1}(x)$, and the path ends with $C^{j_{x}, x}, \Delta^{j_{x}, x} \subset C_{n_{0}+1}^{0}(x)$, cf. (3.22).

- When $x \in \mathcal{D}_{1}$;

$$
\begin{equation*}
j_{x} \leq c \ell_{n_{0}} \tag{3.47}
\end{equation*}
$$

after at most $2 d$ steps $C^{j, x}, j<j_{x}-1$, remains at least at $|\cdot|_{\infty}$-distance $2 L_{n_{0}}$ from $\partial C_{n_{0}+1}(x)$, and the path ends with $C^{j_{x}, x}, \Delta^{j_{x}, x}$, so that $\Delta^{j_{x}, x} \subseteq C_{n_{0}+1, \gamma_{x}}(x)$.

We will use the fact that when $L_{0}$ is large we can select $\pi_{x}$, when $x \in \mathcal{D} \backslash\left(\mathcal{D}_{2} \cup \mathcal{D}_{4}\right)$ and then complete it into $\pi_{x}, x \in \mathcal{D}$, so that $\gamma_{j_{x}, x}$ is arbitrary and $y_{j_{x}, x}$ is an arbitrary point of, cf. (3.22), $L_{n_{0}} \mathbb{Z}^{d} \cap C_{n_{0}+1}^{0}(x) \cap B_{\infty}\left(z_{x}, 3 L_{n_{0}}^{1+\nu_{x}^{\prime}}\right)$, when $x \in \mathcal{D}_{2}$, when $x \in \mathcal{D}_{4}, C_{n_{0}}\left(y_{0, x}\right)$ is an arbitrary adjacent box of $C_{n_{0}}\left(y_{x}\right) \supseteq C_{n_{0}-1}\left(z_{x}\right), \gamma_{0, x}$ is arbitrary in $\left\{1, \ldots, 2 d 5^{(d-1)}\right\}$, and $\pi_{x}, x \in \mathcal{D}$ fulfills all the above properties.

We will now derive lower bounds on the exit probabilities of $C_{n_{0}+1}(x)$ before time $L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}$, via $C_{n_{0}+1, \gamma_{x}}(x)$, when starting in $C_{n_{0}-1}\left(z_{x}\right)$, for $x \in \mathcal{D}$. We only need to consider $\omega$ such that $\omega \in \mathcal{C}_{x}$, for $x \in \mathcal{D}$, cf. (3.38). These lower bounds will yield upper bounds on the variables $J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}}, x \in \mathcal{D}$, in terms of $J_{n_{0}, r, r}$, variables to which we will apply the induction assumption (1.48). In what follows $\pi_{x}, x \in \mathcal{D}$, always stand for a family of finite sequences satisfying (3.41) - (3.47). We also introduce the shorthand notation

$$
\begin{equation*}
J_{j, x}=J_{n_{0}, y_{j, x}, \Delta^{j-1}, \gamma_{j, x}}, 0 \leq j \leq j_{x}, x \in \mathcal{D} . \tag{3.48}
\end{equation*}
$$

When $x \in \mathcal{D}_{1}$ : we use the path of boxes $C^{j, x}$ and "boundary boxes" $\Delta^{j, x}, 0 \leq j \leq j_{x}$, to let the path exit. Noting that $c l_{n_{0}} L_{n_{0}}^{2}<L_{n_{0}+1}^{2}$, when $L_{0}$ is large, the strong Markov property implies that for $\omega \in \Omega$ :

$$
\begin{equation*}
\inf _{y \in C_{n_{0}-1}\left(z_{x}\right)} P_{y, \omega}\left[H_{C_{n_{0}+1, \gamma_{x}}(x)} \leq L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}^{\prime}\right] \geq \prod_{0 \leq j \leq j_{x}} c_{1} L_{n_{0}}^{-\zeta J_{j, x}} . \tag{3.49}
\end{equation*}
$$

Using that for large $L_{0}$, cf. (1.15), $L_{n_{0}} \leq 2 L_{n_{0}+1}^{(1+a)^{-1}}$, we now find the desired upper bound:

$$
\begin{equation*}
J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \leq c \ell_{n_{0}}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \sum_{0 \leq j \leq j_{x}} J_{j, x} . \tag{3.50}
\end{equation*}
$$

When $x \in \mathcal{D}_{2}, \omega \in \mathcal{C}_{x}$ : the event $\mathcal{C}_{x}$, cf. (3.24), ensures the presence of many channels made of at most $4 L_{n_{0}}^{a-\nu_{x}} n_{0}$-good boxes of size $L_{n_{0}}^{1+\nu_{x}}$, along which, as we now explain, the diffusion travels well.

Indeed consider $B_{0}$ and $B_{1}=B_{0}+L_{n_{0}}^{1+\nu_{x}} e$, with $|e|=1, e \in \mathbb{Z}^{d}$, two neighboring $n_{0^{-}}$ good boxes. Denote with $U$ the interior of $B_{0} \cup B_{1}$, with $V_{0}$ the concentric sub-cube of $B_{0}$ with half-side length, with $V_{1}=V_{0}+L_{n_{0}}^{1+\nu_{x}} e$, the corresponding sub-cube of $B_{1}$, and with $W_{1}$ the concentric sub-cube of $B_{1}$ with quarter side-length. Denote with $h$ a continuous [ 0,1$]$-valued function, equal to 1 on $U$ and vanishing outside an $L_{n_{0}-1}$-neighborhood of $U$. We can consider the coupling measure $Q_{n_{0}-1, y}$, for $y \in V_{0}$, constructed in Proposition 3.1. Choosing in the notations of Proposition 3.1:

$$
k_{0}=\left[\frac{L_{n_{0}}^{1+\nu_{x}}}{L_{n_{0}-1}}\right]^{2}\left(\leq L_{n_{0}-1}^{2\left(1+\nu_{x}\right)(1+a)-2} \stackrel{(3.21)}{\leq} L_{n_{0}-1}^{4 a+a^{2}}\right), \text { and } \gamma=L_{n_{0}-1}
$$

it follows from standard Brownian estimates and Remark 3.2, that

$$
\begin{equation*}
\inf _{y \in V_{0}} Q_{n_{0}-1, y}\left(\bar{X}_{k_{0}}^{0} \in W_{1}, \text { and } d\left(\bar{X}_{k}, U^{c}\right) \geq L_{n_{0}}, \text { for } 0 \leq k \leq k_{0}\right) \geq c \tag{3.51}
\end{equation*}
$$

By construction, see above (3.22), in the notations of (3.6), we have for large $L_{0}$ :

$$
\begin{equation*}
k_{0}^{2}\left(\kappa_{n_{0}-1} \Gamma_{n_{0}-1, h}+e^{-\kappa_{n_{0}-1}}\right) \leq \kappa_{n_{0}-1} L_{n_{0}-1}^{8 a+4 a^{2}} L_{n_{0}-1}^{-\delta} \stackrel{(1.14),(1.40)}{\leq} L_{n_{0}-1}^{-\delta / 2} \tag{3.52}
\end{equation*}
$$

So in the notations of (1.8), (1.19) we find for large $L_{0}$ :

$$
\begin{align*}
& \inf _{y \in V_{0}} P_{y, \omega}\left[H_{V_{1}}<T_{U} \wedge\left(k_{0} L_{n_{0}-1}^{2}\right)\right] \geq \\
& \inf _{y \in V_{0}} P_{y, \omega}\left[X_{k L_{n_{0}-1}^{2}} \in V_{1}, \text { and for } 0 \leq k<k_{0}, d\left(X_{k L_{n_{0}-1}^{2}}, U^{c}\right) \geq \frac{L_{n_{0}}}{2},\right. \text { and } \\
& \left.T_{n_{0}-1} \circ \theta_{k L_{n_{0}-1}^{2}}>L_{n_{0}-1}^{2}\right] \stackrel{(2.2)}{\geq} \inf _{y \in V_{0}} Q_{n_{0}-1, y}\left(\bar{X}_{k_{0}} \in V_{1}, d\left(\bar{X}_{k}, U^{c}\right) \geq \frac{L_{n_{0}}}{2},\right. \text { for }  \tag{3.53}\\
& \left.0 \leq k \leq k_{0}\right)-k_{0} e^{-\kappa_{n_{0}-1}} \stackrel{(3.6),(3.51),(3.52)}{\geq} c .
\end{align*}
$$

So (3.53) shows in a quantitative way that the diffusion "travels well" from $V_{0}$ to $V_{1}$ without leaving $U$. We now explain how this is used to construct an exit strategy from $C_{n_{0}-1}\left(z_{x}\right)$ to $C_{n_{0}+1, \gamma_{x}}(x)$, before time $L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}$.

We use the path of boxes $C^{j, x}$ with boundary boxes $\Delta^{j, x}, 0 \leq j \leq j_{x}$, to go from $C_{n_{0}-1}\left(z_{x}\right)$ to $\Delta^{j_{x}, x} \subset B_{\infty}\left(z_{x}, 2 L_{n_{0}}^{1+\nu_{x}^{\prime}}\right) \cap C_{n_{0}+1}^{0}(x)$, and $\Delta^{j_{x}, x}$ is chosen to be inside a channel of $n_{0}$-good boxes that exit $C_{n_{0}+1}(x)$ in $C_{n_{0}+1, \gamma_{x}}(x)$. We use repeatedly (3.53) to control how the diffusion travels in the channel. The above strategy brings the path to its goal before time

$$
\left[c L_{n_{0}}^{\nu_{x}^{\prime}} L_{n_{0}}^{2}+\frac{c L_{n_{0}+1}}{L_{n_{0}}^{1+\nu_{x}}}\left(L_{n_{0}}^{1+\nu_{x}}\right)^{2}\right] \wedge T_{C_{n_{0}+1}^{\prime}(x)} \leq L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}
$$

We thus find that, cf. (3.46), (1.44),

$$
\begin{align*}
& \inf _{y \in C_{n_{0}-1(z x)}} P_{y, \omega}\left[H_{C_{n_{0}+1, \gamma_{x}}}(x) \leq L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}\right] \geq \\
& c^{L_{n_{0}+1} / L_{n_{0}}^{1+\nu_{x}}} c_{1}^{c L_{n_{0}^{x}}^{\prime}} \widetilde{\sim} \widetilde{\inf }\left\{L_{n_{0}}^{-\zeta \sum_{j=0}^{j_{x}} J_{j, x}}\right\}, \tag{3.54}
\end{align*}
$$

where $\widetilde{\text { inf }}$ refers to the fact that one takes the infimum over a collection of finite sequences $\pi_{x}$, with all possible end points $y_{j_{x}, x} \in L_{n_{0}} \mathbb{Z}^{d} \cap C_{n_{0}+1}^{0}(x) \cap B_{\infty}\left(z_{x}, 2 L_{n_{0}}^{1+\nu_{x}^{\prime}}\right)$. This is an infimum over a set of cardinality smaller than

$$
\begin{equation*}
c L_{n_{0}}^{d \nu_{x}^{\prime}} \stackrel{(3.21)}{\leq} L_{n_{0}}^{d a}, L_{0} \text { large } . \tag{3.55}
\end{equation*}
$$

Further from our choice in (3.21), we see that

$$
\begin{align*}
\alpha_{x}-\left(a-\nu_{x}\right) & =\frac{3}{4} \alpha_{x}-\frac{a}{2} \stackrel{(3.21)}{\geq} \frac{9}{16} a-\frac{a}{2}=\frac{a}{16}  \tag{3.56}\\
\alpha_{x}-\nu_{x}^{\prime} & =\frac{3}{8} \alpha_{x}-\frac{a}{4} \stackrel{(3.21)}{\geq} \frac{9}{32} a-\frac{a}{4}=\frac{a}{32} . \tag{3.57}
\end{align*}
$$

As a result of (3.54), analogously to (3.50), we find that for $x \in \mathcal{D}_{2}, \omega \in \mathcal{C}_{x}$,

$$
\begin{align*}
& J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \leq  \tag{3.58}\\
& c\left(L_{n_{0}}^{\left(a-\nu_{x}\right)}+L_{n_{0}}^{\nu_{x}^{\prime}}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \widetilde{\inf }\left\{\sum_{0 \leq j \leq j_{x}} J_{j, x}\right\} \leq\right. \\
& c L_{n_{0}}^{\alpha_{x}-\frac{a}{32}}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \widetilde{\inf }\left\{\sum_{0 \leq j \leq x} J_{j, x}\right\},
\end{align*}
$$

and $\widetilde{\text { inf }}$ has the same meaning as in (3.54) and involves the infimum over a set of cardinality bounded by (3.55).

We now turn to the discussion of $x \in \mathcal{D}_{3}$ and $x \in \mathcal{D}_{4}$, beginning with some considerations on $\mathcal{C}_{x}$, when $x \in \mathcal{D}_{3} \cup \mathcal{D}_{4}$. We thus consider an $x \in \mathcal{D}_{3} \cup \mathcal{D}_{4}, \omega \in \mathcal{C}_{x}$, and $y \in C_{n_{0}}\left(y_{0}\right)$ with $d_{\infty}\left(C_{n_{0}}\left(y_{0}\right), C_{n_{0}+1}(x)\right) \leq L_{n_{0}}$, such that in the notations of (3.31):

$$
\begin{equation*}
d\left(y, K_{x, \omega}\right) \stackrel{\text { def }}{=} r>0 \tag{3.59}
\end{equation*}
$$

For $m \geq 1$, we define

$$
\begin{align*}
D_{m}= & \widehat{y}_{0}+2^{m}\left(\left[-\frac{L_{n_{0}}}{2}, \frac{L_{n_{0}}}{2}\right]^{d} \backslash\left(-\frac{L_{n_{0}}}{4}, \frac{L_{n_{0}}}{4}\right)^{d}\right), \text { with }  \tag{3.60}\\
& \widehat{y}_{0} \text { the center of } C_{n_{0}}\left(y_{0}\right), \\
K_{m}= & K_{x, \omega} \cap D_{m}, K_{0}=K_{x, \omega} \cap C_{n_{0}}\left(y_{0}\right) . \tag{3.61}
\end{align*}
$$

Keeping in mind $L_{n_{0}+1}$ as a unit scale, we consider for $m \geq 0$, the Newtonian capacity of $L_{n_{0}+1}^{-1} K_{m}$ :

$$
\begin{equation*}
\operatorname{cap}_{m}=\operatorname{cap}\left(L_{n_{0}+1}^{-1} K_{m}\right) \stackrel{(3.32)}{\leq} \kappa_{n_{0}-1} \frac{N_{x}}{\left(\ell_{n_{0}-1} \ell_{n_{0}}\right)^{d-2}} \tag{3.62}
\end{equation*}
$$

We now consider an arbitrary continuous, $[0,1]$-valued, function $h$ such that:

$$
\begin{align*}
& h=1 \text { on } C_{n_{0}+1}^{\prime}(x) \backslash \widetilde{K}_{x, \omega} \stackrel{(3.31)}{\supseteq} C_{n_{0}+1}^{\prime}(x) \backslash K_{x, \omega}, \text { and }  \tag{3.63}\\
& h \chi_{n_{0}-1, z} \equiv 0, \text { for all } z \in L_{n_{0}-1} \mathbb{Z}^{d} \cap \widetilde{\mathcal{B}}_{n_{0}-1}^{c}(\omega) .
\end{align*}
$$

We can now consider the coupling measure $Q_{n_{0}-1, y}$ from Proposition 3.1. Keeping in mind that under this measure $\bar{X}_{k}^{0}, k \geq 0$, is a Brownian motion starting from $y$ sampled at times $\alpha_{n_{0}-1} k L_{n_{0}-1}^{2}$, we see from an analogous calculation as for the classical Wiener test, cf. [28, p. 72-74], that

$$
\begin{align*}
& Q_{n_{0}-1, y}\left[\bar{X}_{k}^{0} \in K_{x, \omega}, \text { for some } k \geq 0\right] \leq \\
& c\left(\sum_{m \geq 2} \operatorname{cap}_{m}\left(2^{m} \ell_{n_{0}}^{-1}\right)^{-(d-2)}+\sum_{m=0,1} \operatorname{cap}_{m}\left(\frac{r}{L_{n_{0}+1}}\right)^{-(d-2)}\right) \stackrel{(3.62)}{\leq}  \tag{3.64}\\
& \kappa_{n_{0}-1} N_{x}\left(\ell_{n_{0}-1}^{-(d-2)}+\left(\frac{r}{L_{n_{0}-1}}\right)^{-(d-2)}\right)
\end{align*}
$$

where we recall the notation (3.59).
We now proceed in a similar fashion as in (3.53), with the help of Proposition 3.1, choosing in (3.6) $\gamma=L_{n_{0}-1}$, and

$$
\begin{equation*}
k_{0}=\left[\frac{1}{10}\left(\frac{L_{n_{0}+1}}{L_{n_{0}-1}}\right)^{2}\right] \leq L_{n_{0}-1}^{4 a+2 a^{2}} . \tag{3.65}
\end{equation*}
$$

We find that for large $L_{0}$ :

$$
\begin{align*}
& P_{y, \omega}\left[H_{C_{n_{0}+1}, \gamma_{x}(x)}<\left(\frac{1}{5} L_{n_{0}+1}^{2}\right) \wedge T_{C_{n_{0}+1}^{\prime}(x)}^{\prime}\right] \geq P_{y, \omega}\left[X_{k_{0} L_{n_{0}-1}^{2}} \in C_{n_{0}+1, \gamma_{x}}(x),\right. \\
& d\left(X_{k L_{n_{0}-1}}^{2}, C_{n_{0}+1}^{\prime}(x)^{c}\right) \geq \frac{L_{n_{0}}}{2}, d\left(X_{k L_{n_{0}-1}}^{2}, \widetilde{\mathcal{B}}_{n_{0}-1}^{c}(\omega) \cap L_{n_{0}-1} \mathbb{Z}^{d}\right) \geq 29 \sqrt{d} L_{n_{0}-1}, \\
& \text { for } \left.0 \leq k \leq k_{0}, T_{n_{0}-1} \circ \theta_{k L_{n_{0}-1}^{2}}^{2}>L_{n_{0}-1}^{2}, \text { for } 0 \leq k<k_{0}\right] \geq  \tag{3.66}\\
& Q_{n_{0}-1, y}\left[\bar{X}_{k_{0}} \in C_{n_{0}+1, \gamma_{x}}(x), d\left(\bar{X}_{k}, C_{n_{0}+1}^{\prime}(x)^{c}\right) \geq \frac{L_{n_{0}}}{2},\right. \\
& \left.d\left(\bar{X}_{k}, \widetilde{\mathcal{B}}_{n_{0}-1}^{c}(\omega) \cap L_{n_{0}-1} \mathbb{Z}^{d}\right) \geq 29 \sqrt{d} L_{n_{0}-1}, \text { for } 0 \leq k \leq k_{0}\right]-k_{0} e^{-\kappa_{n_{0}-1}},
\end{align*}
$$

where we used that $h=1$ on $C_{n_{0}+1}^{\prime}(x) \backslash \widetilde{K}_{x, \omega}$, cf. (3.63), as well as the localization part of (2.2). Then with (3.6), denoting with $\widetilde{C}_{n_{0}+1, \gamma_{x}}(x)$ the concentric box to $C_{n_{0}+1, \gamma_{x}}(x)$, with half-size, we find

$$
\begin{aligned}
& \geq Q_{n_{0}-1, u}\left[\bar{X}_{k_{0}}^{0} \in \widetilde{C}_{n_{0}+1, \gamma_{x}}(x), d\left(\bar{X}_{k}^{0}, C_{n_{0}+1}^{\prime}(x)^{c}\right) \geq L_{n_{0}} \text { for } 0 \leq k \leq k_{0},\right. \\
& \left.\quad \bar{X}_{k}^{0} \notin K_{x, \omega}, \text { for } 0 \leq k \leq k_{0}\right]-k_{0} e^{-\kappa_{n_{0}-1}}-k_{0}^{2}\left(\kappa_{n_{0}-1} L_{n_{0}-1}^{-\delta}+e^{-\kappa_{n_{0}-1}}\right),
\end{aligned}
$$

where we have used that $h \chi_{n_{0}-1, z} \equiv 0$, for $z \in L_{n_{0}-1} \mathbb{Z}^{d} \backslash \widetilde{\mathcal{B}}_{n_{0}-1}(\omega)$, as well as (2.2) in estimating $\Gamma_{n_{0}-1, h}$ of (3.6).

Combining this with (3.64), (3.65), and the inequality

$$
Q_{n_{0}-1, y}\left[\bar{X}_{k_{0}}^{0} \in \widetilde{C}_{n_{0}+1, \gamma_{x}}(x), d\left(\bar{X}_{k}^{0}, C_{n_{0}+1}^{\prime}(x)^{c}\right) \geq L_{n_{0}}, \text { for } 0 \leq k \leq k_{0}\right] \geq 4 c_{1}
$$

that follows from the definition of $c_{1}$ below (1.44), and (1.49), we conclude with (1.14), (1.40) that

$$
\begin{align*}
& P_{y, \omega}\left[H_{C_{n_{0}+1}, \gamma_{x}(x)}<\left(\frac{1}{5} L_{n_{0}+1}^{2}\right) \wedge T_{C_{n_{0}+1}^{\prime}(x)}\right] \geq \\
& 4 c_{1}-\kappa_{n_{0}-1} N_{x}\left(\ell_{n_{0}-1}^{-(d-2)}+\left(\frac{r}{L_{n_{0}-1}}\right)^{-(d-2)}\right) . \tag{3.67}
\end{align*}
$$

This will be a crucial estimate to control exit strategies of the path starting in $C_{n_{0}-1}\left(z_{x}\right)$ and landing in $C_{n_{0}+1, \gamma_{x}}(x)$ before time $L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}$, when $x$ belongs to $\mathcal{D}_{3} \cup \mathcal{D}_{4}$.

When $x \in \mathcal{D}_{3}, \omega \in \mathcal{C}_{x}$ : we describe the exit strategy. First consider the boxes $C^{j, x}$, with boundary boxes $\Delta^{j, x}, 0 \leq j \leq j_{x} \stackrel{(3.45)}{=} n_{x}+3 d$. Consider a path of the diffusion starting in $C_{n_{0}-1}\left(z_{x}\right)$ successively entering the $\Delta^{j, x} \subset C^{j+1, x}$ before time $L_{n_{0}}^{2} \wedge T_{C_{n_{0}}^{\prime}\left(y_{j} ; x\right)}, 0 \leq j \leq j_{x}$. From the time it enters $C^{2 d, x}$ until it enters $\Delta^{j_{x}, x}$, the path remains in $C_{n_{0}+1}(x)$, and has diameter at least $n_{x} L_{n_{0}}$.

If $\theta>0$ is such that the above mentioned portion of the path remains in the open set

$$
\begin{equation*}
U_{\theta}=\left\{y \in \mathbb{R}^{d}, d\left(y, K_{x, \omega}\right)<\theta\right\}, \tag{3.68}
\end{equation*}
$$

in view of (3.32), the fact that $\omega \in \mathcal{C}_{x}$ then implies

$$
n_{x} L_{n_{0}}<2 N_{x}\left(4 \widetilde{D}_{n_{0}-1}+\theta\right)
$$

Choosing

$$
\begin{equation*}
r=\frac{n_{x} L_{n_{0}}}{2 N_{x}}-4 \widetilde{D}_{n_{0}-1}>0, \text { when } L_{0} \text { is large, cf. (3.32), (3.45), } \tag{3.69}
\end{equation*}
$$

the path enters $C_{n_{0}+1}(x) \cap U_{r}^{c}$ before time $\left(j_{x}+1\right) L_{n_{0}}^{2} \wedge T_{C_{n_{0}+1}(x)}^{\prime} \leq \frac{1}{4} L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}$. Letting this entrance point in $C_{n_{0}+1}(x) \cap U_{r}^{c}$ play the role of $y$ in (3.67), we can use the strong Markov property and find that for large $L_{0}$ :

$$
\begin{equation*}
\inf _{w \in C_{n_{0}-1}\left(z_{x}\right)} P_{w, \omega}\left[H_{C_{n_{0}+1, \gamma_{x}}(x)} \leq L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}\right] \geq c_{1}^{j_{x}+1} L_{n_{0}}^{\left(-\zeta \sum_{0 \leq j \leq j x} J_{j, x}\right)} 2 c_{1} \tag{3.70}
\end{equation*}
$$

where we used that thanks to (1.14), (3.20), (3.32), (3.45), (3.69), the last term of (3.67) is arbitrarily small, when $L_{0}$ is large. As a result we thus see that when $L_{0}$ is large:

$$
\begin{equation*}
J_{n_{0}+1, x, C_{n_{0}-1}\left(z_{x}\right), \gamma_{x}} \leq c n_{x}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \sum_{0 \leq j \leq j_{x}} J_{j, x} . \tag{3.71}
\end{equation*}
$$

When $x \in \mathcal{D}_{4}, \omega \in \mathcal{C}_{x}$ : we denote with $\widetilde{C}_{n_{0}, x}$ the union, (we recall that $C_{n_{0}}\left(y_{x}\right) \supseteq$ $C_{n_{0}-1}\left(z_{x}\right)$, cf. (3.42)):

$$
\begin{equation*}
\widetilde{C}_{n_{0}, x}=\bigcup_{\left|\bar{y}-y_{x}\right|_{\infty} \leq L_{n_{0}}} \overline{C_{n_{0}}^{\prime}}(\bar{y}) . \tag{3.72}
\end{equation*}
$$

A path of the diffusion inside $\widetilde{C}_{n_{0}, x}$ starting in $C_{n_{0}-1}\left(z_{x}\right)$ with diameter at least $\frac{1}{2} L_{n_{0}}$ before time $\left(\frac{1}{2} L_{n_{0}+1}^{2}\right) \wedge T_{C_{n_{0}+1}^{\prime}(x)}$, by the same argument as below (3.68), enters before that time the set $\widetilde{C}_{n_{0}, x} \cap U_{r}^{c}$, with

$$
\begin{equation*}
r=\frac{L_{n_{0}}}{4 N_{x}}-4 \widetilde{D}_{n_{0}-1}>0, \text { when } L_{0} \text { is large } . \tag{3.73}
\end{equation*}
$$

If this entrance point in $\widetilde{C}_{n_{0}, x} \cap U_{r}^{c}$ plays the role of $y$, (3.67) provides a lower bound on the probability that the path then reaches $C_{n_{0}+1, \gamma_{x}}(x)$ before $\left(\frac{1}{5} L_{n_{0}+1}^{2}\right) \wedge T_{C_{n_{0}+1}^{\prime}(x)}$.

Note that when starting at $u$ in $C_{n_{0}}(\bar{y})$, with $\left|\bar{y}-y_{x}\right|_{\infty} \leq L_{n_{0}}$ :

$$
\begin{align*}
& P_{u, \omega}\left[X_{L_{n_{0}}^{2} \wedge T_{C_{n_{0}}^{\prime}(\bar{y})}^{*}} \geq \frac{1}{2} L_{n_{0}}\right] \geq c_{1} L_{n_{0}}^{-\zeta J_{x}}, \text { where }  \tag{3.74}\\
& J_{x}=\inf \left\{J_{n_{0}, y^{\prime}, C_{n_{0}}\left(y^{\prime}\right), \gamma^{\prime}} ;\left|y^{\prime}-y_{x}\right|_{\infty} \leq L_{n_{0}}, \gamma^{\prime} \in\left\{1, \ldots, 2 d 5^{(d-1)}\right\}\right\} . \tag{3.75}
\end{align*}
$$

With the strong Markov property, we thus see that for large $L_{0}$,

$$
\begin{align*}
& \inf _{y \in C_{n_{0}-1}\left(z_{x}\right)} P_{y, \omega}\left[H_{C_{n_{0}+1, \gamma x}(x)} \leq L_{n_{0}+1}^{2} \wedge T_{C_{n_{0}+1}^{\prime}(x)}\right] \geq \\
& 2 c_{1}\left(1-\left(1-c_{1} L_{n_{0}}^{-\zeta J_{x}}\right)^{\left.\ell \ell_{n_{0}}^{2} / 3\right]}\right) \geq 2 c_{1}\left(1-\left(1-c_{1} L_{n_{0}}^{-\zeta J_{x}}\right)^{\ell_{n_{0}}^{2} / 4}\right)=  \tag{3.76}\\
& \inf _{y^{\prime}, \gamma^{\prime}} 2 c_{1}\left(1-\left(1-c_{1} L_{n_{0}}^{\left.\left.-\zeta J_{n_{0}, y^{\prime}, C_{n_{0}}\left(y^{\prime}\right), \gamma^{\prime}}\right)^{\ell_{n_{0}}^{2} / 4}\right),}\right.\right.
\end{align*}
$$

where the infimum is over the same set as in (3.75).
We will now employ the bounds (3.50), (3.58), (3.71), (3.76) to bound $I$ in (3.38) and prove the claim (3.39). To keep track of the infima that enter (3.58), (3.76), we introduce a set $\Pi$ of $\pi=\left(\pi_{x}\right)_{x \in \mathcal{D}}$, such that for any $x_{0} \in \mathcal{D}_{2}, \pi \in \Pi$, the set of $\pi^{\prime} \in \Pi$ that coincide with $\pi$ for $x \neq x_{0}$ is such that all points of $L_{n_{0}} \mathbb{Z}^{d} \cap C_{n_{0}+1}^{0}(x) \cap B_{\infty}\left(x, 3 L_{n_{0}}^{1+\nu_{x}^{\prime}}\right)$ and all $\gamma$ in $\left\{1, \ldots, 2 d 5^{(d-1)}\right\}$ occur as $y_{j_{x_{0}, x_{0}}}$ and $\gamma_{j_{x_{0}, x},}$, and similarly for any $x_{0} \in \mathcal{D}_{4}, \pi \in \Pi$, the set of $\pi^{\prime} \in \widetilde{\Pi}$ that coincide with $\pi$ for $x \neq x_{0}$ is such that all $y^{\prime} \in L_{n_{0}} \mathbb{Z}^{d}$ with $C_{n_{0}}\left(y^{\prime}\right) \subset \widetilde{C}_{n_{0}, x}$ and $\gamma^{\prime} \in\left\{1, \ldots, 2 d 5^{(d-1)}\right\}$ occur as $y_{0, x}$ and $\gamma_{0, x}$. With (3.55) we see that when $L_{0}$ is large we can choose such a $\Pi$ with cardinality

$$
\begin{equation*}
|\Pi| \leq L_{n_{0}}^{d a\left|\mathcal{D}_{2}\right|} c^{\left|\mathcal{D}_{4}\right|} . \tag{3.77}
\end{equation*}
$$

Note that for any $\pi \in \Pi$, the sets $\Delta^{j, x},-1 \leq j<j_{x}, x \in \mathcal{D}$, lie at mutual $|\cdot|_{\infty}$-distance at least $10 d L_{n_{0}-1}$, cf. (3.43), so that in view of (1.48), for any choice of $v_{j, x} \geq 0$, where $(j, x) \in \mathcal{J} \xlongequal{\text { def }}\left\{\left(j^{\prime}, x^{\prime}\right) ; x \in \mathcal{D}, 0 \leq j^{\prime} \leq j_{x}\right\}$,

$$
\begin{equation*}
\mathbb{P}\left[\text { for all }(j, x) \in \mathcal{J}, J_{j, x} \geq v_{j, x}\right] \leq \prod_{(j, x) \in \mathcal{J}} P\left[Z \geq v_{j, x}\right] \tag{3.78}
\end{equation*}
$$

where $Z=Z_{1}$, and $Z_{k}, k \geq 1$, is an i.i.d. family of non-negative random variables defined in some auxiliary probability space such that

$$
\begin{equation*}
P[Z>v]=L_{n_{0}}^{-\bar{M}_{n_{0}}(1+v)} \text { for } v>0 \tag{3.79}
\end{equation*}
$$

(so $P[Z=0]=1-L_{n_{0}}^{-\bar{M}_{n_{0}}}$, and we assume from now on that $L_{0} \geq$ const $\left(c_{2}\right)$ so that $\left.\prod_{n \geq 0}\left(1-c_{2}\left(\log L_{n}\right)^{-1}\right) \geq \frac{1}{2}\right)$. Let us mention that (3.78) can be rephrased in terms of upper orthant order, see Shaked-Shanthikumar [24, p. 140]. We denote with $\Sigma_{k}, k \geq 0$, the partial sums

$$
\begin{equation*}
\Sigma_{0}=0, \Sigma_{k}=Z_{1}+\cdots+Z_{k}, \text { for } k \geq 1 \tag{3.80}
\end{equation*}
$$

Note that for $0 \leq \lambda<\bar{M}_{n_{0}} \log L_{n_{0}}$, one has

$$
\begin{equation*}
E\left[e^{\lambda Z}\right]=1+\int_{0}^{\infty} \lambda e^{\lambda v} L_{n_{0}}^{-\bar{M}_{n_{0}}(v+1)} d v=1+\frac{\lambda}{\left(\bar{M}_{n_{0}} \log L_{n_{0}}-\lambda\right)} L_{n_{0}}^{-\bar{M}_{n_{0}}} \tag{3.81}
\end{equation*}
$$

Analogously for an arbitrary collection $v_{x} \geq 0, x \in \mathcal{D}$, and $\lambda_{x} \in\left[0, \bar{M}_{n_{0}} \log L_{n_{0}}\right), x \in$ $\mathcal{D} \backslash \mathcal{D}_{4}$, it follows from (3.78), see also [24, Theorem 5.G.1, p. 141], that:

$$
\begin{align*}
& \mathbb{P}\left[\sum_{0 \leq j \leq j_{x}} J_{j, x} \geq v_{x}, \text { for } x \in \mathcal{D}\right] \leq \\
& \exp \left\{-\sum_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} \lambda_{x} v_{x}\right\} E\left[\exp \left\{\sum_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} \lambda_{x} \sum_{0 \leq j \leq j_{x}} J_{j, x}\right\}, \text { for } x \in \mathcal{D}_{4}, J_{0, x} \geq v_{x}\right] \stackrel{(3.78)}{\leq} \\
& \exp \left\{-\sum_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} \lambda_{x} v_{x}\right\} \prod_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} E\left[e^{\left.\lambda_{x} \Sigma_{\left(j_{x}+1\right)}\right]} \prod_{x \in \mathcal{D}_{4}} P\left[Z \geq v_{x}\right] \stackrel{(3.81)}{=}\right.  \tag{3.82}\\
& \exp \left\{-\sum_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} \lambda_{x} v_{x}\right\} \prod_{x \in \mathcal{D} \backslash \mathcal{D}_{4}}\left(1+\frac{\lambda_{x} L_{n_{0}}^{-\bar{M}_{n_{0}}}}{\bar{M}_{n_{0}} \log L_{n_{0}}-\lambda_{x}}\right)^{j_{x}+1} \prod_{x \in \mathcal{D}_{4}} P\left[Z \geq v_{x}\right] .
\end{align*}
$$

We will now use (3.82) to bound $I$ in (3.38). Indeed for large $L_{0}$, with (3.50), (3.58), (3.71), (3.76) we have

$$
\begin{align*}
I \leq \mathbb{P}[ & \bigcup_{\pi \in \Pi}\left\{c \ell_{n_{0}}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \sum_{0 \leq j \leq j_{x}} J_{j, x} \geq u_{x}, \text { for } x \in \mathcal{D}_{1},\right. \\
& c L_{n_{0}}^{\alpha_{x}-\frac{a}{32}}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \sum_{0 \leq j \leq j_{x}} J_{j, x} \geq u_{x}, \text { for } x \in \mathcal{D}_{2},  \tag{3.83}\\
& c n_{x}\left(\log L_{n_{0}+1}\right)^{-1}+(1+a)^{-1} \sum_{0 \leq j \leq j_{x}} J_{j, x} \geq u_{x}, \text { for } x \in \mathcal{D}_{3}, \\
& \left.\left.1-\frac{1}{2} L_{n_{0}+1}^{-\zeta u_{x}} \leq\left(1-c_{1} L_{n_{0}}^{-\zeta J_{0, x}}\right)^{\ell_{n_{0}}^{2} / 4}, \text { for } x \in \mathcal{D}_{4}\right\}\right] .
\end{align*}
$$

From (3.45) - (3.47), $j_{x} \leq c \ell_{n_{0}}$, for $x \in \mathcal{D} \backslash \mathcal{D}_{4}$, so using (3.77) and (3.82) with $\lambda_{x}=\lambda_{*} \stackrel{\text { def }}{=}$ $\bar{M}_{n_{0}} \log L_{n_{0}}-1$, for all $x \in \mathcal{D} \backslash \mathcal{D}_{4}$, and $(1+u) \leq e^{u}$, we find

$$
\begin{align*}
I \leq & L_{n_{0}}^{d a\left|\mathcal{D}_{2}\right|} c^{\left|\mathcal{D}_{4}\right|} \exp \left\{-\sum_{x \in \mathcal{D} \backslash \mathcal{D}_{4}} \bar{M}_{n_{0}}\left(\log L_{n_{0}+1}\right)\left(u_{x}+1\right)\left(1-c\left(\log L_{n_{0}}\right)^{-1}\right)\right\}  \tag{3.84}\\
& \prod_{x \in \mathcal{D}_{4}} P\left[1-\frac{1}{2} L_{n_{0}+1}^{-\zeta u_{x}} \leq\left(1-c_{1} L_{n_{0}}^{-\zeta Z}\right)^{\ell_{n_{0}}^{2} / 4}\right] .
\end{align*}
$$

Note that with $L_{0}$ large and $\Pi_{n \geq 0}\left(1-c_{2}\left(\log L_{n}\right)^{-1}\right) \geq \frac{1}{2}$, cf. below (3.79), each individual term of the last product is smaller than

$$
\begin{align*}
& P\left[c \ell_{n_{0}}^{-2} L_{n_{0}+1}^{-\zeta u_{x}} \geq L_{n_{0}}^{-\zeta Z}\right] \leq P\left[\zeta Z \geq \zeta u_{x}(1+a)+2 a-c\left(\log L_{n_{0}}\right)^{-1}\right] \\
& \quad \begin{array}{l}
(3.79) \\
\leq \exp \left\{-\left(\log L_{n_{0}}\right) \bar{M}_{n_{0}}\left(u_{x}(1+a)+\frac{2}{\zeta} a-c\left(\log L_{n_{0}}\right)^{-1}+1\right)\right\} \\
\quad \leq \exp \left\{-\left(\log L_{n_{0}}\right) \bar{M}_{n_{0}}\left[(1+a)\left(u_{x}+1\right)+\frac{3}{4}\left(\frac{2}{\zeta}-1\right) a\right]\right\} .
\end{array} . \tag{3.85}
\end{align*}
$$

Coming back to (3.84), we obtain

$$
I \leq L_{n_{0}+1}^{\left[-\sum_{x \in \mathcal{D}} \bar{M}_{n_{0}}\left(1-c\left(\log L_{n_{0}}\right)^{-1}\right)\left(u_{x}+1\right)\right]} L_{n_{0}}^{-\left|\mathcal{D}_{4}\right|\left(\frac{1}{\zeta}-\frac{1}{2}\right) a \bar{M}_{n_{0}}},
$$

and in view of (1.43), this proves (3.39).
We can now conclude the proof of Proposition 3.3. Coming back to (3.17), (3.36), (3.37), (3.39), we observe that when $L_{0}$ is large,

$$
\left(c \ell_{n_{0}-1}^{d} \ell_{n_{0}}^{d}\right)^{|\mathcal{A}|} \leq L_{n_{0}}^{3 d a|\mathcal{A}|} \stackrel{(3.18)}{\leq} L_{n_{0}}^{d 3^{d+1} a|\widetilde{\mathcal{A}}|}
$$

and hence

$$
\begin{align*}
& L_{n_{0}}^{d 3^{d+1} a\left|\widetilde{\mathcal{A}} \backslash \widetilde{\mathcal{A}}_{4}\right|} L_{n_{0}+1}^{-\bar{M}_{n_{0}}\left(1-c^{\prime} / \log L_{n_{0}}\right) \sum_{x \in \tilde{\mathcal{A}}}\left(u_{x}+1\right)} \leq \tag{3.86}
\end{align*}
$$

where $L_{0} \geq \operatorname{const}\left(c_{2}\right)$, so that $\prod_{n \geq 0}\left(1-c_{2}\left(\log L_{n}\right)^{-1}\right) \geq \frac{1}{2}$ and in particular $\bar{M}_{n_{0}} \geq 1$. We then see that if $c_{2}$ is chosen to be constant bigger than the constant $c^{\prime \prime \prime}$ that appears in the last member of (3.86), then (1.48) holds for $n=n_{0}+1$. This proves Proposition 3.3.

## 4 Surgery and contraction of Hölder-norms

We continue the proof of Theorem 1.1. The aim is now to propagate the part of (1.47) that concerns Hölder-norms at level $n_{0}+1$, cf. (1.39). The part of (1.47) that concerns localization controls has been taken care of in Proposition 2.2. The control of Höldernorms will be carried out in the present and in the next section. Here we first perform "surgery" and remove "Hölder-norm defects" at level $n_{0}-m_{0}-1$ that occur in the large box $5 \mathcal{T}_{n_{0}+1}$, see (2.1). We show that with overwhelming $\mathbb{P}$-probability the kernel $R_{n}$ of the diffusion in the modified environment, when starting in $\mathcal{T}_{n_{0}+1}$, gets closer and closer in $\|\cdot\|_{n}$-norm to $R_{n}^{0}$ as $n$ goes from $n_{0}-m_{0}-1$ to $n_{0}$, cf. Proposition 4.11. The crucial step comes in Proposition 4.1, where Hölder-norm estimates are derived on what is in essence the linearization of the evolution after surgery at level $n+1$, when expressed in terms of the one at level $n$, as $n$ varies from $n_{0}-m_{0}-1$ to $n_{0}$.

As a shorthand notation, we write, cf. Theorem 1.1,

$$
\begin{equation*}
n_{0}^{\prime}=n_{0}-m_{0}-1 \geq 0 \tag{4.1}
\end{equation*}
$$

and introduce the event, cf. (2.2)

$$
\begin{align*}
G= & \left\{\omega \in \Omega ; L_{n_{0}^{\prime}} \mathbb{Z}^{d} \cap \widetilde{\mathcal{B}}_{n_{0}^{\prime}}(\omega)^{c} \cap\left(5 \mathcal{T}_{n_{0}+1}\right)\right. \text { is contained in the union of } \\
& \text { at most } \left.\widetilde{\ell}_{0} \text { open balls with radius } 3 \widetilde{D}_{n_{0}^{\prime}} \text { and center in } L_{n_{0}^{\prime}} \mathbb{Z}^{d}\right\}, \tag{4.2}
\end{align*}
$$ where $\tilde{\ell}_{0}=\left[\frac{2 M_{0}}{M_{0}(1+a)^{-\left(m_{0}+2\right)}-2 d}\right]+1$.

By analogous considerations as in (3.31), (3.32), we see that on $G^{c}$, we can find $\tilde{\ell}_{0}$ disjoint open balls with centers in $L_{n_{0}^{\prime}} \mathbb{Z}^{d} \cap \widetilde{\mathcal{B}}_{n_{0}^{\prime}}(\omega)^{c} \cap\left(5 \mathcal{T}_{n_{0}+1}\right)$ and radius $\frac{3}{2} \widetilde{D}_{n_{0}^{\prime}}$, so that with (2.3), (1.7), (1.46), (1.47), we see that when $L_{0}$ is large

$$
\begin{align*}
\mathbb{P}\left[G^{c}\right] & \leq c\left(\frac{L_{n_{0}+1}^{2}}{L_{n_{0}^{\prime}}}\right)^{\widetilde{\ell}_{0} d} L_{n_{0}^{\prime}}^{-M_{0} \widetilde{\ell}_{0}} \leq c L_{n_{0}+1}^{\widetilde{\ell}_{0} d\left(2-(1+a)^{-\left(m_{0}+2\right)}\right)-M_{0} \widetilde{\ell}_{0}(1+a)^{-\left(m_{0}+2\right)}}  \tag{4.3}\\
& \leq\left(100\left(m_{0}+2\right)\right)^{-1} L_{n_{0}+1}^{-M_{0}} .
\end{align*}
$$

We introduce the set of finite sequences of length at most $\widetilde{\ell}_{0}$ :

$$
\begin{align*}
\Sigma=\{ & \sigma=\left(\sigma_{1}, \sigma_{2}, \ldots, \sigma_{\widetilde{\ell}}\right) ; \text { with } \\
& \left.0 \leq \widetilde{\ell} \leq \widetilde{\ell}_{0}, \sigma_{i} \in L_{n_{0}^{\prime}} \mathbb{Z}^{d}, B\left(\sigma_{i}, 3 \widetilde{D}_{n_{0}^{\prime}}\right) \cap 5 \mathcal{T}_{n_{0}+1} \neq \emptyset, \text { for } 1 \leq i \leq \widetilde{\ell}\right\} \tag{4.4}
\end{align*}
$$

we denote with $\emptyset$ the only element of $\Sigma$ with length $\widetilde{\ell}=0$. We can now write

$$
\begin{equation*}
G \subset \bigcup_{\sigma \in \Sigma} G_{\sigma}, \text { with } G_{\sigma}=\left\{\omega \in \Omega ;\left(5 \mathcal{T}_{n_{0}+1} \cap L_{n_{0}^{\prime}} \mathbb{Z}^{d}\right) \backslash \bigcup_{i=1}^{\tilde{l}} B\left(\sigma_{i}, 3 \widetilde{D}_{n_{0}^{\prime}}\right) \subseteq \widetilde{\mathcal{B}}_{n_{0}^{\prime}}(\omega)\right\} \tag{4.5}
\end{equation*}
$$

for $\sigma=\left(\sigma_{1}, \ldots, \sigma_{\widetilde{\ell}}\right)$, with $0 \leq \tilde{\ell} \leq \widetilde{\ell}_{0}$.
Loosely speaking, on $G_{\sigma}$ the defects at level $n_{0}^{\prime}$ occurring within $5 \mathcal{T}_{n_{0}+1}$ are contained in the "small set" $\bigcup_{1}^{\tilde{\ell}} B\left(\sigma_{i}, 3 \widetilde{D}_{n_{0}^{\prime}}\right)$. We are now going to perform surgery on these defects. To this end for each $\sigma \in \Sigma$, we choose a $[0,1]$-valued function $g_{\sigma}$ such that with $\sigma=$ $\left(\sigma_{1}, \ldots, \sigma_{\ell}\right), 0 \leq \tilde{\ell} \leq \widetilde{\ell}_{0}$,

$$
\left\{\begin{array}{l}
g_{\sigma}=0 \quad \text { on } \quad \bigcup_{1 \leq i \leq \widetilde{\ell}} \bar{B}\left(\sigma_{i}, 5 \widetilde{D}_{n_{0}^{\prime}}\right) \cup\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}  \tag{4.6}\\
\quad=1 \text { on }\left\{d_{\infty}\left(\cdot,\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}\right) \geq 2 \widetilde{D}_{n_{0}^{\prime}}\right\} \backslash \bigcup_{1 \leq i \leq \tilde{\ell}} B\left(\sigma_{i}, 7 \widetilde{D}_{n_{0}^{\prime}}\right) \\
\left|g_{\sigma}(y)-g_{\sigma}(z)\right| \leq c\left|\frac{y-z}{L_{n_{0}^{\prime}}}\right|, \text { for all } y, z \in \mathbb{R}^{d},
\end{array}\right.
$$

(with the $\beta=1$ analogue to (1.29), one can for instance construct $g_{\sigma}$ as a product of functions attached to each $\sigma_{i}, 1 \leq i \leq \tilde{\ell}$, when $\tilde{\ell} \geq 1$ ).

One can then define the corrected transition kernels for $\sigma \in \Sigma, \omega \in \Omega$ :

$$
\begin{equation*}
R_{n_{0}^{\prime}, \sigma}^{*}=\widetilde{R}_{n_{0}^{\prime}}^{0}+g_{\sigma}\left(\widetilde{R}_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right), \text { cf. }(1.20),(1.21), \tag{4.7}
\end{equation*}
$$

and by induction for $n \in\left[n_{0}^{\prime}, n_{0}\right]$ :

$$
\begin{align*}
& R_{n+1, \sigma}^{*}=\left(R_{n}^{0}+h_{n} S_{n, \sigma}^{*}\right)^{\ell_{n}^{2}} \text {, with } S_{n, \sigma}^{*}=R_{n, \sigma}^{*}-R_{n}^{0} \text {, and } h_{n} \text { functions } \\
& \text { with values in }[0,1] \text {, taking the value } 1 \text { on }\left\{d_{\infty}\left(\cdot,\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}\right) \geq 2 L_{n+1}^{2}\right\} \text {, }  \tag{4.8}\\
& \text { the value } 0 \text { on }\left\{d_{\infty}\left(\cdot,\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}\right) \leq L_{n+1}^{2}\right\} \text {, such that } \sup _{n_{0}^{\prime} \leq n \leq n_{0}}\left|h_{n}\right|_{(n)} \leq c .
\end{align*}
$$

Note that $R_{n_{0}^{\prime}, \sigma}^{*}(x, d y)$ is supported in $\bar{B}\left(x, \widetilde{D}_{n_{0}^{\prime}}\right)$, and when $L_{0}$ is large, it follows by induction that for $n_{0}^{\prime} \leq n \leq n_{0}$ :

$$
\begin{equation*}
R_{n+1, \sigma}^{*}(x, d y)=\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{L_{n+1}^{2} / L_{n_{0}^{\prime}}^{2}}(x, d y), \text { if } d_{\infty}\left(x,\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}\right) \geq 3 L_{n+1}^{2} \tag{4.9}
\end{equation*}
$$

It is also convenient to introduce some further kernels $\widetilde{R}_{n, \sigma}^{*}$ that have a well-localized dependence on the environment, and intuitively are "stopped versions" of the kernels $R_{n, \sigma}^{*}$. For our purpose the difference between these two kernels will be "negligible", cf. (4.140), and (4.12). More precisely, for $x \in \mathbb{R}^{d}$, we consider, (see (1.14) for the notation)

$$
\begin{align*}
& \psi_{n, x}(z) \text { a piecewise linear function of }|z-x| \text {, with value } 1 \text { for } \\
& |z-x| \leq D_{n}^{*} \stackrel{\text { def }}{=} L_{n} e^{3 c_{0}\left(\log \log L_{n}\right)^{2}}, \text { and value } 0 \text { for }|z-x| \geq D_{n}^{*}+1 \tag{4.10}
\end{align*}
$$

We define the probability kernels $\widetilde{R}_{n, \sigma}^{*}$, for $n_{0}^{\prime} \leq n \leq n_{0}$, as

$$
\begin{align*}
\left(\widetilde{R}_{n, \sigma}^{*} f\right)(x)= & \sum_{0 \leq m<L_{n}^{2} / L_{n_{0}^{\prime}}^{2}}\left[\left(\psi_{n, x} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{m}\left(1-\psi_{n, x}\right) f\right](x)  \tag{4.11}\\
& +\left[\left(\psi_{n, x} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{L_{n}^{2} / L_{n_{0}^{\prime}}^{2}} f\right](x), \text { with } f \text { bounded measurable } .
\end{align*}
$$

The kernel $\widetilde{R}_{n, \sigma}^{*}(x, d y)$ corresponds to a "soft stopping" with the function $\psi_{n, \sigma}$ of the Markov chain with kernel $R_{n_{0}^{\prime}, \sigma}^{*}$ starting at $x$, at time $L_{n}^{2} / L_{n_{0}^{\prime}}^{2}$, see also (4.138) for a trajectorial interpretation. In particular $\widetilde{R}_{n_{0}^{\prime}, \sigma}^{*}$ coincides with $R_{n_{0}^{\prime}, \sigma}^{*}$ and $\widetilde{R}_{n, \sigma}^{*}(x, d y)$ is supported in $\bar{B}\left(x, D_{n}^{*}+1+\widetilde{D}_{n_{0}^{\prime}}\right)$. It is also convenient to introduce

$$
\begin{equation*}
\widetilde{S}_{n, \sigma}^{*}=\widetilde{R}_{n, \sigma}^{*}-R_{n}^{0}, \tag{4.12}
\end{equation*}
$$

and we now see that for $n_{0}^{\prime} \leq n \leq n_{0}, x \in \mathbb{R}^{d}$,

$$
\begin{equation*}
\widetilde{R}_{n, \sigma}^{*}(x, d y) \text { or } \widetilde{S}_{n, \sigma}^{*}(x, d y) \text { depend in a } \mathcal{G}_{\bar{B}\left(x, D_{n}^{*}+1+\widetilde{D}_{n_{0}^{\prime}}\right)} \text {-measurable fashion in } \omega . \tag{4.13}
\end{equation*}
$$

In analogy with (1.24), we also define for $\sigma \in \Sigma, \omega \in \Omega, n_{0}^{\prime} \leq n \leq n_{0}, x \in \mathbb{R}^{d}$,

$$
\begin{align*}
\widetilde{d}_{n, \sigma}^{*}(x, \omega) & =\int(y-x) \widetilde{R}_{n, \sigma}^{*}(x, d y)=\int(y-x) \widetilde{S}_{n, \sigma}^{*}(x, d y) \\
\left(\widetilde{\gamma}_{n, \sigma}^{*}\right)^{i, j}(x, \omega) & =\int(y-x)_{i}(y-x)_{j} \widetilde{S}_{n, \sigma}^{*}(x, d y), 1 \leq i, j \leq d \tag{4.14}
\end{align*}
$$

We want to compare $R_{n, \sigma}^{*}$ with $R_{n}^{0}$ on the event $G_{\sigma}$, when starting reasonably away from ( $\left.5 \mathcal{I}_{n_{0}+1}\right)^{c}$, for $n_{0}^{\prime} \leq n \leq n_{0}+1$. Note that with (4.8), using perturbation expansion for $n_{0}^{\prime} \leq n \leq n_{0}$.

$$
\begin{align*}
& S_{n+1, \sigma}^{*}=\left(R_{n}^{0}+h_{n} S_{n, \sigma}^{*}\right)^{\ell_{n}^{2}}-\left(R_{n}^{0}\right)^{\ell_{n}^{2}}+\left(R_{n}^{0}\right)^{\ell_{n}^{2}}-R_{n+1}^{0} \\
& \stackrel{(1.21),(1.54)}{=} \sum_{\substack{k_{0}+\cdots+k_{m}+m=\ell_{n}^{2} \\
k_{i} \geq 0, m \geq 1}}\left(R_{n}^{0}\right)^{k_{0}} h_{n} S_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{k_{1}} h_{n} S_{n, \sigma}^{*} \ldots h_{n} S_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{k_{m}}  \tag{4.15}\\
&+P_{\alpha_{n} L_{n+1}^{2}}-P_{\alpha_{n+1} L_{n+1}^{2}} .
\end{align*}
$$

In essence we are going to first study the "linearized" term corresponding to $m=1$ in the above series, however replacing $S_{n, \sigma}^{*}$, with the more convenient $\widetilde{S}_{n, \sigma}^{*}$, due to their better localization properties. With this in mind, we introduce for $\sigma \in \Sigma, n_{0}^{\prime} \leq n \leq n_{0}$, $v \in L_{n+1} \mathbb{Z}^{d}$, with the notation (1.38), the operator

$$
\begin{equation*}
\widetilde{\mathcal{L}}_{\sigma, n, v}=\sum_{0 \leq k<\ell_{n}^{2}} \chi_{n+1, v}\left(R_{n}^{0}\right)^{k} h_{n, v} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \widetilde{\chi}_{n+1, v} \tag{4.16}
\end{equation*}
$$

where we have used the shorthand notations $h_{n, v}(\cdot)=h_{n}(\cdot) \chi_{D_{n+1}}(\cdot-v)$, and $\widetilde{\chi}_{n+1, v}(\cdot)=$ $\chi_{\widetilde{D}_{n+1}}(\cdot-v)$, cf. (1.37). We also introduce, cf. (1.13), (1.40),

$$
\begin{equation*}
\nu_{n}=2 \bar{\kappa}_{n_{0}^{\prime}}\left(L_{n_{0}^{\prime}}\right)^{-\delta}\left(\frac{L_{n}}{L_{n_{0}^{\prime}}}\right)^{-\beta / 4}, \text { for } n_{0}^{\prime} \leq n \leq n_{0}+1 \tag{4.17}
\end{equation*}
$$

where it should be observed that $\frac{\beta}{4}>\delta$, and $\bar{\kappa}_{n}$ is defined in (2.4). Our first important step comes with

Proposition 4.1. When $L_{0}$ is large, if for some $n \in\left[n_{0}^{\prime}, n_{0}\right]$,

$$
\begin{equation*}
\mathbb{P}\left[\sup _{y \in\left[0, L_{n}\right]^{d}}\left\{\left|\frac{\widetilde{d}_{n, \sigma=\emptyset}^{*}}{L_{n}}(y, \omega)\right| \vee\left|\frac{\widetilde{\widetilde{\gamma}}_{n, \sigma=\emptyset}^{*}}{L_{n}^{2}}(y, \omega)\right|>\nu_{n}\right\}\right] \leq L_{n 0}^{-2} \tag{4.18}
\end{equation*}
$$

then for any $\sigma \in \Sigma, v \in L_{n+1} \mathbb{Z}^{d}$, and event $G_{\sigma, n, v} \subseteq G_{\sigma}$ on which

$$
\begin{equation*}
\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n} \leq \nu_{n}, \text { with } \mathcal{S}_{n, v} \stackrel{\text { def }}{=} L_{n} \mathbb{Z}^{d} \cap\left\{d\left(\cdot, \operatorname{Supp} h_{n, v}\right) \leq 20 \sqrt{d} L_{n}\right\} \tag{4.19}
\end{equation*}
$$

one has

$$
\begin{equation*}
\mathbb{P}\left[G_{\sigma, n, v} \cap\left\{\left\|\widetilde{\mathcal{L}}_{\sigma, n, v}\right\|_{n+1}>\frac{\kappa_{n} \nu_{n}}{\ell_{n}^{\beta / 3}}\right\}\right] \leq e^{-\kappa_{n_{0}}} . \tag{4.20}
\end{equation*}
$$

Proof. Without loss of generality we assume that

$$
\begin{equation*}
h_{n, v} \text { is not identically } 0, \tag{4.21}
\end{equation*}
$$

otherwise there is nothing to prove. We decompose $\widetilde{\mathcal{L}}_{\sigma, n, v}$ into

$$
\begin{equation*}
\widetilde{\mathcal{L}}_{\sigma, n, v}=\mathcal{L}_{A}+\mathcal{L}_{B}+\mathcal{L}_{C}+\mathcal{L}_{D} \tag{4.22}
\end{equation*}
$$

where the operators on the right-hand side of (4.22) are respectively obtained by restricting the summation over $k$ in (4.16) to

$$
\begin{align*}
& I_{A}=\{0\}, I_{B}=\left\{k: 0<k \leq \frac{\ell_{n}^{2}}{2}\right\}, I_{C}=\left\{k: \frac{\ell_{n}^{2}}{2}<k \leq \ell_{n}^{2}-\ell_{n}^{\frac{2}{3} \beta}\right\}  \tag{4.23}\\
& I_{D}=\left\{k: \ell_{n}^{2}-\ell_{n}^{\frac{2}{3} \beta}<k \leq \ell_{n}^{2}-1\right\} .
\end{align*}
$$

We will obtain controls like (4.20) on each term of the decomposition, with the role of $\ell_{n}^{\beta / 3}$ replaced with $\ell_{n}^{1-\beta}$ for $\mathcal{L}_{A}, \ell_{n}^{(1-\beta) \wedge\left(\frac{d}{2}-1\right)}$ for $\mathcal{L}_{B}, \ell_{n}^{\beta / 3}=\ell_{n}^{\beta / 3 \wedge \beta \wedge\left(\frac{d}{2}-1\right)}$ for $\mathcal{L}_{C}$, and $\ell_{n}^{\beta / 3}$ for $\mathcal{L}_{D}$, cf. Lemmas 4.2, 4.3, 4.5, 4.6. We begin with the control of $\mathcal{L}_{A}$.

Lemma 4.2. When $L_{0}$ is large, for $\sigma \in \Sigma$, $n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}$, with (4.21), $\omega \in \Omega$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{A}\right\|_{n+1} \leq \frac{\kappa_{n}}{\ell_{n}^{1-\beta}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) \tag{4.24}
\end{equation*}
$$

Proof. By construction, cf. (4.19), Supp $h_{n, v} \subseteq \bigcup_{x \in \mathcal{S}_{n, v}} \bar{B}\left(x, \sqrt{d} L_{n}\right)$, and for $x \in \mathcal{S}_{n, v}$, $y \in \bar{B}\left(x, 20 \sqrt{d} L_{n}\right), f$ with $|f|_{(n+1)} \leq 1$, one has

$$
\begin{align*}
& \left(\widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-1} \widetilde{\chi}_{n+1, v} f\right)(y)=\left(\widetilde{S}_{n, \sigma}^{*} H\right)(y), \text { with } \\
& H(z)=\left(P_{\alpha_{n}\left(\ell_{n}^{2}-1\right) L_{n}^{2}} \widetilde{\chi}_{n+1, v} f\right)(z)-\left(P_{\alpha_{n}\left(\ell_{n}^{2}-1\right) L_{n}^{2}} \widetilde{\chi}_{n+1, v} f\right)(x), \tag{4.25}
\end{align*}
$$

simply because $\widetilde{S}_{n}^{*} 1=0$. With the help of (1.49), (1.56), we find

$$
\begin{equation*}
|\nabla H| \leq c L_{n+1}^{-1}, \text { and } H(x)=0 \tag{4.26}
\end{equation*}
$$

Using a cut-off function and (A.6) from the Appendix, we can thus find $\widetilde{H}$ such that

$$
\begin{equation*}
\text { Supp } \widetilde{H} \subset B\left(x, 4 D_{n}^{*}\right),|\widetilde{H}| \leq|H|, \widetilde{H}=H \text { on } \bar{B}\left(x, 3 D_{n}^{*}\right), \text { and }|\widetilde{H}|_{(n)} \leq \frac{\kappa_{n}}{\ell_{n}} \tag{4.27}
\end{equation*}
$$

With the remark above (4.12) on the support of $\widetilde{R}_{n, \sigma}^{*}(y, \cdot)$, we see that

$$
\begin{equation*}
\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}(H-\widetilde{H})=-\chi_{n, x} R_{n}^{0}(H-\widetilde{H}) \tag{4.28}
\end{equation*}
$$

and with (1.49), (1.56) and (4.27), we find

$$
\begin{equation*}
\left|\chi_{n, x}\left(\widetilde{S}_{n, \sigma}^{*} H-\widetilde{S}_{n, \sigma}^{*} \widetilde{H}\right)\right|_{(n)} \leq e^{-\kappa_{n}} \tag{4.29}
\end{equation*}
$$

As a result of (4.25), (4.29), we obtain

$$
\left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-1} \widetilde{\chi}_{n+1, v} f\right|_{(n)} \leq\left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*} \widetilde{H}\right|_{(n)}+e^{-\kappa_{n}} \stackrel{(4.27)}{\leq}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n} \frac{\kappa_{n}}{\ell_{n}}+e^{-\kappa_{n}}
$$

Letting the family of functions $h_{n, v} \chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-1} \widetilde{\chi}_{n+1, v} f, x \in \mathcal{S}_{n, v}$ play the role of the $\left(g_{i}\right)_{i \in I}$ in Lemma A. 1 of the Appendix, with (1.29) we find for large $L_{0}$ :

$$
\begin{equation*}
\left|\chi_{n+1, v} h_{n, v} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-1} \widetilde{\chi}_{n+1, v} f\right|_{(n)} \leq \frac{\kappa_{n}}{\ell_{n}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) \tag{4.30}
\end{equation*}
$$

and since $\left\|\mathcal{L}_{A}\right\|_{n+1} \leq \ell_{n}^{\beta}\left\|\mathcal{L}_{A}\right\|_{n},(4.24)$ follows.
We now turn to the control of $\mathcal{L}_{D}$.
Lemma 4.3. When $L_{0}$ is large, for $\sigma \in \Sigma, n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}$ with (4.21), $\omega \in \Omega$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{D}\right\|_{n+1} \leq \frac{\kappa_{n}}{\ell_{n}^{\beta / 3}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) \tag{4.31}
\end{equation*}
$$

Proof. Note that for $k<\ell_{n}^{2}$, cf. (4.22), and $f$ with $|f|_{(n+1)} \leq 1$, with (1.55), (1.49), $\left|\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \widetilde{\chi}_{n+1, v} f\right|_{(n+1)} \leq c$. Hence for $x \in \mathcal{S}_{n, v}$, repeating the construction used in Lemma 4.2, we can find $\widetilde{H}$ with $\operatorname{Supp} \widetilde{H} \subset B\left(x, 4 D_{n}^{*}\right),|\widetilde{H}|_{(n)} \leq \kappa_{n} \ell_{n}^{-\beta}$ such that

$$
\begin{equation*}
\left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \widetilde{\chi}_{n+1, v} f-\chi_{n, x} \widetilde{S}_{n, \sigma}^{*} \widetilde{H}\right|_{(n)} \leq e^{-\kappa_{n}} \tag{4.32}
\end{equation*}
$$

With $L_{0}$ large we thus find that

$$
\begin{equation*}
\left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \widetilde{\chi}_{n+1, v} f\right|_{(n)} \leq \frac{\kappa_{n}}{\ell_{n}^{\beta}}\left(\sup _{x^{\prime} \in \mathcal{S}_{n, v}}\left\|\chi_{n, x^{\prime}}, \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.33}
\end{equation*}
$$

Note also that with (1.49), (1.56), for $t \geq \alpha_{n} L_{n+1}^{2} / 2$,

$$
\begin{equation*}
\left|P_{t} g\right|_{(n+1)} \leq c|g|_{\infty}, \text { when } g \text { is bounded measurable } \tag{4.34}
\end{equation*}
$$

so that for each $k \in I_{D}$,

$$
\begin{equation*}
\left|\chi_{n+1, v}\left(R_{n}^{0}\right)^{k} h_{n, v} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \widetilde{\chi}_{n+1, v} f\right|_{(n+1)} \leq \frac{\kappa_{n}}{\ell_{n}^{\beta}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.35}
\end{equation*}
$$

Since $\left|I_{D}\right| \leq \ell_{n}^{\frac{2}{3} \beta}$, summing over $k \in I_{D}$, we obtain (4.31).
We continue with the analysis of $\mathcal{L}_{C}$ and $\mathcal{L}_{B}$. We first need to recall some facts related to Taylor's formula. For $g$ a $C^{2}$-function on $\mathbb{R}^{d}$, Taylor's formula with integral remainder of order 2 states that for $y, z \in \mathbb{R}^{d}$ :

$$
\begin{equation*}
g(y+z)=g(y)+\sum_{|\gamma| \leq 2} \frac{1}{\gamma!} D^{\gamma} g(y) z^{\gamma}+r_{g}(y, z) \tag{4.36}
\end{equation*}
$$

where $\gamma=\left(\gamma_{1}, \ldots, \gamma_{d}\right)$ is a multi-index, $|\gamma|=\gamma_{1}+\cdots+\gamma_{d}, \gamma!=\gamma_{1}!\ldots \gamma_{d}!, z^{\gamma}=z_{1}^{\gamma_{1}} \ldots z_{d}^{\gamma_{d}}$, and

$$
\begin{equation*}
r_{g}(y, z)=\int_{0}^{1} 3(1-t)^{2} \sum_{|\gamma|=3} \frac{1}{\gamma!} D^{\gamma} g(y+t z) z^{\gamma} d t \tag{4.37}
\end{equation*}
$$

and otherwise hopefully obvious notations. We recall the definition (4.14), and the notation (1.54).
Lemma 4.4. When $L_{0}$ is large, $\sigma \in \Sigma$, $n_{0}^{\prime} \leq n \leq n_{0}, \omega \in \Omega$, for $1 \leq j \leq \ell_{n}^{2}, x \in L_{n} \mathbb{Z}^{d}$, $|y-x| \leq 10 \sqrt{d} L_{n}, f$ bounded measurable,

$$
\begin{align*}
\int \widetilde{S}_{n, \sigma}^{*}(y, d z)\left[\left(R_{n}^{0}\right)^{j} f\right](z) & =\widetilde{d}_{n, \sigma}^{*}(y, \omega) \cdot\left(D P_{\alpha_{n} j L_{n}^{2}} f\right)(y)  \tag{4.38}\\
& +\frac{1}{2} \widetilde{\gamma}_{n, \sigma}^{*}(y, \omega) \cdot\left(D^{2} P_{\alpha_{n} j L_{n}^{2}} f\right)(y)+H_{j, f}(y),
\end{align*}
$$

and

$$
\begin{equation*}
\left|H_{j, f}(y)\right| \leq c\left(1+\frac{\left(D_{n}^{*}\right)^{1-\beta}}{\sqrt{j} L_{n}^{1-\beta}}\right) \frac{D_{n}^{* 3}}{j^{\frac{3}{2}} L_{n}^{3}}\left(\frac{|f|_{1}}{\left(\sqrt{j} L_{n}\right)^{d}} \wedge|f|_{\infty}\right)\left(\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.39}
\end{equation*}
$$

Proof. With (4.36), (1.21), we can write:

$$
\left(R_{n}^{0}\right)^{j} f(y+z)=P_{\alpha_{n} j L_{n}^{2}} f(y)+\left(D P_{\alpha_{n} j L_{n}^{2}} f\right)(y) \cdot z+\frac{1}{2}\left(D^{2} P_{\alpha_{n} j L_{n}^{2}} f(y) \cdot z \otimes z+r_{j, f, y}(z),\right.
$$

and $r_{j, f, y}(\cdot-y)$ coincides in $\bar{B}\left(y, 3 D_{n}^{*}\right)$ with $\widetilde{r}(\cdot)$ which is supported in $B\left(y, 4 D_{n}^{*}\right)$, and such that

$$
\begin{equation*}
|\widetilde{r}|_{(n)} \leq c \frac{D_{n}^{* 3}}{j^{\frac{3}{2}} L_{n}^{3}}\left[\frac{|f|_{1}}{\left(\sqrt{j} L_{n}\right)^{d}} \wedge|f|_{\infty}\right]\left(1+\frac{D_{n}^{*(1-\beta)}}{\sqrt{j} L_{n}^{(1-\beta)}}\right) . \tag{4.40}
\end{equation*}
$$

Indeed with (4.37), (1.49), (1.56):

$$
\begin{equation*}
\sup _{x \in B\left(y, 5 D_{n}^{*}\right)}\left|r_{j, f, y}(x-y)\right| \leq a_{j, n} \stackrel{\text { def }}{=} c \frac{D_{n}^{* 3}}{j^{\frac{3}{2}} L_{n}^{3}}\left[\frac{|f|_{1}}{\left(\sqrt{j} L_{n}\right)^{d}} \wedge|f|_{\infty}\right], \tag{4.41}
\end{equation*}
$$

and for $w, w^{\prime} \in B\left(y, 5 D_{n}^{*}\right)$,

$$
\begin{align*}
& \left|r_{j, f, y}(w-y)-r_{j, f, y}\left(w^{\prime}-y\right)\right| \leq c \sup _{0 \leq t \leq 1,|\gamma|=3} \mid\left(D^{\gamma} P_{\alpha_{n} j L_{n}^{2}} f\right)(y+t(w-y)) \cdot \\
& (w-y)^{\gamma}-\left(D^{\gamma} P_{\alpha_{n} j L_{n}^{2}} f\right)\left(y+t\left(w^{\prime}-y\right)\right) \cdot\left(w^{\prime}-y\right)^{\gamma} \mid \leq \\
& c \sup _{\substack{z \in B\left(y, 5 D_{n}^{*}\right) \\
|\gamma|=3}}\left|D^{\gamma} P_{\alpha_{n} j L_{n}^{2}} f(z)\right| D_{n}^{* 2}\left|w-w^{\prime}\right|+  \tag{4.42}\\
& c\left|w-w^{\prime}\right| \sup _{\substack{z \in B\left(y, 5 D_{n}^{*}\right) \\
|\gamma|=4}}\left|D^{\gamma} P_{\alpha_{n} j L_{n}^{2}} f(z)\right| D_{n}^{* 3} \stackrel{(1.56)}{\leq} \\
& \frac{c\left|w-w^{\prime}\right|}{D_{n}^{*}}\left(\frac{D_{n}^{* 3}}{j^{\frac{3}{2}} L_{n}^{3}}+\frac{D_{n}^{* 4}}{j^{2} L_{n}^{4}}\right)\left(\frac{|f|_{1}}{\left(\sqrt{j} L_{n}\right)^{d}} \wedge|f|_{\infty}\right) \leq c\left|\frac{w-w^{\prime}}{D_{n}^{*}}\right|^{\beta} a_{j, n}\left(1+\frac{D_{n}^{*}}{\sqrt{j} L_{n}}\right) .
\end{align*}
$$

So using a cut-off function, we obtain the claim (4.40). Since $\widetilde{R}_{n, \sigma}^{*}(y, d z)$ is supported in $\bar{B}\left(y, 3 D_{n}^{*}\right)$, cf. above (4.12),

$$
\begin{equation*}
\left|\int \widetilde{S}_{n, \sigma}^{*}(y, d z)\left(r_{j, f, y}(z-y)-\widetilde{r}(z)\right)\right|=\left|\int R_{n}^{0}(y, d z)\left(r_{j, f, y}(z-y)-\widetilde{r}(z)\right)\right| \leq c a_{j, n} e^{-\kappa_{n}} \tag{4.43}
\end{equation*}
$$

using Cauchy-Schwarz's inequality, (1.49), (1.59) in the last step.
Taking into account that $\chi_{n, x}(y)=1$, and

$$
H_{j, f}(y)=\int \widetilde{S}_{n, \sigma}^{*}(y, d z) r_{j, f, y}(z-y)
$$

the claim (4.39) now follows from the above inequality and (4.40).
We now decompose $\mathcal{L}_{C}$, cf. (4.22), into

$$
\begin{equation*}
\mathcal{L}_{C}=\mathcal{L}_{C}^{1}+\mathcal{L}_{C}^{2} \tag{4.44}
\end{equation*}
$$

where in the notations of (4.14), (4.38)

$$
\begin{aligned}
\mathcal{L}_{C}^{1} f(y)=\sum_{k \in I_{C}} \chi_{n+1, v}(y)\{ & \left(R_{n}^{0}\right)^{k}\left(h _ { n , v } ( \cdot ) \left[\widetilde{d}_{n, \sigma}^{*}(\cdot, \omega) \cdot\left(D P_{\alpha_{n}\left(\ell_{n}^{2}-k-1\right)} \widetilde{\chi}_{n+1, v} f\right)(\cdot)\right.\right. \\
& \left.\left.\left.+\frac{1}{2} \widetilde{\gamma}_{n, \sigma}^{*}(\cdot, \omega) \cdot\left(D^{2} P_{\alpha_{n}\left(\ell_{n}^{2}-k-1\right)} \widetilde{\chi}_{n+1, v} f\right)(\cdot)\right]\right)\right\}(y)
\end{aligned}
$$

and

$$
\mathcal{L}_{C}^{2} f(y)=\sum_{k \in I_{C}} \chi_{n+1, v}(y) \int\left(R_{n}^{0}\right)^{k}(y, d z) h_{n, v}(z) H_{\ell_{n}^{2}-k-1, \tilde{\chi}_{n+1, v}}(z)
$$

Our next step comes with

Lemma 4.5. When $L_{0}$ is large, $\sigma \in \Sigma$, $n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}$ with (4.21), $\omega \in \Omega$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{C}^{2}\right\|_{n+1} \leq \frac{\kappa_{n}}{\ell_{n}^{\beta / 3}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.45}
\end{equation*}
$$

Moreover, if $n$ is as in (4.18), with the notation (4.17) and above (4.19)

$$
\begin{equation*}
\mathbb{P}\left[G_{\sigma, n, v} \cap\left\{\left\|\mathcal{L}_{C}^{1}\right\|_{n+1} \geq \frac{\kappa_{n}}{\ell_{n}^{\beta \wedge\left(\frac{d}{2}-1\right)}} \nu_{n}\right\}\right] \leq e^{-\kappa_{n_{0}}} \tag{4.46}
\end{equation*}
$$

Proof. We begin with the proof of (4.45). We choose $f$ with $|f|_{(n+1)} \leq 1$, and deduce from (4.39) and (4.34), that

$$
\begin{equation*}
\left\|\mathcal{L}_{C}^{2}\right\|_{n+1} \leq \sum_{k \in I_{C}} \kappa_{n}\left(\ell_{n}^{2}-k-1\right)^{-\frac{3}{2}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.47}
\end{equation*}
$$

Noting that $\sum_{j \geq \ell_{n}^{2 \beta / 3}} j^{-3 / 2} \leq c \ell_{n}^{-\beta / 3}$, we find (4.45).
We then turn to the proof of (4.46). We further decompose $\mathcal{L}_{C}^{1}$ into

$$
\begin{equation*}
\mathcal{L}_{C}^{1}=\sum_{\gamma \in\{0,1\}^{d}} \mathcal{L}_{C, \gamma}+\mathcal{L}_{C}^{\prime}, \tag{4.48}
\end{equation*}
$$

where

$$
\begin{align*}
\mathcal{L}_{C, \gamma} f(y) & =\sum_{q \in \Lambda_{\gamma}, k \in I_{C}} \chi_{n+1, v}(y) \Phi_{q, k}(f)(y),  \tag{4.49}\\
\mathcal{L}_{C}^{\prime} f(y) & =\sum_{q \in \Lambda^{\prime}, k \in I_{C}} \chi_{n+1, v}(y) \Phi_{q, k}(f)(y), \tag{4.50}
\end{align*}
$$

and we have used the notations for $q \in \mathbb{Z}^{d}, k \geq 0$,

$$
\begin{align*}
\Phi_{q, k}(f)(y)=\int_{B_{q}} P_{\alpha_{n} k L_{n}^{2}}(y, d z) & h_{n, v}(z)\left[\widetilde{d}_{n, \sigma}^{*}(z, \omega) \cdot\left(D P_{\alpha_{n}\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}} \widetilde{\chi}_{n+1, v} f\right)(z)\right.  \tag{4.51}\\
& \left.+\frac{1}{2} \widetilde{\gamma}_{n, \sigma}^{*}(z, \omega) \cdot\left(D^{2} P_{\alpha_{n}\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}} \widetilde{\chi}_{n+1, v} f\right)(z)\right]
\end{align*}
$$

Note that in view of (1.7) and (4.13), for $f$ bounded measurable and $\gamma \in\{0,1\}^{d}$,

$$
\begin{equation*}
\left\{\left(\Phi_{q, k}(f)\right)_{0 \leq k \leq \ell_{n}^{2}}\right\} \text { are independent under } \mathbb{P} \text {, as } q \text { varies over } \Lambda_{\gamma} . \tag{4.55}
\end{equation*}
$$

Note also that when $L_{0}$ is large, with $\sigma, n, v$ as above (4.45) and $\gamma \in\{0,1\}^{d}$, by the properties of the support of $h_{n, v}$, cf. below (4.16),

$$
\begin{equation*}
\left|\Lambda_{\gamma}\right| \leq c\left(\frac{D_{n+1}}{D_{n}^{*}}\right)^{d} \leq \kappa_{n} \ell_{n}^{d},\left|\Lambda^{\prime}\right| \leq c \tag{4.56}
\end{equation*}
$$

We use wavelets, see (1.34), to control $\left\|\mathcal{L}_{C}^{1}\right\|_{n+1}$, and recall from Proposition A. 2 in the Appendix that for $\gamma \in\{0,1\}^{d}$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{C, \gamma}\right\|_{n+1} \leq c \sup _{\alpha, \ell, p} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \frac{1}{2^{d \ell}}\left|\left\langle\theta_{\alpha, \ell, p}, \mathcal{L}_{C, \gamma} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right|, \tag{4.57}
\end{equation*}
$$

where the supremum runs over $\alpha \in\{0,1\}^{d}, \ell \leq J_{n+1}$, cf. (A.7), $p \in \mathbb{Z}^{d}$, with $\alpha \neq 0$, when $\ell<J_{n+1}$, and similar constraints for $\alpha^{\prime}, \ell^{\prime}, p^{\prime}$ in the sum. An analogous inequality holds for $\mathcal{L}_{C}^{\prime}$ in place of $\mathcal{L}_{C, \gamma}$. From now we consider triplets

$$
\begin{align*}
& (\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right) \text { satisfying the above conditions and such that } \\
& \operatorname{Supp} \theta_{\alpha, \ell, p} \cap \operatorname{Supp} \chi_{n+1, v} \neq \emptyset, \text { and } \operatorname{Supp} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \cap \operatorname{Supp} \widetilde{\chi}_{n+1, v} \neq \emptyset, \tag{4.58}
\end{align*}
$$

cf. below (4.16) for the notation.
Given $\gamma \in\{0,1\}^{d}$, we introduce an enumeration $q_{j}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$, of $\Lambda_{\gamma}$. We then define for $0 \leq j \leq\left|\Lambda_{\gamma}\right|$

$$
\begin{equation*}
M_{j}=2^{-d \ell} \sum_{k \in I_{C}, j^{\prime} \leq j}\left\langle\theta_{\alpha, \ell, p}, \chi_{n+1, v} \psi_{j^{\prime}, k}\right\rangle, \text { for } j \geq 1, M_{0}=0 \tag{4.59}
\end{equation*}
$$

where in the notations of (4.51),

$$
\begin{equation*}
\psi_{j, k}(y) \stackrel{\text { def }}{=} \Phi_{q_{j}, k}\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)(y) \tag{4.60}
\end{equation*}
$$

We now bound $\left|M_{j}-M_{j-1}\right|$, first when $\omega \in G_{\sigma, n, v}$, cf. above (4.19), and then for a general $\omega \in \Omega$. Note that with analogous arguments as in the proof of Lemma 2.1, in view of (4.14), (4.19), (1.49), for $\omega \in G_{\sigma, n, v}$,

$$
\begin{equation*}
\text { for } y \in \operatorname{Supp} h_{n, v},\left|\widetilde{d}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} \nu_{n} L_{n},\left|\widetilde{\gamma}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} \nu_{n} L_{n}^{2} \tag{4.61}
\end{equation*}
$$

In addition to (4.58), let us first assume that

$$
\begin{equation*}
2^{\ell^{\prime}} \leq L_{n} \tag{4.62}
\end{equation*}
$$

Then for $y, y^{\prime} \in B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap \operatorname{Supp} \theta_{\alpha, \ell, p}, \omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$, with the help of (1.49), (1.56), (4.61), in view of (4.51), (4.60), we find when $L_{0}$ is large:

$$
\begin{align*}
& \sum_{k \in I_{C}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \frac{c 2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \sum_{k \in I_{C}} \\
& {\left[\kappa_{n} \nu_{n} L_{n}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} L_{n}^{-1} \exp \left\{-\frac{c\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|^{2}}{\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\}\right.}  \tag{4.63}\\
& \left(2^{\ell^{\prime}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} L_{n}^{-1}\right)^{d}+\kappa_{n} \nu_{n} L_{n}^{2}\left(\ell_{n}^{2}-k-1\right)^{-1} L_{n}^{-2} \\
& \left.\exp \left\{-\frac{c \mid 2^{\ell^{\prime} p^{\prime}-\left.10 D_{n}^{*} q_{j}\right|^{2}}}{\left.\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\}\left(2^{\ell^{\prime}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} L_{n}^{-1}\right)^{d}\right],
\end{align*}
$$

where in the expression inside the exponential we made use of (4.62), of Supp $\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \subseteq$ $B\left(2^{\ell^{\prime}} p^{\prime}, c 2^{\ell^{\prime}}\right)$, and of $\left(\ell_{n}^{2}-k-1\right)^{1 / 2} L_{n} \geq D_{n}^{*}$, for $k \in I_{C}$. Hence the left-hand side of (4.63) is smaller than:

$$
\begin{align*}
& \frac{c 2^{\ell}}{L_{n+1}} \kappa_{n} \nu_{n}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \sum_{k \in I_{C}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}}  \tag{4.64}\\
& \exp \left\{-\frac{c\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|^{2}}{\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\}\left(2^{\ell^{\prime}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} L_{n}^{-1}\right)^{d}
\end{align*}
$$

Using a comparison with $\int_{0}^{\infty} s^{-\rho} e^{-u / s} d s$, we find that

$$
\begin{equation*}
\text { for } \rho>1, u>0, \sum_{1 \leq k<\infty} k^{-\rho} \exp \left\{-\frac{u}{k}\right\} \leq c(\rho)\left(u^{-(\rho-1)} \wedge 1\right) \tag{4.65}
\end{equation*}
$$

so that

$$
\sum_{k \in I_{C}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{(d+1)}{2}} \exp \left\{-\frac{c \mid 2^{\ell^{\prime} p^{\prime}-\left.10 D_{n}^{*} q_{j}\right|^{2}}}{\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\} \leq c\left[\left(\frac{L_{n}}{\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|}\right)^{d-1} \wedge 1\right],
$$

and coming back to (4.63), (4.64), we find that for $\omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|, y, y^{\prime} \in$ $B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap \operatorname{Supp} \theta_{\alpha, \ell, p}:$
and with entirely analogous bounds

$$
\begin{equation*}
\sum_{k \in I_{C}}\left|\psi_{j, k}(y)\right| \leq \kappa_{n} \nu_{n}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{d}\left[\left(\frac{L_{n}}{\left.\left.\mid 2^{\ell^{\prime} p^{\prime}-10 D_{n}^{*} q_{j} \mid}\right)^{d-1} \wedge 1\right] . . ~ . ~}\right.\right. \tag{4.67}
\end{equation*}
$$

We now replace (4.62) with:

$$
\begin{equation*}
L_{n}<2^{\ell^{\prime}} \leq L_{n+1} \tag{4.68}
\end{equation*}
$$

We then write for $y, y^{\prime} \in B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap \operatorname{Supp} \theta_{\alpha, \ell, p}, \omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{align*}
& \sum_{k \in I_{C}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \frac{c 2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \kappa_{n} \nu_{n} \\
& {\left[\sum_{k \in I_{C},}\left(\ell_{n}^{2}-k-1\right)^{-\frac{(d+1)}{2}} \exp \left\{-\frac{c\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|^{2}}{\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{d}+\right.}  \tag{4.69}\\
& 2^{2^{\prime} \leq\left(\ell_{n}^{2}-k-1\right)^{1 / 2} L_{n}} \\
& \left.\sum_{k \in I_{C},}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} \exp \left\{-\frac{c\left(\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|-c 2^{\ell^{\prime}}\right)_{+}^{2}}{\left(\ell_{n}^{2}-k-1\right) L_{n}^{2}}\right\}\right], \\
& 2^{\ell^{\prime}>\left(\ell_{n}^{2}-k-1\right)^{1 / 2} L_{n}}
\end{align*}
$$

where we omitted the intermediary step, cf. (4.63), where terms corresponding to $\widetilde{d_{n, \sigma}^{*}}$ and $\widetilde{\gamma}_{n, \sigma}^{*}$ are separately bounded. Note that

$$
\begin{equation*}
\quad \sum_{\substack{k \in I_{C}, 2^{\ell^{\prime}} \leq\left(\ell_{n}^{2}-k-1\right)^{1 / 2} L_{n}}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{(d+1)}{2}} \leq c\left(\frac{L_{n}}{2^{\ell^{\prime}}}\right)^{d-1}, \sum_{\substack{k \in I_{C}, 2^{\ell^{\prime}}>\left(\ell_{n}^{2}-k-1\right)^{1 / 2} L_{n}}}\left(\ell_{n}^{2}-k-1\right)^{-\frac{1}{2}} \leq c \frac{2^{\ell^{\prime}}}{L_{n}} . \tag{4.70}
\end{equation*}
$$

These inequalities together with (4.65) show that for $\omega \in G_{\sigma, n, v}, y, y^{\prime} \in B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap$ $\operatorname{Supp} \theta_{\alpha, \ell, p}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$, with (4.68) we have:

$$
\begin{align*}
& \sum_{k \in I_{C}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \kappa_{n} \nu_{n}\left(\frac{2^{\ell}}{L_{n+1}}\right)\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d}\left[\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{d}\right. \\
& \left.\left\{\left(\frac{L_{n}}{2^{\ell^{\prime}}}\right)^{d-1} \wedge\left(\frac{L_{n}}{\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|}\right)^{d-1}\right\}+\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right) \exp \left\{-c \frac{\left(\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|-c 2^{\ell^{\prime}}\right)_{+}^{2}}{2^{2 \ell^{\prime}}}\right\}\right] \leq  \tag{4.71}\\
& \kappa_{n} \nu_{n}\left(\frac{2^{\ell}}{L_{n+1}}\right)\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \frac{2^{\ell^{\prime}}}{L_{n}}\left[1 \wedge \left(\frac{2^{\ell^{\prime}}}{\left.\left.\mid 2^{\ell^{\prime} p^{\prime}-10 D_{n}^{*} q_{j} \mid}\right)^{d-1}+\exp \left\{-c\left(\frac{\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|}{2^{\ell^{\prime}}}\right)^{2}\right\}\right]}\right.\right.
\end{align*}
$$

and with entirely similar estimates we also have in this situation

$$
\begin{align*}
& \sum_{k \in I_{C}}\left|\psi_{j, k}(y)\right| \leq \kappa_{n} \nu_{n}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \frac{2^{\ell^{\prime}}}{L_{n}}\left[1 \wedge\left(\frac{2^{\ell^{\prime}}}{\left|2^{\ell^{\prime}} p^{\prime}-10 D_{n}^{*} q_{j}\right|}\right)^{d-1}+\right.  \tag{4.72}\\
&\left.\exp \left\{-c\left(\frac{\mid 2^{2^{\prime} p^{\prime}-10 D_{n}^{*} q_{j} \mid}}{2^{\ell^{\prime}}}\right)^{2}\right\}\right] .
\end{align*}
$$

Using the fact that $\int \theta_{\alpha, \ell, p}(y) d y=0$, when $\alpha \neq 0$, cf. (A.12), we see collecting (4.66), (4.67), (4.71), (4.72) that for large $L_{0}, \omega \in G_{\sigma, n, v}, \gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58),

$$
\begin{equation*}
\left|M_{j}-M_{j-1}\right| \leq \delta_{\ell, \ell^{\prime}}(j), 1 \leq j \leq\left|\Lambda_{\gamma}\right|, \tag{4.73}
\end{equation*}
$$

where up to a constant multiplicative factor, $\delta_{\ell, \ell^{\prime}}(j)$ is given by the right-hand side of (4.66) when $2^{\ell^{\prime}} \leq L_{n}$, and by the last member of (4.71) when $L_{n}<2^{\ell^{\prime}} \leq L_{n+1}$.

Observe that when we consider a general $\omega$ in place of $\omega \in G_{\sigma, n, v}$, as above, we can use analogous bounds with the only difference that (4.61) is now replaced with:

$$
\begin{equation*}
\left|\widetilde{d}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} L_{n},\left|\widetilde{\gamma}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} L_{n}^{2}, \text { for } \sigma \in \Sigma, y \in \mathbb{R}^{d}, \omega \in \Omega \tag{4.74}
\end{equation*}
$$

Hence we find that for $\omega \in \Omega, \gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58),

$$
\begin{equation*}
\left|M_{j}-M_{j-1}\right| \leq \kappa_{n} \nu_{n}^{-1} \delta_{\ell^{\prime}, \ell^{\prime}}(j), 1 \leq j \leq\left|\Lambda_{\gamma}\right| \tag{4.75}
\end{equation*}
$$

Now for $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), we introduce the conditional probability:

$$
\begin{equation*}
\widetilde{\mathbb{P}}(\cdot)=\mathbb{P}\left[\cdot| | M_{j}-M_{j-1}\left|\leq \delta_{\ell, \ell^{\prime}}(j), 1 \leq j \leq\left|\Lambda_{\gamma}\right|\right],\right. \tag{4.76}
\end{equation*}
$$

and denote with $\widetilde{\mathbb{E}}$ the corresponding expectation. We note that thanks to the independence under $\mathbb{P}$ of the increments $M_{j}-M_{j-1}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$, cf. (4.55), these increments are independent under $\widetilde{\mathbb{P}}$ as well. We will now bound

$$
\begin{equation*}
\Delta_{j} \stackrel{\text { def }}{=} \widetilde{\mathbb{E}}\left[M_{j}-M_{j-1}\right], 1 \leq j \leq\left|\Lambda_{\gamma}\right| \tag{4.77}
\end{equation*}
$$

First note that for $y \in \bigcup_{q \in \Lambda_{\gamma}} B_{q}$, cf. (4.52), (4.54), with $L_{0}$ large, we can replace $R_{n_{0}^{\prime}, \sigma}^{*}$ with $\widetilde{R}_{n_{0}^{\prime}}$ in the right-hand side of (4.11), when calculating $\widetilde{d}_{n, \sigma}^{*}(y, \omega), \widetilde{\gamma}_{n, \sigma}^{*}(y, \omega)$ in (4.14). So by isotropy, cf. (1.12), for $y \in \bigcup_{q \in \Lambda_{\gamma}} B_{q}$ :

$$
\begin{equation*}
\mathbb{E}\left[\widetilde{d_{n, \sigma}^{*}}(y, \omega)\right]=0 \tag{4.78}
\end{equation*}
$$

Moreover for $y$ in the same set, with $1 \leq i, j \leq d$, we have

$$
\begin{equation*}
\mathbb{E}\left[\left(\widetilde{\gamma}_{n, \sigma}^{*}\right)^{i, j}(y, \omega)\right] \stackrel{(1.25)}{=} \mathbb{E}\left[\left(\widetilde{\gamma}_{n, \sigma}^{*}\right)^{i, j}(y, \omega)-\left(\widetilde{\gamma}_{n}\right)^{i, j}(y, \omega)\right] . \tag{4.79}
\end{equation*}
$$

On the event where for all $x \in L_{n} \mathbb{Z}^{d}$ with Supp $\chi_{n, x} \cap B\left(y, 3 D_{n}^{*}\right) \neq \emptyset, x \in \widetilde{B}_{n}(\omega)$, and all $x^{\prime} \in L_{n_{0}^{\prime}} \mathbb{Z}^{d}$ with Supp $\chi_{n_{0}^{\prime}, x} \cap B\left(y, 3 D_{n}^{*}\right) \neq \emptyset, x^{\prime} \in \mathcal{B}_{n_{0}^{\prime}}(\omega)$, the integrand in the righthand side of (4.79), using the remark above (4.78), the strong Markov property, and the localization estimate in (2.2), is bounded in absolute value by

$$
c D_{n}^{* 2}\left(\left(\frac{L_{n}}{L_{n_{0}^{\prime}}}\right)^{2} e^{-\kappa_{n_{0}^{\prime}}}+e^{-\kappa_{n}}\right) \leq e^{-\kappa_{n_{0}}}, \text { with } L_{0} \text { large } .
$$

Bounding with (1.47) the probability of the complement of this event, we see that for large $L_{0}, \gamma \in\{0,1\}^{d}, y \in \bigcup_{q \in \Lambda_{\gamma}} B_{q}$,

$$
\begin{align*}
\left|\mathbb{E}\left[\widetilde{\gamma}_{n, \sigma}^{*}(y, \omega)\right]\right| & \leq c D_{n}^{* 2}\left[\left(\frac{D_{n}^{*}}{L_{n_{0}^{\prime}}}\right)^{d} L_{n_{0}^{\prime}}^{-M_{0}}+\kappa_{n} L_{n}^{-M_{0}}\right]+e^{-\kappa_{n_{0}}}  \tag{4.80}\\
& \leq \kappa_{n_{0}} L_{n_{0}}^{(2+d)-M_{0}(1+a)^{-\left(m_{0}+1\right)}} \stackrel{(1.46)}{\leq} L_{n_{0}}^{-10} .
\end{align*}
$$

We then observe that the bounds we derived below (4.61) until (4.73) show that when $1 \leq j \leq\left|\Lambda_{\gamma}\right|$, with $\kappa_{n}$ as in (4.61), $(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58),

$$
\begin{align*}
& \text { if }\left|\widetilde{d}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} \nu_{n} L_{n},\left|\widetilde{\gamma}_{n, \sigma}^{*}(y, \omega)\right| \leq \kappa_{n} \nu_{n} L_{n}^{2} \text {, for all } y \in B_{q_{j}} \text {, then }  \tag{4.81}\\
& \left|M_{j}-M_{j-1}\right| \leq \delta_{\ell, \ell^{\prime}}(j)
\end{align*}
$$

Hence on the event $\left\{\left|M_{j}-M_{j-1}\right|>\delta_{\ell, \ell^{\prime}}(j)\right\}$, for some $y \in B_{q_{j}}$, (4.81) does not hold, and by the remark above (4.78), we can replace $\sigma$ with $\emptyset(\in \Sigma)$, when negating (4.81). We thus find with (4.18) that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ with (4.58), for $1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\mathbb{P}\left[\left|M_{j}-M_{j-1}\right|>\delta_{\ell, \ell^{\prime}}(j)\right] \leq c \frac{\left|B_{q_{j}}\right|}{L_{n}^{d}} L_{n_{0}}^{-2} \stackrel{(4.52)}{\leq} \kappa_{n} L_{n_{0}}^{-2} \tag{4.82}
\end{equation*}
$$

Coming back to (4.78), (4.80), to replace (4.61), the estimates (4.61) until (4.73) now show that with $1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\left|\mathbb{E}\left[M_{j}-M_{j-1}\right]\right| \leq\left(\kappa_{n} \nu_{n} L_{n}^{2}\right)^{-1} L_{n_{0}}^{-10} \delta_{\ell, \ell^{\prime}}(j) \leq L_{n_{0}}^{-10} \delta_{\ell, \ell^{\prime}}(j) \tag{4.83}
\end{equation*}
$$

and noting that

$$
\mathbb{E}\left[M_{j}-M_{j-1}\right] \stackrel{(4.77)}{=} \Delta_{j} \mathbb{P}\left[\left|M_{j}-M_{j-1}\right| \leq \delta_{\ell, \ell^{\prime}}(j)\right]+\mathbb{E}\left[M_{j}-M_{j-1},\left|M_{j}-M_{j-1}\right|>\delta_{\ell, \ell^{\prime}}(j)\right],
$$

we obtain from (4.75), (4.82), (4.83), that for $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), $1 \leq j \leq\left|\Lambda_{j}\right|:$

$$
\begin{equation*}
\left|\Delta_{j}\right| \leq 2\left(L_{n_{0}}^{-10}+\kappa_{n} L_{n_{0}}^{-2} \nu_{n}^{-1}\right) \delta_{\ell, \ell^{\prime}}(j) \leq L_{n_{0}}^{-1} \delta_{\ell, \ell^{\prime}}(j) \stackrel{\text { def }}{=} \widetilde{\delta}_{\ell, \ell^{\prime}}(j) . \tag{4.84}
\end{equation*}
$$

Observe that under $\widetilde{\mathbb{P}}, M_{\left|\Lambda_{\gamma}\right|}-\sum_{j=1}^{\left|\Lambda_{\gamma}\right|} \Delta_{j}$ is a sum of $\left|\Lambda_{\gamma}\right|$ independent variables respectively bounded by $2 \delta_{\ell, \ell^{\prime}}(j)$. Note also that when $2^{\ell^{\prime}} \leq L_{n}$, by (4.66), (4.67)

$$
\begin{align*}
\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}} & \leq \kappa_{n} \nu_{n} \frac{2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{d}  \tag{4.85}\\
& \leq \kappa_{n} \nu_{n} \frac{2^{\ell}}{L_{n+1}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d} \stackrel{\text { def }}{=} \sigma_{n}\left(\ell, \ell^{\prime}\right),
\end{align*}
$$

whereas for $L_{n}<2^{\ell^{\prime}} \leq L_{n+1}$, with (4.71), (4.72)

$$
\begin{align*}
\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell^{\prime}, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}} & \leq \kappa_{n} \nu_{n} \frac{2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \frac{2^{\ell^{\prime}}}{L_{n}}\left[\left(\frac{2^{\ell^{\prime}}}{D_{n}^{*}}\right)^{d}+\left(\frac{2^{\ell^{\prime}}}{D_{n}^{*}}\right)^{d}\right]^{\frac{1}{2}}  \tag{4.86}\\
& \leq \kappa_{n} \nu_{n} \ell_{n}^{-d} \frac{2^{\ell}}{L_{n+1}}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{\frac{d}{2}+1} \stackrel{\text { def }}{=} \sigma_{n}\left(\ell, \ell^{\prime}\right)
\end{align*}
$$

Note also that when $L_{0}$ is large, for $\ell, \ell^{\prime} \leq J_{n+1}, \gamma \in\{0,1\}^{d}$ :

$$
\begin{align*}
& \sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \widetilde{\delta}_{\ell, \ell^{\prime}}(j) \underset{\text { Cauchy-Schwarz }}{\stackrel{(4.84)}{\leq}}\left|\Lambda_{\gamma}\right|^{\frac{1}{2}} L_{n_{0}}^{-1}\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \sigma_{\ell, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}}  \tag{4.87}\\
& \stackrel{(4.56),(1.14)}{\leq} \frac{1}{2} \sigma_{n}\left(\ell, \ell^{\prime}\right) .
\end{align*}
$$

We thus see that for $u \geq \sigma_{n}\left(\ell, \ell^{\prime}\right)$, with a slight variation of Azuma's inequality, cf. [1], or [30], p. 308,

$$
\begin{align*}
& \mathbb{P}\left[\left|M_{\left|\Lambda_{\gamma}\right|}\right| \geq u, G_{\sigma, n, v}\right] \leq \widetilde{\mathbb{P}}\left[\left|M_{\left|\Lambda_{\gamma}\right|}\right| \geq u\right] \leq \\
& \widetilde{\mathbb{P}}\left[\left|M_{\left|\Lambda_{\gamma}\right|}-\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \Delta_{j}\right| \geq u-\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \widetilde{\delta}_{\ell, \ell^{\prime}}(j)\right] \leq 2 \exp \left\{-\frac{1}{32}\left(\frac{u}{\sigma_{n}\left(\ell, \ell^{\prime}\right)}\right)^{2}\right\} . \tag{4.88}
\end{align*}
$$

If we define for $\gamma \in\{0,1\}^{d}$ the event

$$
\begin{align*}
& G_{\sigma, n, v, C, \gamma}=G_{\sigma, n, v} \cap\left\{\text { for }(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right), \text { as in }(4.58),\right. \\
& \left.\frac{1}{2^{d \ell}}\left|\left\langle\theta_{\alpha, \ell, p}, \mathcal{L}_{C, \gamma} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right| \leq \sigma_{n}\left(\ell, \ell^{\prime}\right)\left(1+\ell_{-}+\ell_{-}^{\prime}\right) e^{\left(\log \log L_{n}\right)^{2}}\right\}, \tag{4.89}
\end{align*}
$$

$\left(\ell_{-}, \ell_{-}^{\prime}\right.$ denote the respective negative parts of $\left.\ell, \ell^{\prime}\right)$, we see that when $L_{0}$ is large

$$
\begin{align*}
& \mathbb{P}\left[G_{\sigma, n, v} \backslash G_{\sigma, n, v, C, \gamma}\right] \leq \\
& \sum_{\ell, \ell^{\prime} \leq J_{n+1}} c\left(\frac{D_{n+1}}{2^{\ell}}\right)^{d}\left(\frac{\widetilde{D}_{n+1}}{2^{\ell^{\prime}}}\right)^{d} \exp \left\{-\frac{1}{32} e^{2\left(\log \log L_{n}\right)^{2}}\left(1+\ell_{-}+\ell_{-}^{\prime}\right)^{2}\right\} \leq  \tag{4.90}\\
& c\left(\sum_{\ell \leq J_{n+1}}\left(\frac{\widetilde{D}_{n+1}}{2^{\ell}}\right)^{d} \exp \left\{-\frac{1}{64} e^{2\left(\log \log L_{n}\right)^{2}}\left(1+\ell_{-}^{2}\right)\right\}\right)^{2} \leq e^{-\kappa_{n_{0}}} .
\end{align*}
$$

Observe that on $G_{\sigma, n, v, C, \gamma}$ in view of (4.57) one has

$$
\begin{equation*}
\left\|\mathcal{L}_{C, \gamma}\right\|_{n+1} \leq \Gamma^{\prime} \stackrel{\text { def }}{=} c \sup _{\alpha, \ell, p} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \sigma_{n}\left(\ell, \ell^{\prime}\right) e^{\left(\log \log L_{n}\right)^{2}}\left(1+\ell_{-}+\ell_{-}^{\prime}\right) \tag{4.91}
\end{equation*}
$$

with $(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ varying over the set described in (4.58). We now write:

$$
\begin{equation*}
\Gamma^{\prime} \leq \Gamma_{1}^{\prime}+\Gamma_{2}^{\prime} \tag{4.92}
\end{equation*}
$$

where $\Gamma_{1}^{\prime}$ corresponds to the expression in the right-hand side of (4.91) with $2^{\ell^{\prime}} \leq L_{n}$, and $\Gamma_{2}^{\prime}$ to the expression with $L_{n}<2^{\ell^{\prime}} \leq L_{n+1}$. We thus see that for large $L_{0}$,

$$
\begin{align*}
\Gamma_{1}^{\prime} & \stackrel{(4.85)}{\leq} \kappa_{n} \nu_{n} \sup _{2^{\ell} \leq L_{n+1}} \frac{2^{\ell}}{L_{n+1}} \sum_{2^{\ell^{\prime}} \leq L_{n}, p^{\prime}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left(1+\ell_{-}+\ell_{-}^{\prime}\right) \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \\
& \leq \kappa_{n} \nu_{n} \sup _{2^{\ell} \leq L_{n+1}}\left(\frac{2^{\ell}}{L_{n+1}}\right)^{1-\beta}\left(1+\ell_{-}\right) \sum_{2^{\ell^{\prime} \leq L_{n}}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{\beta}\left(1+\ell_{-}^{\prime}\right) \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{\beta}} . \tag{4.93}
\end{align*}
$$

On the other hand, (recall $\ell_{-}^{\prime}=0$, when $L_{n}<2^{\ell^{\prime}} \leq L_{n+1}$ ):

$$
\begin{align*}
\Gamma_{2}^{\prime} & \stackrel{(4.86)}{\leq} \kappa_{n} \nu_{n} \sup _{2^{\ell} \leq L_{n+1}} \frac{2^{\ell}}{L_{n+1}}\left(1+\ell_{-}\right) \sum_{L_{n}<2^{\ell^{\prime}} \leq L_{n+1}, p^{\prime}} \ell_{n}^{-d}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{\frac{d}{2}+1} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \\
& \leq \kappa_{n} \nu_{n} \sup _{2^{\ell} \leq L_{n+1}}\left(\frac{2^{\ell}}{L_{n+1}}\right)^{1-\beta}\left(1+\ell_{-}\right) \sum_{L_{n}<2^{\ell^{\prime}} \leq L_{n+1}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{\beta} \ell_{n}^{-d}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{\frac{d}{2}+1}\left(\frac{\widetilde{D}_{n+1}}{2^{\ell^{\prime}}}\right)^{d}  \tag{4.94}\\
& \leq \kappa_{n} \nu_{n} \sup _{2^{\ell} \leq L_{n+1}}\left(\frac{2^{\ell}}{L_{n+1}}\right)^{1-\beta}\left(1+\ell_{-}\right) \sum_{L_{n}<2^{\ell^{\prime}} \leq L_{n+1}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{\beta}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{-\frac{d}{2}+1} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{\beta \wedge\left(\frac{d}{2}-1\right)}} .
\end{align*}
$$

Combining (4.93), (4.94), we see that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d}$, on $G_{\sigma, n, v, C, \gamma}$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{C, \gamma}\right\|_{n+1} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{\beta \wedge\left(\frac{d}{2}-1\right)}} \tag{4.95}
\end{equation*}
$$

We now turn to the study of $\mathcal{L}_{C}^{\prime}$. Keeping in mind that $\left|\Lambda^{\prime}\right| \leq c$, cf. (4.53), using similar estimates as in (4.66), (4.67), (4.71), (4.72), we see that for large $L_{0}$, with $(\alpha, \ell, p)$, $\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), and for $\omega \in G_{\sigma, n, v}$ :
(4.96) $\frac{1}{2^{d \ell}}\left|\left\langle\theta_{\alpha, \ell, p}, \mathcal{L}_{C}^{\prime} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right| \leq \begin{cases}\kappa_{n} \nu_{n} \frac{2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{d}, & \text { for } 2^{\ell^{\prime}} \leq L_{n}, \\ \kappa_{n} \nu_{n} \frac{2^{\ell}}{L_{n+1}}\left(\frac{D_{n}^{*}}{L_{n+1}}\right)^{d} \frac{2^{\ell^{\prime}}}{L_{n}}, & \text { for } L_{n}<2^{\ell^{\prime}} \leq L_{n+1} .\end{cases}$

By direct inspection in (4.85), (4.86), we see that the above right-hand side is bounded by $\kappa_{n} \sigma_{n}\left(\ell, \ell^{\prime}\right)$. Hence the analogous bound as in (4.57), for $\mathcal{L}_{C}^{\prime}$, as well as (4.91) - (4.94), now prove that when $L_{0}$ is large, for $\omega \in G_{\sigma, n, v}$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{C}^{\prime}\right\|_{n+1} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{\beta \wedge\left(\frac{d}{2}-1\right)}} \tag{4.97}
\end{equation*}
$$

Collecting (4.90), (4.95), (4.97), we have completed the proof of (4.46).
We continue with the analysis of $\mathcal{L}_{B}$. In analogy with (4.44), and with $I_{B}$ replacing $I_{C}$ there, we write:

$$
\begin{equation*}
\mathcal{L}_{B}=\mathcal{L}_{B}^{1}+\mathcal{L}_{B}^{2} \tag{4.98}
\end{equation*}
$$

Lemma 4.6. When $L_{0}$ is large, $\sigma \in \Sigma$, $n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}$ with (4.21), $\omega \in \Omega$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{B}^{2}\right\|_{n+1} \leq \frac{\kappa_{n}}{\ell_{n}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right) . \tag{4.99}
\end{equation*}
$$

Moreover if $n$ is as in (4.18) with the notations (4.17) and above (4.19),

$$
\begin{equation*}
\mathbb{P}\left[G_{\sigma, n, v} \cap\left\{\left\|\mathcal{L}_{B}^{1}\right\|_{n+1} \geq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{(1-\beta) \wedge\left(\frac{d}{2}-1\right)}}\right\}\right] \leq e^{-\kappa_{n_{0}}} . \tag{4.100}
\end{equation*}
$$

Proof. We begin with the proof of (4.99). Note that with (1.49), (1.56), for $g$ bounded measurable,

$$
\begin{equation*}
\left|\chi_{n+1, v} P_{\alpha_{n} k L_{n}^{2}} g\right|_{(n+1)} \leq \frac{c \ell_{n}}{\sqrt{k}}|g|_{\infty}, \text { for } 1 \leq k \leq \ell_{n}^{2} \tag{4.101}
\end{equation*}
$$

hence with (4.39) we find

$$
\begin{align*}
\left\|\mathcal{L}_{B}^{2}\right\|_{n+1} & \leq \sum_{k \in I_{B}} \frac{c \ell_{n}}{\sqrt{k}} \frac{\kappa_{n}}{\ell_{n}^{3}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right)  \tag{4.102}\\
& \leq \frac{\kappa_{n}}{\ell_{n}}\left(\sup _{x \in \mathcal{S}_{n, v}}\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}+e^{-\kappa_{n}}\right)
\end{align*}
$$

This proves (4.99).
We continue with the proof of (4.100). In analogy to (4.48), and with $I_{B}$ replacing $I_{C}$ there, we decompose $\mathcal{L}_{B}^{1}$ into:

$$
\begin{equation*}
\mathcal{L}_{B}^{1}=\sum_{\gamma \in\{0,1\}^{d}} \mathcal{L}_{B, \gamma}+\mathcal{L}_{B}^{\prime}, \tag{4.103}
\end{equation*}
$$

For $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ satisfying (4.58), we introduce in full analogy with (4.59), with $I_{B}$ replacing $I_{C}$ there, $M_{j}, 0 \leq j \leq\left|\Lambda_{\gamma}\right|$. With the definition (4.60), we observe that for large $L_{0}$, when

$$
\begin{equation*}
2^{\ell} \leq L_{n} \tag{4.104}
\end{equation*}
$$

for $y, y^{\prime} \in B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap \operatorname{Supp} \theta_{\alpha, \ell, p}, \omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{j}\right|$, with the help of (1.56), (4.61),

$$
\begin{align*}
& \sum_{k \in I_{B}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \\
& \sum_{k \in I_{B}} c \frac{D_{n}^{* d}}{L_{n}^{d}} \frac{1}{k^{\frac{d+1}{2}}} \frac{\left|y-y^{\prime}\right|}{L_{n}} \exp \left\{-\frac{c A_{j}\left(y, y^{\prime}\right)^{2}}{k L_{n}^{2}}\right\} \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}, \tag{4.105}
\end{align*}
$$

with

$$
\begin{equation*}
A_{j}\left(y, y^{\prime}\right)=\inf \left\{|w-\widetilde{w}|, w \in B_{q_{j}}, \widetilde{w}=\lambda y+(1-\lambda) y^{\prime}, 0 \leq \lambda \leq 1\right\} \tag{4.106}
\end{equation*}
$$

As a result of (4.65), under the above assumptions:

$$
\begin{equation*}
\sum_{k \in I_{B}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left[\left(\frac{L_{n}}{A_{j}\left(y, y^{\prime}\right)}\right)^{d-1} \wedge 1\right] \frac{\left|y-y^{\prime}\right|}{L_{n}} \tag{4.107}
\end{equation*}
$$

and by an analogous calculation

$$
\begin{equation*}
\sum_{k \in I_{B}}\left|\psi_{j, k}(y)\right| \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left[\left(\frac{L_{n}}{A_{j}(y)}\right)^{d-2} \wedge 1\right], \text { with } A_{j}(y) \stackrel{\text { def }}{=} d\left(y, B_{q_{j}}\right) \tag{4.108}
\end{equation*}
$$

If we now turn to the case

$$
\begin{equation*}
L_{n}<2^{\ell} \leq L_{n+1} \tag{4.109}
\end{equation*}
$$

under the same conditions as stated above (4.105), we find

$$
\begin{align*}
& \quad \sum_{k \in I_{B}, \sqrt{k} L_{n}>2^{\ell}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left|\frac{y-y^{\prime}}{L_{n}}\right| \\
& \sum_{2^{\ell}<\sqrt{k} L_{n} \leq L_{n+1}} k^{-\frac{(d+1)}{2}} \exp \left\{-\frac{c A_{j}\left(y, y^{\prime}\right)^{2}}{k L_{n}^{2}}\right\} . \tag{4.110}
\end{align*}
$$

Note that one has the following refinement of (4.65):

$$
\begin{equation*}
\sum_{v<k} k^{-\rho} \exp \left\{-\frac{u}{k}\right\} \leq c(\rho)\left\{(u \vee v)^{-(\rho-1)} \wedge 1\right\}, \text { for } u, v>0, \rho>1 \tag{4.111}
\end{equation*}
$$

that is obtained by considering the case $u=0$, and using (4.65). Hence for large $L_{0}$, when (4.109) holds, for $y, y^{\prime} \in B\left(v, 20 \sqrt{d} L_{n+1}\right) \cap \operatorname{Supp} \theta_{\alpha, \ell, p}, \omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\sum_{k \in I_{B}, \sqrt{k} L_{n}>2^{\ell}}\left|\psi_{j, k}(y)-\psi_{j, k}\left(y^{\prime}\right)\right| \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d} \frac{\left|y-y^{\prime}\right|}{L_{n}}\left\{\left(\frac{L_{n}}{2^{\ell} \vee A_{j}\left(y, y^{\prime}\right)}\right)^{d-1} \wedge 1\right\} \tag{4.112}
\end{equation*}
$$

and in an analogous fashion:

$$
\begin{equation*}
\sum_{k \in I_{B}, \sqrt{k} L_{n}>2^{\ell}}\left|\psi_{j, k}(y)\right| \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left\{\left(\frac{L_{n}}{2^{\ell} \vee A_{j}(y)}\right)^{d-2} \wedge 1\right\} . \tag{4.113}
\end{equation*}
$$

On the other hand with (4.60), (4.51):

$$
\begin{aligned}
& \sum_{k \in I_{B}, \sqrt{k} L_{n} \leq 2^{\ell}} \frac{1}{2^{\ell d}} \int_{\operatorname{Supp} \theta_{\alpha, \ell, p}}\left|\psi_{j, k}(y)\right| d y \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d} \\
& \sum_{k \in I_{B}, \sqrt{k} L_{n} \leq 2^{\ell}} \frac{1}{2^{\ell^{\ell d}}} \int_{B_{q_{j}}} d z \int_{B\left(2^{\ell} p, 2^{\ell}\right)} d y \frac{c}{\left(k L_{n}^{2}\right)^{d / 2}} \exp \left\{-\frac{c(z-y)^{2}}{k L_{n}^{2}}\right\} \leq \\
& \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell}}{L_{n}}\right)^{2}\left(\frac{D_{n}^{*}}{2^{\ell}}\right)^{d} \exp \left\{-c\left(\frac{A_{j}\left(2^{\ell} p\right)}{2^{\ell}}\right)^{2}\right\} .
\end{aligned}
$$

Collecting our bounds, we thus see that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), $\omega \in G_{\sigma, n, v}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\left|M_{j}-M_{j-1}\right| \leq \delta_{\ell, p, \ell^{\prime}}(j) \tag{4.115}
\end{equation*}
$$

where for $2^{\ell} \leq L_{n}, 2^{\ell^{\prime}} \leq L_{n+1}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\delta_{\ell, p, \ell^{\prime}}(j)=\frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left[\frac{2^{\ell}}{L_{n+1}}\left\{\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{d-2} \wedge 1\right\}+\frac{2^{\ell}}{L_{n}}\left\{\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{d-1} \wedge 1\right\}\right] \tag{4.116}
\end{equation*}
$$

with

$$
A_{j, \ell, p}=\inf \left\{|w-\widetilde{w}|, w \in B_{q_{j}}, \widetilde{w} \in B\left(2^{\ell} p, c 2^{\ell}\right)\right\}
$$

with $c$ such that $\operatorname{Supp} \theta_{\alpha}(\cdot) \subseteq B(0, c)$, for all $\alpha \in\{0,1\}^{d}$, and we have made use of the fact that since $2^{\ell} \leq L_{n}, \alpha \neq 0$, and in view of (A.12), $\int \theta_{\alpha, \ell, p}(y) d y=0$.

On the other hand when $L_{n}<2^{\ell} \leq L_{n+1}, 2^{\ell^{\prime}} \leq L_{n+1}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{align*}
& \delta_{\ell, p, \ell^{\prime}}(j)=\frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left[\frac{2^{\ell}}{L_{n+1}}\left(\frac{L_{n}}{2^{\ell} \vee A_{j, \ell, p}}\right)^{d-2}+\frac{2^{\ell}}{L_{n}}\left(\frac{L_{n}}{2^{\ell} \vee A_{j, \ell, p}}\right)^{d-1}\right.  \tag{4.117}\\
& \left.+\left(\frac{L_{n}}{2^{\ell}}\right)^{d-2} \exp \left\{-c\left(\frac{A_{j, \ell, p}}{2^{\ell}}\right)^{2}\right\}\right] .
\end{align*}
$$

Arguing as above (4.75), we see that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d},(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), for $\omega \in \Omega, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$ :

$$
\begin{equation*}
\left|M_{j}-M_{j-1}\right| \leq \kappa_{n} \nu_{n}^{-1} \delta_{\ell, p, \ell^{\prime}}(j) . \tag{4.118}
\end{equation*}
$$

Keeping the same notation $\widetilde{\mathbb{P}}$ and $\Delta_{j}, 1 \leq j \leq\left|\Lambda_{\gamma}\right|$, as in (4.76), (4.77), with the only difference that $\delta_{\ell, p, \ell^{\prime}}(j)$ replaces $\delta_{\ell, \ell^{\prime}}(j)$ in (4.76), repeating the argument leading to (4.84), we see that for large $L_{0}$, under the same conditions as above (4.118)

$$
\begin{equation*}
\left|\Delta_{j}\right| \leq L_{n_{0}}^{-1} \delta_{\ell, p, \ell^{\prime}}(j) \stackrel{\text { def }}{=} \widetilde{\delta}_{\ell, p, \ell^{\prime}}(j), 1 \leq j \leq\left|\Lambda_{\gamma}\right| \tag{4.119}
\end{equation*}
$$

Keeping in mind the objective of deriving bounds that parallel (4.88), we now bound $\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell, p, \ell^{\prime}}^{2}(j)\right)^{1 / 2}$. To this end note first that for $2^{\ell} \leq L_{n}, p, \ell^{\prime}$ compatible with (4.58), cf. (4.116),

$$
\begin{align*}
& \left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell, p, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}} \leq \\
& \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left[\frac{2^{\ell}}{L_{n+1}}\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{2(d-2)} \wedge 1\right)^{\frac{1}{2}}+\right.  \tag{4.120}\\
& \left.\quad\left(\frac{2^{\ell}}{L_{n}}\right)\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{2(d-1)} \wedge 1\right)^{\frac{1}{2}}\right] .
\end{align*}
$$

Observe that with (4.54), and the notations below (4.116),

$$
\text { i) } \begin{align*}
\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{2(d-2)} \wedge 1 \leq \kappa_{n} \ell_{n}^{2 \nu(d)}, \text { with } \nu(d) & =\frac{1}{2}, \text { when } d=3 \\
& =0, \text { when } d \geq 4 \tag{4.121}
\end{align*}
$$

$$
\text { ii) } \sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{A_{j, \ell, p}}\right)^{2(d-1)} \wedge 1 \leq c
$$

As a result we obtain that for $2^{\ell} \leq L_{n}, p, \ell^{\prime}$ compatible with (4.58):

$$
\begin{equation*}
\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell, p, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}} \leq \sigma_{n}\left(\ell, \ell^{\prime}\right) \stackrel{\text { def }}{=} \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d} \frac{2^{\ell}}{L_{n}} . \tag{4.122}
\end{equation*}
$$

To handle the case $L_{n}<2^{\ell} \leq L_{n+1}$, observe that:

$$
\begin{align*}
& \text { i) } \sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \exp \left\{-c\left(\frac{A_{j, \ell, p}}{2^{\ell}}\right)^{2}\right\} \leq c\left(\frac{2^{\ell}}{L_{n}}\right)^{d} \\
& \text { ii) } \sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{2^{\ell} \vee A_{j, \ell, p}}\right)^{2(d-1)} \leq c\left(\frac{L_{n}}{2^{\ell}}\right)^{2(d-1)}\left(\frac{2^{\ell}}{D_{n}^{*}}\right)^{d}+c\left(\frac{L_{n}}{2^{\ell}}\right)^{d-2} \leq c\left(\frac{L_{n}}{2^{\ell}}\right)^{d-2}  \tag{4.123}\\
& \text { iii) } \sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|}\left(\frac{L_{n}}{2^{\ell} \vee A_{j, \ell, p}}\right)^{2(d-2)} \leq c\left(\frac{L_{n}}{2^{\ell}}\right)^{2(d-2)}\left(\frac{2^{\ell}}{D_{n}^{*}}\right)^{d}+\sum_{c 2^{\ell}<i D_{n}^{*} \leq c \widetilde{D}_{n+1}} c i^{-(d-3)} \\
& \quad \leq \kappa_{n} \ell_{n}^{2 \nu(d)}, \text { with the notation of }(4.121) .
\end{align*}
$$

Coming back to (4.117), we obtain for $L_{n}<2^{\ell} \leq L_{n+1}, p, \ell^{\prime}$ compatible with (4.58):

$$
\begin{align*}
\left(\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \delta_{\ell, p, \ell^{\prime}}(j)^{2}\right)^{\frac{1}{2}} & \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left[\left(\frac{2^{\ell}}{L_{n+1}}\right) \ell_{n}^{\nu(d)}+\frac{2^{\ell}}{L_{n}}\left(\frac{L_{n}}{2^{\ell}}\right)^{\frac{d}{2}-1}+\left(\frac{L_{n}}{2^{\ell}}\right)^{\frac{d}{2}-2}\right]  \tag{4.124}\\
& \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left[\left(\frac{2^{\ell}}{L_{n+1}}\right) \ell_{n}^{\nu(d)}+\left(\frac{2^{\ell}}{L_{n}}\right)^{\nu(d)}\right] \stackrel{\text { def }}{=} \sigma_{n}\left(\ell, \ell^{\prime}\right)
\end{align*}
$$

The same argument leading to (4.87), (4.88) shows that when $L_{0}$ is large, $\ell, \ell^{\prime} \leq J_{n+1}$, $\gamma \in\{0,1\}^{d}$ :

$$
\begin{equation*}
\sum_{1 \leq j \leq\left|\Lambda_{\gamma}\right|} \widetilde{\delta}_{\ell, p, \ell^{\prime}}(j) \leq \frac{1}{2} \sigma_{n}\left(\ell, \ell^{\prime}\right) \tag{4.125}
\end{equation*}
$$

and for $u \geq \sigma_{n}\left(\ell, \ell^{\prime}\right)$,

$$
\begin{equation*}
\mathbb{P}\left[\left|M_{\left|\Lambda_{\gamma}\right|}\right| \geq u, G_{\sigma, n, v}\right] \leq 2 \exp \left\{-\frac{1}{32}\left(\frac{u}{\sigma_{n}\left(\ell, \ell^{\prime}\right)}\right)^{2}\right\} \tag{4.126}
\end{equation*}
$$

We can now introduce for $\gamma \in\{0,1\}^{d}$ the event

$$
\begin{align*}
G_{\sigma, n, v, B, \gamma}=G_{\sigma, n, v} & \cap\left\{\text { for }(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right) \text { as in }(4.58),\right. \\
& \left.\frac{1}{2^{d \ell}}\left|\left\langle\theta_{\alpha, \ell, p}, \mathcal{L}_{B, \gamma}^{1} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right| \leq \sigma_{n}\left(\ell, \ell^{\prime}\right)\left(1+\ell_{-}+\ell_{-}^{\prime}\right) e^{\left(\log \log L_{n}\right)^{2}}\right\}, \tag{4.127}
\end{align*}
$$

and find that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d}$, similarly to (4.90),

$$
\begin{equation*}
\mathbb{P}\left[G_{\sigma, n, v} \backslash G_{\sigma, n, v, B, \gamma}\right] \leq e^{-\kappa_{n_{0}}} \tag{4.128}
\end{equation*}
$$

Moreover on the event $G_{\sigma, n, v, B, \gamma}$, we have

$$
\begin{equation*}
\left\|\mathcal{L}_{B, \gamma}^{1}\right\|_{n+1} \leq \Gamma \stackrel{\text { def }}{=} c \sup _{\alpha, \ell, p} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \sigma_{n}\left(\ell, \ell^{\prime}\right)\left(1+\ell_{-}+\ell_{-}^{\prime}\right) e^{\left(\log \log L_{n}\right)^{2}} \tag{4.129}
\end{equation*}
$$

with $(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ varying over the set described in (4.58). We now write:

$$
\begin{equation*}
\Gamma \leq \Gamma_{1} \vee \Gamma_{2} \tag{4.130}
\end{equation*}
$$

with $\Gamma_{1}$ defined as $\Gamma$ with the additional requirement $2^{\ell} \leq L_{n}$, and $\Gamma_{2}$ with the additional requirement $L_{n}<2^{\ell} \leq L_{n+1}$, instead. With (4.122), we find for large $L_{0}$ :

$$
\begin{align*}
\Gamma_{1} & \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}} \sup _{2^{\ell} \leq L_{n}} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell}}{L_{n}}\right)\left(1+\ell_{-}^{\prime}+\ell_{-}^{\prime}\right) \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \\
& \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}} \sup _{2^{\ell} \leq L_{n}}\left(\frac{2^{\ell}}{L_{n}}\right)^{1-\beta}\left(1+\ell_{-}\right) \sum_{2^{\ell^{\prime}} \leq L_{n+1}}\left(1+\ell_{-}^{\prime}\right)\left(\frac{2^{\ell^{\prime}}}{L_{n}}\right)^{\beta} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{(1-\beta)}}, \tag{4.131}
\end{align*}
$$

whereas with (4.124), we find, (recall $\ell_{-}=0$, when $L_{n}<2^{\ell} \leq L_{n+1}$ ):

$$
\begin{align*}
\Gamma_{2} & \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}} \sup _{L_{n}<2^{\ell} \leq L_{n+1}} \sum_{2^{\ell^{\prime}} \leq L_{n+1}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}}\left[\frac{2^{\ell}}{L_{n+1}} \ell_{n}^{\nu(d)}+\left(\frac{2^{\ell}}{L_{n}}\right)^{\nu(d)}\right]\left(1+\ell_{-}^{\prime}\right) \\
& \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}} \sup _{L_{n}<2^{\ell} \leq L_{n+1}}\left[\ell_{n}^{\nu(d)}\left(\frac{2^{\ell}}{L_{n+1}}\right)^{1-\beta}+\left(\frac{L_{n+1}}{2^{\ell}}\right)^{\beta}\left(\frac{2^{\ell}}{L_{n}}\right)^{\nu(d)}\right]  \tag{4.132}\\
& \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{1-(\beta \vee \nu(d))}}=\frac{\kappa_{n} \nu_{n}}{\ell_{n}^{(1-\beta) \wedge\left(\frac{d}{2}-1\right)}} .
\end{align*}
$$

Coming back to (4.129), we see that when $L_{0}$ is large, for $\gamma \in\{0,1\}^{d}$, on $G_{\sigma, n, v, B, \gamma}$, we have

$$
\begin{equation*}
\left\|\mathcal{L}_{B, \gamma}^{1}\right\|_{n+1} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{(1-\beta) \wedge\left(\frac{d}{2}-1\right)}} \tag{4.133}
\end{equation*}
$$

We now turn to the study of $\mathcal{L}_{B}^{\prime}$. Keeping in mind that $\left|\Lambda^{\prime}\right| \leq c$, cf. (4.53), using similar estimates as in (4.115), (4.116), (4.117), we see that for large $L_{0}$, with $(\alpha, \ell, p),\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ as in (4.58), and for $\omega \in G_{\sigma, n, v}$ :

$$
2^{-d \ell}\left|\left\langle\theta_{\alpha, \ell, p}, \mathcal{L}_{B}^{\prime} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right| \leq\left\{\begin{array}{l}
\frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell^{\prime}}}{L_{n+1}}\right)^{d} \frac{2^{\ell}}{L_{n}}, \text { if } 2^{\ell} \leq L_{n},  \tag{4.134}\\
\frac{\kappa_{n} \nu_{n}}{\ell_{n}}\left(\frac{2^{\ell}}{L_{n+1}}\right)^{d}\left(\frac{2^{\ell}}{L_{n+1}}+\left(\frac{L_{n}}{2^{\ell}}\right)^{d-2}\right), \\
\\
\text { if } L_{n}<2^{\ell} \leq L_{n+1} .
\end{array}\right.
$$

By direct inspection, cf. (4.122), (4.124), we see that the right-hand side above is bounded by $\kappa_{n} \sigma_{n}\left(\ell, \ell^{\prime}\right)$. Hence the analogous bound to (4.129) for $\mathcal{L}_{B}^{\prime}$, as well as (4.131), (4.132) show that when $L_{0}$ is large, for $\omega \in G_{\sigma, n, v}$ :

$$
\begin{equation*}
\left\|\mathcal{L}_{B}^{\prime}\right\|_{n+1} \leq \frac{\kappa_{n} \nu_{n}}{\ell_{n}^{(1-\beta) \wedge\left(\frac{d}{2}-1\right)}} \tag{4.135}
\end{equation*}
$$

Combining (4.128), (4.133), (4.135), we have proved (4.100).
Collecting Lemmas 4.2, 4.3, 4.5, 4.6, we see that we have proved Proposition 4.1.
Remark 4.7. As a result of Lemmas 4.2, 4.3, 4.5, 4.6, we see that with high probability on $G_{\sigma, n, v},\left\|\widetilde{\mathcal{L}}_{\sigma, n, v}\right\|_{n+1}$ is smaller than $\kappa_{n} \nu_{n}$ by the crucial contraction factor $\ell_{n}^{-\beta / 3 \wedge(1-\beta) \wedge(d / 2+1)}$ $\left(=\ell_{n}^{-\beta / 3}\right.$, with our choice $\beta \in\left(0, \frac{1}{2}\right]$ in (1.13)). In the proof of Proposition 4.1, there is an asymmetry in the role of $k$ close to 0 and $k$ close to $\ell_{n}^{2}-1$ in the decomposition (4.22), which stems from the use of Taylor's formula to second order, cf. (4.38). In a loose sense, if the $\widetilde{S}_{n, \sigma}^{*}$ in the definition of $\widetilde{\mathcal{L}}_{\sigma, n, v}$ in (4.16) had been centered under $\mathbb{P}$, we could have avoided Taylor's expansion, and chosen in (4.22), $I_{A}=\{0\}, I_{B}=\left\{k: 0<k \leq \ell_{n}^{2} / 2\right\}$, $I_{C}=\left\{k: \ell_{n}^{2} / 2<k<\ell_{n}^{2}-1\right\}, I_{D}=\left\{\ell_{n}^{2}-1\right\}$. With the proper assumptions, the role of $\ell_{n}^{-\beta / 3 \wedge(1-\beta) \wedge(d / 2-1)}$ would then have been replaced with $\ell_{n}^{-\beta \wedge(1-\beta) \wedge\left(\frac{d}{2}-1\right)}$, displaying a higher symmetry between the role of small $k$ and $k$ close to $\ell_{n}^{2}-1$. Ultimately the asymmetry in the proof results from the fact that we work with $\widetilde{S}_{n}$ which compares $\widetilde{R}_{n}$ to the Gaussian kernel $R_{n}^{0}$, rather than separately analyzing $\widetilde{R}_{n}-\mathbb{E}\left[\widetilde{R}_{n}\right]$ and $\mathbb{E}\left[\widetilde{R}_{n}\right]-R_{n}^{0}$.

Our next objective, see the comments above (4.10), is to control $\left\|h_{n}\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right)\right\|_{n}$. To this end we introduce the event, cf. (4.2):

$$
\begin{gather*}
\widetilde{G}=G \cap \bigcap_{n_{0}^{\prime}<n \leq n_{0}}\left\{\omega \in \Omega ; L_{n} \mathbb{Z}^{d} \cap \widetilde{\mathcal{B}}_{n}(\omega)^{c} \cap\left(5 \mathcal{T}_{n_{0}+1}\right)\right. \text { is contained in the }  \tag{4.136}\\
\text { union of at most } \widetilde{\ell}_{0} \text { open balls with radius } 3 \widetilde{D}_{n} \text { and } \\
\text { center in } \left.L_{n} \mathbb{Z}^{d}\right\} .
\end{gather*}
$$

The same estimates as in (4.3), show that for large $L_{0}$,

$$
\begin{equation*}
\mathbb{P}\left[\widetilde{G}^{c}\right] \leq\left(n_{0}-n_{0}^{\prime}+1\right)\left(100\left(m_{0}+2\right)\right)^{-1} L_{n_{0}+1}^{-M_{0}} \leq \frac{1}{100} L_{n_{0}+1}^{-M_{0}} \tag{4.137}
\end{equation*}
$$

It is also convenient for $\sigma \in \Sigma, \omega \in \Omega$, to introduce the laws $P_{y, \omega}^{\sigma}, y \in \mathbb{R}^{d}$, of the canonical Markov chain on $\left(\mathbb{R}^{d}\right)^{\mathbf{N}}$, with transition kernel $R_{n_{0}^{\prime}, \sigma}^{*}$, cf. (4.7). We denote with $E_{y, \omega}^{\sigma}$ the
corresponding expectation and with $Z_{k}, k \geq 0$, the canonical process on $\left(\mathbb{R}^{d}\right)^{\mathbb{N}}$. So for instance for bounded measurable $f$ and $n \in\left[n_{0}^{\prime}, n_{0}\right], y \in \mathbb{R}^{d}$, in view of (4.11),

$$
\begin{align*}
& \widetilde{R}_{n, \sigma}^{*} f(y)= \\
& E_{y, \omega}^{\sigma}\left[\sum_{0 \leq m<k_{n}} \prod_{0 \leq k<m} \psi_{n, y}\left(Z_{k}\right)\left(1-\psi_{n, y}\left(Z_{m}\right)\right) f\left(Z_{m}\right)+\prod_{0 \leq k<k_{n}} \psi_{n, y}\left(Z_{k}\right) f\left(Z_{k_{n}}\right)\right], \tag{4.138}
\end{align*}
$$

with $k_{n}=\left(L_{n} / L_{n_{0}^{\prime}}\right)^{2}$.
Lemma 4.8. When $L_{0}$ is large, for $\sigma \in \Sigma$, $n_{0}^{\prime} \leq n \leq n_{0}, y \in\left\{d\left(\cdot, \operatorname{Supp} h_{n}\right) \leq 50 \sqrt{d} L_{n}\right\}$, $x \in L_{n} \mathbb{Z}^{d} \cap\left\{d\left(\cdot, \operatorname{Supp} h_{n}\right) \leq 20 \sqrt{d} L_{n}\right\}, \omega \in \widetilde{G}$,

$$
\begin{align*}
& P_{y, \omega}^{\sigma}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \geq 30 \widetilde{\ell}_{0} \widetilde{D}_{n}\right] \leq e^{-\kappa_{n_{0}}}  \tag{4.139}\\
& \left\|\chi_{n, x}\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right)\right\|_{n} \leq e^{-\kappa_{n_{0}}} \tag{4.140}
\end{align*}
$$

and

$$
\begin{equation*}
\left\|\chi_{n, x} S_{n, \sigma}^{*}\right\|_{n} \leq c L_{n}^{\beta} \tag{4.141}
\end{equation*}
$$

Proof. We begin with the proof of (4.139). The case $n=n_{0}^{\prime}$ is obvious since $k_{n}=1$, and the steps of $Z$. have length at most $\widetilde{D}_{n_{0}^{\prime}}, P_{y, \omega^{-}}^{\sigma}$ a.s. cf. (4.7). Since $\omega \in \widetilde{G}$, we can find a collection $w_{i} \in L_{n} \mathbb{Z}^{d}, 1 \leq i \leq \widetilde{\ell}_{0}$, with $B\left(w_{i}, 3 \widetilde{D}_{n}\right) \cap 5 \mathcal{T}_{n_{0}+1} \neq \emptyset$, such that

$$
\begin{equation*}
\widetilde{\mathcal{B}}_{n}(\omega) \supseteq\left(\left(5 \mathcal{T}_{n_{0}+1}\right) \cap L_{n} \mathbb{Z}^{d}\right) \backslash \bigcup_{1 \leq i \leq \widetilde{\ell}_{0}} B\left(w_{i}, 3 \widetilde{D}_{n}\right) \tag{4.142}
\end{equation*}
$$

Let us write $\sigma=\left(\sigma_{1}, \ldots, \sigma_{\widetilde{\ell}}\right)$, where $0 \leq \widetilde{\ell} \leq \widetilde{\ell}_{0}$, and introduce the open set

$$
U=\left(\bigcup_{1 \leq i \leq \widetilde{\ell}_{0}} B\left(w_{i}, 6 \widetilde{D}_{n}\right)\right) \cup\left(\bigcup_{1 \leq i \leq \widetilde{\ell}} B\left(\sigma_{i}, 6 \widetilde{D}_{n}\right)\right) .
$$

Since $P_{y, \omega^{-}}^{\sigma}$ a.s., $Z$. has steps of length at most $\widetilde{D}_{n_{0}^{\prime}}$, and $U$ is a union of at most $2 \widetilde{\ell}_{0}$ balls of radius $6 \widetilde{D}_{n}$, using a connectedness argument we see that $P_{y, \omega}^{\sigma}$-a.s., on the event $\bigcap_{0 \leq k \leq k_{n}}\left\{Z_{k} \in U\right\}$, one has $\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \leq 7 \times 2 \widetilde{\ell}_{0} \widetilde{D}_{n}$. Therefore $P_{y, \omega}^{\sigma}$ a.s., on the event in (4.139), $Z$. exits $U$ before times $k_{n}$. If we now define:

$$
\begin{equation*}
\tau=\inf \left\{k \geq 0 ; \inf _{z} d\left(Z_{k}, z\right) \geq 4 \widetilde{D}_{n}\right\},\left(z \text { runs over }\left\{w_{1}, \ldots, w_{\widetilde{\ell}_{0}}, \sigma_{1}, \ldots, \sigma_{\ell}\right\}\right) \tag{4.143}
\end{equation*}
$$

we see that the probability in (4.139) is smaller than:

$$
\begin{equation*}
E_{y, \omega}^{\sigma}\left[\tau<k_{n}, P_{Z_{\tau}, \omega}^{\sigma}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right|>\frac{\widetilde{D}_{n}}{2}\right]\right] \tag{4.144}
\end{equation*}
$$

where we have used the strong Markov property. With our choice of $y$, see also below (4.8), we see that $P_{y, \omega}^{\sigma}$-a.s., on $\left\{\tau<k_{n}\right\}, d_{\infty}\left(Z_{\tau},\left(5 \mathcal{T}_{n_{0}+1}\right)^{c}\right) \geq L_{n+1}^{2}-c L_{n}-\left(L_{n} / L_{n_{0}^{\prime}}\right)^{2} \widetilde{D}_{n_{0}^{\prime}} \geq$ $\widetilde{D}_{n}+2 \widetilde{D}_{n_{0}^{\prime}}$, when $L_{0}$ is large.

So in view of (4.6), with the notation (1.18), we obtain that $P_{y, \omega}^{\sigma}$-a.s., on $\left\{\tau<k_{n}\right\}$,

$$
P_{Z_{\tau}, \omega}^{\sigma}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right|>\frac{\widetilde{D}_{n}}{2}\right] \leq P_{Z_{\tau}, \omega}\left[X_{L_{n}^{2}}^{*}>\frac{\widetilde{D}_{n}}{2}\right]^{(2.2),(4.142)} \leq e^{-\kappa_{n}}
$$

Coming back to (4.144), we obtain (4.139).
We now prove (4.140). Once again the case $n=n_{0}^{\prime}$ is immediate since $\widetilde{R}_{n_{0}^{\prime}, \sigma}^{*}$ coincides with $R_{n_{0}^{\prime}, \sigma^{\prime}}^{*}$. We thus assume $n_{0}^{\prime}<n \leq n_{0}$, and choose $f$ with $|f|_{(n)} \leq 1, \omega \in \widetilde{G}$. With large $L_{0}$, we see that, cf. (4.9), (4.138), for $x$ as in (4.140), $y \in \mathbb{R}^{d}$,

$$
\begin{align*}
& \chi_{n, x}(y)\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right) f(y) \stackrel{\text { def }}{=} \chi_{n, x}(y) \Delta_{n} f(y), \text { with } \\
& \Delta_{n} f(y)=E_{y, \omega}^{\sigma}\left[f\left(Z_{k_{n}}\right)-\sum_{0 \leq m<k_{n}} \prod_{0 \leq k<m} \psi_{n, y}\left(Z_{k}\right)\left(1-\psi_{n, y}\left(Z_{m}\right)\right) f\left(Z_{m}\right)-\right.  \tag{4.145}\\
& \left.\prod_{0 \leq k<k_{n}} \psi_{n, y}\left(Z_{k}\right) f\left(Z_{k_{n}}\right)\right],
\end{align*}
$$

and hence by the choice of $\psi_{n, y}$, cf. (4.10),

$$
\begin{equation*}
\left|\chi_{n, x}(y)\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right) f(y)\right| \leq 2 \chi_{n, x}(y) P_{y, \omega}^{\sigma}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \geq D_{n}^{*}\right] \stackrel{(4.139)}{\leq} e^{-\kappa_{n_{0}}} \tag{4.146}
\end{equation*}
$$

Then for $y, y^{\prime}$ in $\left\{d\left(\cdot, \operatorname{Supp} \chi_{n, x}\right) \leq L_{n}\right\}$, we see that when $\left|y-y^{\prime}\right| \geq e^{-\kappa_{n_{0}}}$,

$$
\begin{equation*}
\left|\chi_{n, x}(y) \Delta_{n} f(y)-\chi_{n, x}\left(y^{\prime}\right) \Delta_{n} f\left(y^{\prime}\right)\right| \leq e^{-\kappa_{n_{0}}} \leq\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} e^{-\kappa_{n_{0}}} \tag{4.147}
\end{equation*}
$$

We thus consider $y, y^{\prime}$ in $\left\{d\left(\cdot, \operatorname{Supp} \chi_{n, x}\right) \leq L_{n}\right\}$, with

$$
\begin{equation*}
\left|y-y^{\prime}\right| \leq e^{-\kappa_{n_{0}}} \tag{4.148}
\end{equation*}
$$

and write in analogy with (2.51):

$$
\begin{equation*}
\left|\Delta_{n} f(y)-\Delta_{n} f\left(y^{\prime}\right)\right| \leq a_{1}+a_{2}, \text { where } \tag{4.149}
\end{equation*}
$$

$$
\begin{aligned}
& a_{1}=\mid E_{y^{\prime}, \omega}^{\sigma}\left[\sum_{0 \leq m<k_{n}} \prod_{0 \leq k<m} \psi_{n, y^{\prime}}\left(Z_{k}\right)\left(1-\psi_{n, y^{\prime}}\left(Z_{m}\right)\right) f\left(Z_{m}\right)+\prod_{0 \leq k<k_{n}} \psi_{n, y^{\prime}}\left(Z_{k}\right) f\left(Z_{k_{n}}\right)\right. \\
& \left.-\sum_{0 \leq m<k_{n}} \prod_{0 \leq k<m} \psi_{n, y}\left(Z_{k}\right)\left(1-\psi_{n, y}\left(Z_{m}\right)\right) f\left(Z_{m}\right)-\prod_{0 \leq k<k_{n}} \psi_{n, y}\left(Z_{k}\right) f\left(Z_{k_{n}}\right)\right] \mid
\end{aligned}
$$

and with hopefully obvious notations

$$
a_{2}=\left|\left(E_{y, \omega}^{\sigma}-E_{y^{\prime}, \omega}^{\sigma}\right)\left[f\left(Z_{k_{n}}\right)-\sum_{0 \leq m<k_{n}} \psi_{n, y}\left(Z_{k}\right)\left(1-\psi_{n, y}\left(Z_{m}\right)\right) f\left(Z_{m}\right)-\prod_{0 \leq k<k_{n}} \psi_{n, y}\left(Z_{k}\right) f\left(Z_{k_{n}}\right)\right]\right| .
$$

In view of (4.10), $\left|\psi_{n, y}(\cdot)-\psi_{n, y^{\prime}}(\cdot)\right| \leq\left|y-y^{\prime}\right|$, and we see that with (4.148) and (1.13),

$$
\begin{equation*}
a_{1} \leq\left(k_{n}^{2}+k_{n}\right)\left|y-y^{\prime}\right| \leq\left(k_{n}^{2}+k_{n}\right) e^{-\kappa_{n_{0}}}\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} \leq e^{-\kappa_{n_{0}}}\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} \tag{4.150}
\end{equation*}
$$

using (4.148), and (1.13). Then using the fact that, cf. (4.6), (4.7),

$$
R_{n_{0}^{\prime}, \sigma}^{*}=\left(1-g_{\sigma}\right) \widetilde{R}_{n_{0}^{\prime}}^{0}+g_{\sigma} \widetilde{R}_{n_{0}^{\prime}}
$$

we can write

$$
\begin{align*}
R_{n_{0}^{\prime}, \sigma}^{*}=A+B, \text { with } A & =\left(1-g_{\sigma}\right) R_{n_{0}^{\prime}}^{0}+g_{\sigma} R_{n_{0}^{\prime}}, \text { and } \\
B & =\left(1-g_{\sigma}\right)\left(\widetilde{R}_{n_{0}^{\prime}}^{0}-R_{n_{0}^{\prime}}^{0}\right)+g_{\sigma}\left(\widetilde{R}_{n_{0}^{\prime}}-R_{n_{0}^{\prime}}\right) . \tag{4.151}
\end{align*}
$$

With (1.60), (1.29), (2.46), we find

$$
\begin{align*}
& \text { i) }\|A\|_{L^{\infty} \rightarrow(n)} \leq\left(\frac{L_{n}}{L_{n_{0}^{\prime}}}\right)^{\beta}\|A\|_{L^{\infty} \rightarrow\left(n_{0}^{\prime}\right)} \leq c L_{n}^{\beta} \text {, and }  \tag{4.152}\\
& \text { ii) }\|B\|_{n} \leq\left(\frac{L_{n}}{L_{n_{0}^{\prime}}}\right)^{\beta}\|B\|_{n_{0}^{\prime}} \leq e^{-\kappa_{n_{0}}} \text {. }
\end{align*}
$$

Denoting with $g(\cdot)$ the function $\chi_{L_{n}}(\cdot-y)$, cf. (1.37), we have

$$
\begin{equation*}
a_{2}=\left|R_{n_{0}^{\prime}, \sigma}^{*}\left(g E_{\cdot, \omega}^{\sigma}[H]\right)(y)-R_{n_{0}^{\prime}, \sigma}^{*}\left(g E_{\cdot, \omega}^{\sigma}[H]\right)\left(y^{\prime}\right)\right| \tag{4.153}
\end{equation*}
$$

where $|H| \leq 21_{\left\{\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \geq D_{n}^{*} / 2\right\}}$, and

$$
E_{z, \omega}^{\sigma}[H]=\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}-1} f(z)-\sum_{0 \leq m<k_{n}-1}\left(\psi_{n, y} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{m}\left(1-\psi_{n, y}\right) f(z)-\left(\psi_{n, y} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}-1} f(z)
$$

Using (4.151) in (4.153), as well as (4.152) i), we thus find

$$
\begin{align*}
a_{2} & \leq\left|\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} c L_{n}^{\beta} \sup _{z \in B\left(y, 2 L_{n}\right)} P_{z, \omega}^{\sigma}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \geq \frac{D_{n}^{*}}{2}\right]+a_{2}^{\prime} \\
& \quad(4.139)  \tag{4.154}\\
\leq & \left.\frac{y-y^{\prime}}{L_{n}}\right|^{\beta} e^{-\kappa_{n}}+a_{2}^{\prime}, \text { where } \\
a_{2}^{\prime} & =\left|B\left(g E_{\cdot, \omega}[H]\right)(y)-B\left(g E_{\cdot, \omega}[H]\right)\left(y^{\prime}\right)\right| .
\end{align*}
$$

In view of (4.152) ii), (4.147) - (4.150), the claim (4.140) will follow once we show that

$$
\begin{equation*}
\left|g E_{\cdot, \omega}[H]\right|_{(n)} \leq c L_{n}^{\beta} k_{n} \tag{4.155}
\end{equation*}
$$

To this end observe that for $m \geq 1$, with (4.151), using perturbation expansion

$$
\begin{equation*}
\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{m}=B^{m}+\sum_{0 \leq m^{\prime}<m} B^{m^{\prime}} A\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{m-m^{\prime}-1}, \tag{4.156}
\end{equation*}
$$

so that with (4.152)

$$
\left|\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}-1} f\right|_{(n)} \leq\|B\|_{n}^{k_{n}-1}+\sum_{0 \leq m^{\prime}<k_{n}-1}\|B\|_{n}^{m^{\prime}} c L_{n}^{\beta} \leq c L_{n}^{\beta} .
$$

Analogously, we see that with $0 \leq m<k_{n}-1$,

$$
\begin{aligned}
\left|\left(\psi_{n, y} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{m}\left(1-\psi_{n, y}\right) f\right|_{(n)} & \leq\left\|\psi_{n, y} B\right\|_{n}^{m}\left|1-\psi_{n, y}\right|_{(n)}+\sum_{0 \leq m^{\prime}<m}\left\|\psi_{n, y} B\right\|_{n}^{m^{\prime}} c L_{n}^{\beta} \\
& \begin{array}{rl}
(4.10),(4.152) \\
& \leq \\
\mid\left(\psi_{n, y} R_{n_{0}^{\prime}, \sigma}^{*}\right)^{\beta}, \text { and } \\
k_{n}-1 & \left.f\right|_{(n)}
\end{array} \leq c L_{n}^{\beta} .
\end{aligned}
$$

The claim (4.155) follows, and this finishes the proof of (4.140).
Let us finally prove (4.141). For large $L_{0}, \sigma \in \Sigma, n_{0}^{\prime} \leq n \leq n_{0}, \omega \in \widetilde{G}$, and $x$ as in (4.141), as a result of (4.8), (4.9):

$$
\chi_{n, x} S_{n, \sigma}^{*}=\chi_{n, x}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}}-\chi_{n, x} R_{n}^{0} .
$$

Using (4.156) and (4.152), the claim (4.141) immediately follows.
Keeping in mind the expansion (4.15), it is convenient to modify (4.16), and introduce for $\sigma \in \Sigma, n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}$ the operator

$$
\begin{equation*}
\mathcal{L}_{\sigma, n, v}=\sum_{0 \leq k<\ell_{n}^{2}} \chi_{n+1, v}\left(R_{n}^{0}\right)^{k} h_{n} S_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} . \tag{4.157}
\end{equation*}
$$

As an application of the previous lemma we have
Lemma 4.9. When $L_{0}$ is large, for $\sigma \in \Sigma, n_{0}^{\prime} \leq n \leq n_{0}, v \in L_{n+1} \mathbb{Z}^{d}, \omega \in \widetilde{G}$

$$
\begin{equation*}
\left\|\mathcal{L}_{\sigma, n, v}-\widetilde{\mathcal{L}}_{\sigma, n, v}\right\|_{n+1} \leq e^{-\kappa_{n_{0}}} . \tag{4.158}
\end{equation*}
$$

Proof. We write, (recall that $\left.h_{n, v}(\cdot)=\chi_{D_{n+1}}(\cdot-v) h_{n}(\cdot)\right)$,

$$
\begin{align*}
& \mathcal{L}_{\sigma, n, v}-\widetilde{\mathcal{L}}_{\sigma, n, v}=\mathcal{L}^{1}+\mathcal{L}^{2}+\mathcal{L}^{3}, \text { with } \\
& \mathcal{L}^{1}=\sum_{0 \leq k<\ell_{n}^{2}} \chi_{n+1, v}\left(R_{n}^{0}\right)^{k}\left(h_{n}-h_{n, v}\right) S_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1} \\
& \mathcal{L}^{2}=\sum_{0 \leq k<\ell_{n}^{2}} \chi_{n+1, v}\left(R_{n}^{0}\right)^{k} h_{n, v}\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right)\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1}  \tag{4.159}\\
& \mathcal{L}^{3}=\sum_{0 \leq k<\ell_{n}^{2}} \chi_{n+1, v}\left(R_{n}^{0}\right)^{k} h_{n, v} \widetilde{S}_{n, \sigma}^{*}\left(R_{n}^{0}\right)^{\ell_{n}^{2}-k-1}\left(1-\widetilde{\chi}_{n+1, v}\right) .
\end{align*}
$$

Keeping in mind (4.140), (4.141), together with (1.55), (1.56), (1.49), (1.29), we see that

$$
\left\|\mathcal{L}^{1}\right\|_{n} \leq \ell_{n}^{2} c e^{-\kappa_{n_{0}}} c L_{n}^{\beta} \leq e^{-\kappa_{n_{0}}}, \quad\left\|\mathcal{L}^{2}\right\|_{n} \leq \ell_{n}^{2} c e^{-\kappa_{n_{0}}} \leq e^{-\kappa_{n_{0}}}
$$

Noting that $h_{n, v} \widetilde{S}_{n, \sigma}^{*} g=-h_{n, v} R_{n}^{0} g$, when $g$ is supported in $B\left(v, 3 D_{n+1}\right)^{c}$, with $L_{0}$ large, we also find

$$
\left\|\mathcal{L}^{3}\right\|_{n} \leq \ell_{n}^{2} c L_{n}^{\beta} e^{-\kappa_{n_{0}}} \leq e^{-\kappa_{n_{0}}} .
$$

Since we also have $\left\|\mathcal{L}^{i}\right\|_{n+1} \leq \ell_{n}^{\beta}\left\|\mathcal{L}^{i}\right\|_{n}$, for $i=1,2,3$, the claim (4.158) follows.
Proposition 4.10. When $L_{0}$ is large, for $n_{0}^{\prime} \leq n \leq n_{0}$, (4.18) is satisfied.
Proof. We use induction over $n \in\left[n_{0}^{\prime}, n_{0}\right]$. First observe that with the notation (4.5) and in analogy with (4.3)

$$
\mathbb{P}\left[G_{\emptyset}\right] \geq 1-c\left(\frac{L_{n_{0}+1}^{2}}{L_{n_{0}^{\prime}}}\right)^{d} L_{n_{0}^{\prime}}^{-M_{0}} \geq 1-c L_{n_{0}+1}^{2 d-M_{0}(1+a)^{-\left(m_{0}+2\right)}} \stackrel{(1.46)}{\geq} 1-c L_{n_{0}+1}^{-98 d} .
$$

Hence with (4.137), we find for large $L_{0}$

$$
\begin{equation*}
\mathbb{P}\left[G_{\emptyset, n_{0}^{\prime}}\right] \geq 1-L_{n_{0}+1}^{-97 d}, \text { with } G_{\emptyset, n_{0}^{\prime}} \stackrel{\text { def }}{=} G_{\emptyset} \cap \widetilde{G} . \tag{4.160}
\end{equation*}
$$

We introduce the notation

$$
\begin{equation*}
\mathcal{S}_{n}=L_{n} \mathbb{Z}^{d} \cap\left\{d\left(\cdot, \operatorname{Supp} h_{n}\right) \leq 20 \sqrt{d} L_{n}\right\}, \text { for } n_{0}^{\prime} \leq n<n_{0} \tag{4.161}
\end{equation*}
$$

Note for later use that with the notation (4.19), for $n_{0}^{\prime} \leq n<n_{0}$,

$$
\begin{equation*}
\mathcal{S}_{n+1} \subseteq\left\{v \in L_{n+1} \mathbb{Z}^{d} ; \mathcal{S}_{n, v} \neq \emptyset\right\}=\left\{v \in L_{n+1} \mathbb{Z}^{d} ; h_{n, v} \not \equiv 0\right\} . \tag{4.162}
\end{equation*}
$$

Further when $L_{0}$ is large, for all $\omega \in G_{\emptyset, n_{0}^{\prime}}, x \in \mathcal{S}_{n_{0}^{\prime}}$, with (4.7)

$$
\begin{equation*}
\left\|\chi_{n_{0}^{\prime}, x} \widetilde{S}_{n_{0}^{\prime}, \forall}^{*}\right\|_{n_{0}^{\prime}}=\left\|\chi_{n_{0}^{\prime}, x} \widetilde{S}_{n_{0}^{\prime}}\right\|_{n_{0}^{\prime}} \stackrel{(2.2)}{\leq} \nu_{n_{0}^{\prime}} \tag{4.163}
\end{equation*}
$$

and for all $y \in\left[0, L_{n_{0}^{\prime}}\right]^{d}$, using (2.2), (2.4)

$$
\begin{equation*}
\left|\frac{\widetilde{d}_{n_{0}^{\prime}, \boldsymbol{\varphi}}^{*}}{L_{n_{0}^{\prime}}^{*}}(y, \omega)\right|\left(=\left|\frac{\widetilde{d}_{n_{0}^{\prime}}}{L_{n_{0}^{\prime}}}(y, \omega)\right|\right) \leq \nu_{n_{0}^{\prime}},\left|\frac{\widetilde{\gamma}_{n_{0}^{\prime}, \boldsymbol{v}}^{*}}{L_{n_{0}^{\prime}}^{2}}(y, \omega)\right|\left(=\left|\frac{\widetilde{\gamma}_{n_{0}^{\prime}}}{L_{n_{0}^{\prime}}^{2}}(y, \omega)\right|\right) \leq \nu_{n_{0}^{\prime}} . \tag{4.164}
\end{equation*}
$$

Let us assume that for $n_{1}$ with $n_{0}^{\prime} \leq n_{1}<n_{0}$, we have a decreasing sequence of events $G_{\emptyset, n}, n_{0}^{\prime} \leq n \leq n_{1}$, such that for $n_{0}^{\prime} \leq n<n_{1}$

$$
\begin{equation*}
\mathbb{P}\left[G_{\emptyset, n} \backslash G_{\emptyset, n+1}\right] \leq e^{-\kappa_{n_{0}}}, \tag{4.165}
\end{equation*}
$$

and for $\omega \in G_{\emptyset, n}, x \in \mathcal{S}_{n}$, (4.163), (4.164) hold with $n$ in place of $n_{0}^{\prime}$ (the expressions in parenthesis in (4.164) being now disregarded). With (4.160), we see that (4.18) is satisfied with $n=n_{1}$, and with (4.20) of Proposition 4.1, where we have set $G_{\emptyset, n_{1}, v} \equiv G_{\emptyset, n_{1}}$, we obtain since $\mathcal{S}_{n_{1}, v} \subseteq \mathcal{S}_{n_{1}}$, for all $v \in L_{n+1} \mathbb{Z}^{d}$,

$$
\begin{align*}
& \mathbb{P}\left[G_{\emptyset, n_{1}} \cap\left\{\sup _{v \in L_{n_{1}+1} \mathbb{Z}^{d}: \mathcal{S}_{n_{1}, v} \neq \emptyset}\left\|\widetilde{\mathcal{L}}_{\emptyset, n_{1}, v}\right\|_{n_{1}+1}>\frac{\kappa_{n_{1}} \nu_{n_{1}}}{\ell_{n_{1}}^{\beta / 3}}\right\}\right] \leq \\
& c\left(\frac{L_{n_{0}+1}^{2}}{L_{n_{1}+1}}\right)^{d} e^{-\kappa_{n_{0}}} \leq e^{-\kappa_{n_{0}}} . \tag{4.166}
\end{align*}
$$

We then define

$$
\begin{equation*}
G_{\emptyset, n_{1}+1}=G_{\emptyset, n_{1}} \cap\left\{\sup _{v \in L_{n_{1}+1} \mathbb{Z}^{d}: \mathcal{S}_{n_{1}, v} \neq \emptyset}\left\|\widetilde{\mathcal{L}}_{\emptyset, n_{1}, v}\right\|_{n_{1}+1} \leq \frac{\kappa_{n_{1}} \nu_{n_{1}}}{\ell_{n_{1}}^{\beta / 3}}\right\}, \tag{4.167}
\end{equation*}
$$

and note from the above that (4.165) is true for $n=n_{1}$. Then with Lemma 4.9, since $G_{\emptyset, n_{1}+1} \subseteq \widetilde{G}$, we have for $\omega \in G_{\emptyset, n_{1}+1}$

$$
\begin{equation*}
\sup _{v \in L_{n_{1}+1} \mathbb{Z}^{d}: \mathcal{S}_{n_{1}, v} \neq \emptyset}\left\|\mathcal{L}_{\emptyset, n_{1}, v}\right\|_{n_{1}+1} \leq 2 \frac{\kappa_{n_{1}} \nu_{n_{1}}}{\ell_{n_{1}}^{\beta / 3}} \tag{4.168}
\end{equation*}
$$

Coming back to (4.15), we see that for $\omega \in G_{\emptyset, n_{1}+1}, v \in \mathcal{S}_{n_{1}+1}$ :

$$
\begin{align*}
& \left\|\chi_{n_{1}+1, v} S_{n_{1}+1, \emptyset}^{*}\right\|_{n_{1}+1} \leq\left\|\mathcal{L}_{\emptyset, n_{1}, v}\right\|_{n_{1}+1}+ \\
& \left\|\sum_{\substack{k_{0}+\cdots+k_{m}+m=\ell_{n_{1}}^{2} \\
k_{i} \geq 0, m \geq}} \chi_{n_{1}+1, v}\left(R_{n_{1}}^{0}\right)^{k_{0}} h_{n_{1}} S_{n_{1}, \emptyset}^{*}\left(R_{n_{1}}^{0}\right)^{k_{1}} \ldots h_{n_{1}} S_{n_{1}, \emptyset}^{*}\left(R_{n_{1}}^{0}\right)^{k_{m}}\right\|_{n_{1}+1}+  \tag{4.169}\\
& c\left\|P_{\alpha_{n_{1}} L_{n_{1}+1}^{2}}-P_{\alpha_{n_{1}+1} L_{n_{1}+1}^{2}}\right\|_{n_{1}+1} \stackrel{\text { def }}{=} a_{1}+a_{2}+a_{3} .
\end{align*}
$$

With (4.162), from (4.168) we find

$$
\begin{equation*}
a_{1} \leq \kappa_{n_{1}} \nu_{n_{1}} \ell_{n_{1}}^{-\beta / 3} \tag{4.170}
\end{equation*}
$$

Then with (4.163), with $n_{1}$ in place of $n_{0}^{\prime}$, (A.3) of the Appendix, and (4.140), we see that for $\omega \in G_{\emptyset, n_{1}+1} \subseteq G_{\emptyset, n_{1}}$ :

$$
\begin{equation*}
\left\|h_{n_{1}} S_{n_{1}, \emptyset}^{*}\right\|_{n_{1}} \leq\left\|h_{n_{1}} \widetilde{S}_{n_{1}, \emptyset}^{*}\right\|_{n_{1}}+e^{-\kappa_{n_{0}}} \leq c \nu_{n_{1}}+e^{-\kappa_{n_{0}}} \leq c_{3} \nu_{n_{1}} \tag{4.171}
\end{equation*}
$$

As a result with the help of (1.55) and the fact that $\|\cdot\|_{n+1} \leq \ell_{n}^{\beta}\|\cdot\|_{n}$, we obtain

$$
\begin{align*}
& a_{2} \leq c \ell_{n_{1}}^{\beta} \sum_{\substack{k_{0}+\cdots+k_{m}+m=\ell_{n_{1}}^{2} \\
k_{i} \geq 0, m \geq 2}}\left(c_{3} \nu_{n_{1}}\right)^{m}=c \ell_{n_{1}}^{\beta}\left[\left(1+c_{3} \nu_{n_{1}}\right)^{\ell_{n_{1}}^{2}}-1-\ell_{n_{1}}^{2} c_{3} \nu_{n_{1}}\right]  \tag{4.172}\\
& \leq c \ell_{n_{1}}^{\beta+4} \nu_{n_{1}}^{2} \exp \left\{c \nu_{n_{1}} \ell_{n_{1}}^{2}\right\} \stackrel{(1.14),(1.40),(4.17)}{\leq} c L_{n_{1}}^{5 a} \nu_{n_{1}}^{2},
\end{align*}
$$

where we used the inequalities $(1+u)^{\ell} \leq e^{u \ell}$ and $e^{v}-1-v \leq v^{2} e^{v}$, for $\ell, u, v$ positive numbers. To bound $a_{3}$, we use the heat equation satisfied by the Brownian semigroup, which implies that for $f$ with $|f|_{\left(n_{1}+1\right)} \leq 1$,

$$
\begin{align*}
& \left|P_{\alpha_{n_{1}} L_{n_{1}+1}^{2}} f-P_{\alpha_{n_{1}+1} L_{n_{1}+1}^{2}} f\right|_{\left(n_{1}+1\right)}=\left|\int_{\alpha_{n_{1}+1} L_{n_{1}+1}^{2}}^{\alpha_{n_{1}} L_{n_{1}+1}^{2}} \frac{1}{2} \Delta P_{s} f d s\right|_{\left(n_{1}+1\right)} \\
& =\left|\int_{\alpha_{n_{1}} L_{n_{1}+1}^{2}}^{\alpha_{n_{1}+1} L_{n_{1}+1}^{2}} \frac{1}{2} P_{s / 2} \Delta P_{s / 2} f d s\right|_{\left(n_{1}+1\right)} \stackrel{(1.56),(1.49) i)}{\leq} c\left|\alpha_{n_{1}+1}-\alpha_{n_{1}}\right|  \tag{4.173}\\
& \stackrel{(1.49) i i)}{\leq} c L_{n_{1}}^{-\frac{19}{10} \delta} \stackrel{(1.14)}{\leq} L_{n_{1}+1}^{-\frac{18}{10} \delta} .
\end{align*}
$$

We have thus shown that when $L_{0}$ is large

$$
\begin{equation*}
a_{3} \leq c L_{n_{1}+1}^{-\frac{18}{10} \delta} . \tag{4.174}
\end{equation*}
$$

Collecting (4.170), (4.172), (4.174), we see that when $L_{0}$ is large, for $\omega \in G_{\emptyset, n_{1}+1}, v \in$ $\mathcal{S}_{n_{1}+1}$ :

$$
\begin{equation*}
\left\|\chi_{n_{1}+1, v} S_{n_{1}+1, \emptyset}^{*}\right\|_{n_{1}+1} \leq c\left(\kappa_{n_{1}} \nu_{n_{1}} \ell_{n_{1}}^{-\beta / 3}+L_{n_{1}}^{5 a} \nu_{n_{1}}^{2}+L_{n_{1}+1}^{-\frac{18}{10} \delta}\right) \tag{4.175}
\end{equation*}
$$

and thank to (4.140), a similar inequality is satisfied by $\chi_{n_{1}+1, v} \widetilde{S}_{n_{1}+1, \emptyset}^{*}$. If we now choose $v=0$, analogous controls as in the derivation of (2.4), using (4.14), and (1.49) i) with $n=n_{1}+1 \leq n_{0}$, and the remark below (4.11), show that

$$
\begin{align*}
& \sup _{y \in\left[0, L_{n_{1}+1}\right]^{d}}\left(\left|\frac{\widetilde{d}_{n_{1}+1}^{*}}{L_{n_{1}+1}}(y, \omega)\right|+\left|\frac{\widetilde{\gamma}_{n_{1}+1}^{*}}{L_{n_{1}+1}^{2}}(y, \omega)\right|\right) \leq  \tag{4.176}\\
& \kappa_{n_{1}+1}\left(\kappa_{n_{1}} \nu_{n_{1}} \ell_{n_{1}}^{-\beta / 3}+L_{n_{1}}^{5 a} \nu_{n_{1}}^{2}+L_{n_{1}+1}^{-\frac{18}{10} \delta}\right) \leq \nu_{n_{1}+1}
\end{align*}
$$

using (1.14), (1.40), (4.17) in the last step.
We thus see that (4.163), (4.164) are satisfied for $\omega \in G_{\emptyset, n_{1}+1}, v \in \mathcal{S}_{n_{1}+1}$, with $n_{1}+1$, in place of $n_{0}^{\prime}$. This completes the induction step, and with (4.160), this is more than enough to prove the claim of Proposition 4.10.

We are now ready to state and prove the main result of this section. We recall the notations introduced in (4.4), (4.5), (4.136).
Proposition 4.11. When $L_{0}$ is large, for each $\sigma \in \Sigma$ there is an event $G_{\sigma, n_{0}+1} \subseteq G_{\sigma} \cap \widetilde{G}$, such that:

$$
\begin{align*}
& \sup _{\sigma \in \Sigma} \mathbb{P}\left[\left(G_{\sigma} \cap \widetilde{G}\right) \backslash G_{\sigma, n_{0}+1}\right] \leq e^{-\kappa_{n_{0}}}  \tag{4.177}\\
& \mathbb{P}\left[\left(\bigcup_{\sigma \in \Sigma} G_{\sigma, n_{0}+1}\right)^{c}\right] \leq \frac{1}{20} L_{n_{0}+1}^{-M_{0}} \tag{4.178}
\end{align*}
$$

and on $G_{\sigma, n_{0}+1}$, for all $n_{0}^{\prime} \leq n \leq n_{0}$, (cf. (4.17), (4.162) for the notation),

$$
\begin{equation*}
\sup _{x \in \mathcal{S}_{n}}\left(\left\|\chi_{n, x} S_{n, \sigma}^{*}\right\|_{n} \vee\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}\right) \leq \nu_{n} \tag{4.179}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1, \sigma}^{*}-\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}}\right)\right\|_{n_{0}+1} \leq \nu_{n_{0}+1} . \tag{4.180}
\end{equation*}
$$

Proof. The argument is similar to the proof of Proposition 4.10. We define for $\sigma \in \Sigma$,

$$
\begin{equation*}
G_{\sigma, n_{0}^{\prime}}=G_{\sigma} \cap \widetilde{G} \tag{4.181}
\end{equation*}
$$

(this is consistent with (4.160), when $\sigma=\emptyset$ ). We then observe with (4.7), (4.11), that when $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in G_{\sigma, n_{0}^{\prime}}, v \in \mathcal{S}_{n_{0}^{\prime}}$ :

$$
\begin{align*}
& \left\|\chi_{n_{0}^{\prime}, x} \widetilde{S}_{n_{0}^{\prime}, \sigma}^{*}\right\|_{n_{0}^{\prime}}=\left\|\chi_{n_{0}^{\prime}, x} S_{n_{0}^{\prime}, \sigma}^{*}\right\|_{n_{0}}=\left\|\chi_{n_{0}^{\prime}, x}\left(g_{\sigma} \widetilde{S}_{n_{0}^{\prime}}+\left(1-g_{\sigma}\right)\left(\widetilde{R}_{n_{0}^{\prime}}^{0}-R_{n_{0}^{\prime}}^{0}\right)\right)\right\|_{n_{0}^{\prime}} \\
& \stackrel{(4.6),(2.2),(2.46)}{\leq} c\left(L_{n_{0}^{\prime}}^{-\delta}+e^{-\kappa_{n_{0}}}\right) \stackrel{(4.17)}{\leq} \nu_{n_{0}^{\prime}} . \tag{4.182}
\end{align*}
$$

Let us now assume that for $n_{1}$ with $n_{0}^{\prime} \leq n_{1}<n_{0}$, and $\sigma \in \Sigma$, we have a decreasing sequence of events, $n_{0}^{\prime} \leq n \leq n_{0}$, such that

$$
\begin{equation*}
\sup _{\sigma \in \Sigma} \mathbb{P}\left[G_{\sigma, n} \backslash G_{\sigma, n+1}\right] \leq e^{-\kappa_{n_{0}}}, \text { for } n_{0}^{\prime} \leq n<n_{1} \tag{4.183}
\end{equation*}
$$

and such that on $G_{\sigma, n}$ :

$$
\begin{equation*}
\sup _{x \in \mathcal{S}_{n}}\left(\left\|\chi_{n, x} S_{n, \sigma}^{*}\right\|_{n} \vee\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n}\right) \leq \nu_{n} \tag{4.184}
\end{equation*}
$$

Then with Proposition 4.1, for $\sigma \in \Sigma$,

$$
\begin{align*}
& \mathbb{P}\left[G_{\sigma, n_{1}} \cap\left\{\sup _{v \in L_{n_{1}+1} \mathbb{Z}^{d}: \mathcal{S}_{n_{1}, v} \neq \emptyset}\left\|\widetilde{\mathcal{L}}_{\sigma, n_{1}, v}\right\|_{n_{1}+1}>\frac{\kappa_{n} \nu_{n_{1}}}{\ell_{n_{1}}^{\beta / 3}}\right\}\right] \leq  \tag{4.185}\\
& c\left(\frac{L_{n_{0}+1}^{2}}{L_{n_{1}+1}}\right)^{d} e^{-\kappa_{n_{0}}} \leq e^{-\kappa_{n_{0}}} .
\end{align*}
$$

We then define for $\sigma \in \Sigma$, (this is consistent with (4.167)):

$$
\begin{equation*}
G_{\sigma, n_{1}+1}=G_{\sigma, n_{1}} \cap\left\{\sup _{v \in L_{n_{1}+1} \mathbb{Z}^{d}: \mathcal{S}_{n_{1}, v} \neq \emptyset}\left\|\widetilde{\mathcal{L}}_{\sigma, n_{1}, v}\right\|_{n_{1}+1} \leq \frac{\kappa_{n} \nu_{n_{1}}}{\ell_{n_{1}}^{\beta / 3}}\right\} \tag{4.186}
\end{equation*}
$$

and see that (4.184) holds with $n_{1}+1$ in place of $n_{1}$. Moreover in a parallel fashion to (4.169), for $\sigma \in \Sigma, \omega \in G_{\sigma, n_{1}+1}, v \in \mathcal{S}_{n_{1}+1}$,

$$
\begin{equation*}
\left\|\chi_{n_{1}+1, v} S_{n_{1}+1, \sigma}^{*}\right\|_{n_{1}+1} \leq a_{1}+a_{2}+a_{3}, \tag{4.187}
\end{equation*}
$$

where $a_{i}, 1 \leq i \leq 3$, are just as in (4.169), with $\sigma$ replacing $\emptyset$ in the expressions entering $a_{1}, a_{2}$. The same reasoning (4.170) - (4.174) shows that when $L_{0}$ is large, for $\sigma \in \Sigma$, $\omega \in G_{\sigma, n_{1}+1}$, and $v \in \mathcal{S}_{n_{1}+1}$ :

$$
\begin{equation*}
\left\|\chi_{n_{1}+1, v} S_{n_{1}+1, \sigma}^{*}\right\|_{n_{1}+1} \leq c\left(\kappa_{n_{1}} \kappa_{n_{1}} \ell_{n_{1}}^{-\beta / 3}+L_{n_{1}}^{5 a} \nu_{n_{1}}^{2}+L_{n_{1}+1}^{-\frac{18}{10} \delta}\right) \tag{4.188}
\end{equation*}
$$

and that with (4.140) a similar inequality holds for $\chi_{n_{1}+1, v} \widetilde{S}_{n_{1}+1, \sigma}^{*}$. This implies that (4.184) is true for $n=n_{1}+1$. This proves by induction (4.183) for $n_{0}^{\prime} \leq n<n_{0}$ and (4.184) for $n_{0}^{\prime} \leq n \leq n_{0}$. We can then define for $\sigma \in \Sigma, G_{\sigma, n_{0}+1}$ via (4.186) with $n_{0}$ in place of $n_{1}$. We then obtain (4.177), (4.180) by writing the analogue of (4.15) for $R_{n_{0}+1}^{*}-\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{0}}$, i.e. without the bottom line of (4.15), (incidentally we recall that (1.50) remains to be proved, cf. Proposition 5.7 below). The claim (4.178) is now a straightforward consequence of (4.5), (4.137), (4.177). This concludes the proof of Proposition 4.11.

## 5 Repairing defects

We conclude the proof of Theorem 1.1 in this section. The main remaining task is to propagate the part of (1.47) concerning Hölder-norm controls at level $n_{0}+1$. In Section 4 we have performed surgery on the environment and removed defects occurring at level $n_{0}^{\prime}=n_{0}-m_{0}-1$. We have shown that the kernels $R_{n, \sigma}^{*}, n_{0}^{\prime} \leq n \leq n_{0}+1, \sigma \in \Sigma$, cf. (4.7), (4.8), describing the evolution at level $n$ "after surgery", were typically well-behaved for Hölder-norms, when $\omega \in G_{\sigma, n_{0}+1}$, and that the complement of $\bigcup_{\sigma \in \Sigma} G_{\sigma, n_{0}+1}$, was "negligible" for our purpose, cf. Proposition 4.11. We now have to show that on "most" of $G_{\sigma, n_{0}+1}, R_{n_{0}+1, \sigma}^{*}$ and $R_{n_{0}+1}$, the true object of our interest, are close in the Höldernorm sense. To this end we will in essence use the smoothing effect of the kernels "after surgery" to repair defects, as well as (1.48) to prevent any trapping effect of the defects. The main step comes with Proposition 5.1. We will also prove (1.50), cf. Proposition 5.7, and thereby complete the proof of Theorem 1.1.

We first introduce some additional notations. We recall that $Z_{k}, k \geq 0$, denotes the canonical process on $\left(\mathbb{R}^{d}\right)^{\mathbb{N}}$, and that the laws $P_{y, \omega}^{\sigma}$, for $\sigma \in \Sigma, \omega \in \Omega, y \in \mathbb{R}^{d}$, with corresponding expectation $E_{y, \omega}^{\sigma}$, have been defined above (4.138). We let $P_{y, \omega}^{e}$ stand for the canonical law on $\left(\mathbb{R}^{d}\right)^{\mathbb{N}}$ of the Markov chain starting at $y \in \mathbb{R}^{d}$, with transition kernel $R_{n_{0}^{\prime}}$. It describes the diffusion in the environment (whence the superscript $e$ ) $\omega \in \Omega$, viewed at times $k L_{n_{0}^{\prime}}^{2}, k \geq 0$, originating from $y$. We let $E_{y, \omega}^{e}$ stand for the corresponding expectation. When no confusion with (1.8) arises, we use the notation

$$
\begin{equation*}
H_{C}=\inf \left\{k \geq 0, Z_{k} \in C\right\}, T_{C}=\inf \left\{k \geq 0, Z_{k} \notin C\right\} \tag{5.1}
\end{equation*}
$$

Likewise we still denote with $\theta_{k}, k \geq 0$, the canonical shift on $\left(\mathbb{R}^{d}\right)^{\mathbb{N}}$. With the notation of (1.44), we introduce the event

$$
\begin{align*}
\bar{G}= & \left\{\omega \in \Omega ; J_{n, x, C_{n}(x), \gamma}=0, \text { for all } n_{0}^{\prime} \leq n \leq n_{0}+1, x \in L_{n} \mathbb{Z}^{d} \cap\left(5 \mathcal{T}_{n_{0}+1}\right),\right.  \tag{5.2}\\
& \text { for all } \left.\gamma \in\left\{1, \ldots, 2 d 5^{(d-1)}\right\}\right\} .
\end{align*}
$$

This is the place where we use the control on traps to make sure that $\bar{G}^{c}$ has negligible probability. With (1.48), for $n \leq n_{0}$ and Proposition 3.3 when $n=n_{0}+1$, (we in fact only need in these controls the case of $\mathcal{A}$ singleton and $u_{x} \rightarrow 0$ ) we see that when $L_{0}$ is large,

$$
\begin{align*}
\mathbb{P}\left[\bar{G}^{c}\right] & \leq \sum_{n_{0}^{\prime} \leq n \leq n_{0}+1} c\left(\frac{L_{n_{0}+1}^{2}}{L_{n}}\right)^{d} L_{n}^{-\bar{M}_{n}} \\
& \leq c\left(m_{0}+2\right) L_{n_{0}+1}^{2 d-(1+a)^{-\left(m_{0}+2\right)} M / 2} \underset{(1.14),(1.17)}{(\leq} L_{n_{0}+1}^{-2 M_{0}} . \tag{5.3}
\end{align*}
$$

With the notations of Proposition 4.11, (4.5), (4.136), we define for each $\sigma \in \Sigma$ :

$$
\begin{equation*}
\bar{G}_{\sigma, n_{0}+1}=G_{\sigma, n_{0}+1} \cap \bar{G} \subseteq G_{\sigma} \cap \widetilde{G} \cap \bar{G} . \tag{5.4}
\end{equation*}
$$

When $L_{0}$ is large with (4.178), (5.3), we find:

$$
\begin{equation*}
\mathbb{P}\left[\left(\bigcup_{\sigma \in \Sigma} \bar{G}_{\sigma, n_{0}+1}\right)^{c}\right] \leq \mathbb{P}\left[\left(\bigcup_{\sigma \in \Sigma} G_{\sigma, n_{0}+1}\right)^{c}\right]+\mathbb{P}\left[\bar{G}^{c}\right] \leq \frac{1}{10} L_{n_{0}+1}^{-M_{0}} \tag{5.5}
\end{equation*}
$$

The next proposition is an important step in our program of "defects repairs". Some elements are reminiscent of Sidoravicius-Sznitman [25], cf. below (2.33) of [25].

Proposition 5.1. When $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$, $f$ with $|f|_{\left(n_{0}+1\right)} \leq 1$,

$$
\begin{align*}
& \sup _{|y| \leq \widetilde{D}_{n_{0}+1}}\left|E_{y, \omega}^{e}\left[f\left(Z_{T}\right)\right]-E_{y, \omega}^{\sigma}\left[f\left(Z_{T}\right)\right]\right| \leq L_{n_{0}+1}^{-(\beta+\delta+a)} \text {, with } \\
& T=\left(\frac{L_{n_{0}+1}}{L_{n_{0}^{\prime}}}\right)^{2}-1 \stackrel{(4.138)}{=} k_{n_{0}+1}-1 . \tag{5.6}
\end{align*}
$$

Proof. We break the difference in (5.6) into three terms that will be separately bounded.
Recall from (4.4) that $\sigma=\left(\sigma_{1}, \ldots, \sigma_{\widetilde{\ell}}\right)$, where $0 \leq \widetilde{\ell} \leq \widetilde{\ell}_{0}$. We introduce

$$
\begin{equation*}
K_{\sigma}=\bigcup_{i=1}^{\tilde{\ell}} \bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right), U_{\sigma}=\bigcup_{i=1}^{\ell} B\left(\sigma_{i}, \frac{1}{5 \widetilde{\ell}_{0}} L_{n_{0}^{\prime}+2}\right), \tag{5.7}
\end{equation*}
$$

and write for $y \in B\left(0, \widetilde{D}_{n_{0}+1}\right)$, (cf. (5.6)),

$$
\begin{align*}
A_{1} & =E_{y, \omega}^{e}\left[f\left(Z_{T}\right), H_{K_{\sigma}}>T\right]-E_{y, \omega}^{\sigma}\left[f\left(Z_{T}\right), H_{K_{\sigma}}>T\right] \\
A_{2} & =E_{y, \omega}^{e}\left[f\left(Z_{T}\right), \frac{T}{2}<H_{K_{\sigma}} \leq T\right]-E_{y, \omega}^{\sigma}\left[f\left(Z_{T}\right), \frac{T}{2}<H_{K_{\sigma}} \leq T\right]  \tag{5.8}\\
A_{3} & =E_{y, \omega}^{e}\left[f\left(Z_{T}\right), H_{K_{\sigma}} \leq \frac{T}{2}\right]-E_{y, \omega}^{\sigma}\left[f\left(Z_{T}\right), H_{K_{\sigma}} \leq \frac{T}{2}\right]
\end{align*}
$$

(incidentally note that $A_{2}=A_{3}=0$, when $\sigma=\emptyset$ ). We thus have

$$
\begin{equation*}
E_{y, \omega}^{e}\left[f\left(Z_{T}\right)\right]-E_{y, \omega}^{\sigma}\left[f\left(Z_{T}\right)\right]=A_{1}+A_{2}+A_{3} \tag{5.9}
\end{equation*}
$$

We first bound $A_{1}$. Note that when $L_{0}$ is large, for $y \in B\left(0, \widetilde{D}_{n_{0}+1}\right), \sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$,

$$
\begin{equation*}
P_{y, \omega}^{\sigma} \text {-a.s. }, T<T_{\frac{1}{5} \tau_{n_{0}+1}} \tag{5.10}
\end{equation*}
$$

indeed, $T \leq\left(L_{n_{0}+1} / L_{n_{0}^{\prime}}\right)^{2}<L_{n_{0}+1}^{2} / 10 \widetilde{D}_{n_{0}^{\prime}}$, when $L_{0}$ is large, see also (4.7). Coming back to the diffusion process, we can write, cf. (4.7):

$$
\begin{align*}
A_{1}=E_{y, \omega}[ & \left.f\left(X_{T L_{n_{0}^{\prime}}^{2}}\right), X_{k L_{n_{0}^{\prime}}^{2}} \notin K_{\sigma}, \text { for } 0 \leq k \leq T\right]  \tag{5.11}\\
& -E_{y, \omega}\left[f\left(X_{V_{T}}\right), X_{V_{k}} \notin K_{\sigma}, \text { for } 0 \leq k \leq T\right],
\end{align*}
$$

where $V_{k}, k \geq 0$, are the iterates of the stopping time $L_{n_{0}^{\prime}}^{2} \wedge T_{n_{0}^{\prime}}$ on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$, cf. (1.19), that is:

$$
\begin{equation*}
V_{0}=0, V_{1}=L_{n_{0}^{\prime}}^{2} \wedge T_{n_{0}^{\prime}}, \text { and } V_{k+1}=V_{1} \circ \theta_{V_{k}}+V_{k}, \text { for } k \geq 1, \tag{5.12}
\end{equation*}
$$

(here of course $\left(\theta_{t}\right)_{t \geq 0}$ stands for the canonical shift on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$ ). With (5.10), (5.11), we see that:

$$
\begin{align*}
& \left|A_{1}\right| \leq 2 \sum_{0 \leq k<T} P_{y, \omega}\left[T_{n_{0}^{\prime}} \circ \theta_{m L_{n_{0}^{\prime}}^{2}}>L_{n_{0}^{\prime}}^{2}, \text { for } 0 \leq m<k, T_{n_{0}^{\prime}} \circ \theta_{k L_{n_{0}^{\prime}}^{2}} \leq L_{n_{0}^{\prime}}^{2},\right. \\
& \left.\quad \text { and } X_{m L_{n_{0}^{\prime}}^{2}} \in \mathcal{T}_{n_{0}+1} \backslash K_{\sigma}, \text { for } 0 \leq m \leq k\right]  \tag{5.13}\\
& \begin{array}{l}
(2.2),(4.5)
\end{array} e^{-\kappa_{n_{0}^{\prime}}} \leq e^{-\kappa_{n_{0}+1}} .
\end{align*}
$$

We now bound $A_{2}$, and by the remark following (5.8), we may and will assume that $\sigma \neq \emptyset$. Note that:

$$
\begin{equation*}
A_{2} \leq P_{y, \omega}^{e}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right]+P_{y, \omega}^{\sigma}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right] . \tag{5.14}
\end{equation*}
$$

We can express both probabilities in the right member of (5.14) in terms of the diffusion process in a similar fashion as in (5.11). Using analogous bounds we see that

$$
\begin{equation*}
\left|P_{y, \omega}^{e}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right]-P_{y, \omega}^{\sigma}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right]\right| \leq e^{-\kappa_{n_{0}+1}} \tag{5.15}
\end{equation*}
$$

Further since $\omega \in \bar{G}_{\sigma, n_{0}+1} \subseteq \bar{G}$, see (5.4), it follows from (5.2), (1.44) with $n=n_{0}$, and the Markov property that for $y$ as in (5.6),

$$
\begin{equation*}
P_{y, \omega}\left[\sup _{0 \leq u \leq v \leq \frac{T}{4} L_{n_{0}^{\prime}}^{2}}\left|X_{v}-X_{u}\right|<\frac{L_{n_{0}}}{2}\right] \leq\left(1-c_{1}\right)^{\ell_{n_{0}}^{2} / 8} \leq e^{-\kappa_{n_{0}}+1} \tag{5.16}
\end{equation*}
$$

With a similar argument as in (3.68), one sees that on the complement of the event that appears in the above probability, $X$. must have exited the open set $\bigcup_{i=1}^{\tilde{\ell}} B\left(\sigma_{i}, \frac{L_{n_{0}}}{4 \ell_{0}}\right)$ by time $\frac{T}{4} L_{n_{0}^{\prime}}^{2}$. We hence find that

$$
\begin{aligned}
P_{y, \omega}^{\sigma}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right] \leq & P_{y, \omega}\left[X_{V_{m}} \notin K_{\sigma}, \text { for all } 0 \leq m \leq \frac{T}{2}, \text { and } X_{V_{k}} \in K_{\sigma},\right. \text { for some } \\
& \left.\frac{T}{2}<k \leq T, \text { and } \sup _{0 \leq u \leq \frac{T}{4} L_{n_{0}^{\prime}}^{2}} d\left(X_{u}, K_{\sigma}\right) \geq \frac{L_{n_{0}}}{4 \widetilde{\ell}_{0}}-10 \widetilde{D}_{n_{0}^{\prime}}\right]+e^{-\kappa_{n_{0}+1}} .
\end{aligned}
$$

Introducing the open set:

$$
\begin{equation*}
\mathcal{U}=\left\{z \in \mathbb{R}^{d} ; d\left(z, K_{\sigma}\right)<\frac{L_{n_{0}}}{4 \widetilde{\ell}_{0}}-11 \widetilde{D}_{n_{0}^{\prime}}\right\}, \tag{5.17}
\end{equation*}
$$

we see with a similar argument as in (5.13), using (5.10), that

$$
\begin{align*}
& P_{y, \omega}^{\sigma}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T\right] \leq P_{y, \omega}^{\sigma}\left[\frac{T}{2}<H_{K_{\sigma}} \leq T \wedge T_{\mathcal{I}_{n_{0}+1}}, T_{\mathcal{U}}<\frac{T}{2}\right]+e^{-\kappa_{n_{0}+1}} \\
& \leq \sup _{z \in \mathcal{I}_{n_{0}+1} \backslash \mathcal{U}} P_{z, \omega}^{\sigma}\left[H_{K_{\sigma}}<T \wedge T_{\mathcal{I}_{n_{0}+1}}\right]+e^{-\kappa_{n_{0}+1}} . \tag{5.18}
\end{align*}
$$

Coming back to (5.14), (5.15), we find

$$
\begin{equation*}
A_{2} \leq 2 \sup _{z \in \mathcal{I}_{n_{0}+1} \backslash \mathcal{U}} P_{z, \omega}^{\sigma}\left[H_{K_{\sigma}}<T \wedge T_{\mathcal{I}_{n_{0}+1}}\right]+e^{-\kappa_{n_{0}+1}} . \tag{5.19}
\end{equation*}
$$

The next step is to bound the first expression in the right-hand side of (5.19). To this end for $w \in 2 \mathcal{T}_{n_{0}+1}$, we introduce the function:

$$
n_{w}(z)= \begin{cases}n_{0}^{\prime}, & \text { if } D_{n_{0}^{\prime}+2}^{*} \geq|z-w|,  \tag{5.20}\\ \sup \left\{n \in\left[n_{0}^{\prime}, n_{0}\right] ;|z-w|>D_{n+1}^{*}\right\}, & \text { else },\end{cases}
$$

and the stopping time (for $Z$.):

$$
\tau_{w}=\left\{\begin{array}{l}
1, \text { when } n_{w}\left(Z_{0}\right)=n_{0}^{\prime}  \tag{5.21}\\
k_{n} \wedge \inf \left\{k \geq 0:\left|Z_{k}-Z_{0}\right| \geq D_{n_{w}\left(Z_{0}\right)}^{*}\right\}, \text { else }
\end{array}\right.
$$

(recall $k_{n}=\left(L_{n} / L_{n_{0}^{\prime}}\right)^{2}$, cf. (4.138), and $D_{n}^{*}$ is defined in (4.10)). We write below $n(z)$ for $n_{w}(z)$. We also introduce the function

$$
\begin{equation*}
f_{w}(z)=\left|\frac{z-w}{D_{n_{0}+1}^{*}}\right|^{-\gamma} \wedge 1, z \in \mathbb{R}^{d}, \text { with } \gamma=d-2-\frac{1}{100} . \tag{5.22}
\end{equation*}
$$

Lemma 5.2. When $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, w \in 2 \mathcal{T}_{n_{0}+1}, z \in\left(2 \mathcal{T}_{n_{0}+1}\right) \cap$ $B\left(w, L_{n_{0}}^{(1+\delta / 2)}\right)$, (cf. (1.40) for the definition of $\delta$ ), we have

$$
\begin{equation*}
E_{z, w}^{\sigma}\left[f_{w}\left(Z_{\tau_{w}}\right)\right] \leq f_{w}(z) \tag{5.23}
\end{equation*}
$$

Proof. When $|z-w| \leq D_{n_{0}+1}^{*},(5.23)$ is immediate. We thus assume that

$$
\begin{equation*}
z_{0} \stackrel{\text { def }}{=} z-w \text { satisfies }\left|z_{0}\right|>D_{n_{0}+1}^{*}, \text { and } z \in\left(2 \mathcal{T}_{n_{0}+1}\right) \cap B\left(w, L_{n_{0}}^{\left(1+\frac{\delta}{2}\right)}\right) \tag{5.24}
\end{equation*}
$$

Consider $x \in \mathbb{R}^{d}$, such that $|x| \leq \frac{1}{2}\left|z_{0}\right|$. Writing $\widehat{z}_{0}=\frac{z_{0}}{\left|z_{0}\right|}$, we have

$$
\begin{align*}
& \left|z_{0}+x\right|^{-\gamma}=\left|z_{0}\right|^{-\gamma}\left|\widehat{z}_{0}+\frac{x}{\left|z_{0}\right|}\right|^{-\gamma}=\left|z_{0}\right|^{-\gamma}\left(1+2 \widehat{z}_{0} \cdot \frac{x}{\left|z_{0}\right|}+\frac{|x|^{2}}{\left|z_{0}\right|^{2}}\right)^{-\frac{\gamma}{2}} \\
& =\left|z_{0}\right|^{-\gamma}\left(1-\frac{\gamma}{2}\left(2 \widehat{z}_{0} \cdot \frac{x}{\left|z_{0}\right|}+\frac{|x|^{2}}{\left|z_{0}\right|^{2}}\right)+\frac{1}{2}\left(\gamma^{2}+2 \gamma\right)\left(\frac{\widehat{z}_{0} \cdot x}{\left|z_{0}\right|}\right)^{2}+r\left(z_{0}, x\right)\right) \text {, }  \tag{5.25}\\
& \text { with }\left|r\left(z_{0}, x\right)\right| \leq c\left(\frac{|x|}{\left|z_{0}\right|}\right)^{3}
\end{align*}
$$

after the application of Taylor's formula to second order in the neighborhood of 0 , to the function $(1+u)^{-\gamma / 2}, u \in(-1,1)$. Coming back to (5.21), with (5.24) in force, we see that

$$
\begin{gather*}
E_{z, \omega}^{\sigma}\left[f_{w}\left(Z_{\tau_{w}}\right)\right] \leq f_{w}(z)\left(1-\frac{\gamma}{\left|z_{0}\right|} \widehat{z}_{0} \cdot E_{z, \omega}^{\sigma}\left[Z_{\tau_{w}}-Z_{0}\right]-\frac{\gamma}{2\left|z_{0}\right|^{2}} E_{z, \omega}^{\sigma}\left[\left|Z_{\tau_{w}}-Z_{0}\right|^{2}\right]\right. \\
\left.+\frac{1}{2} \frac{\left(\gamma^{2}+2 \gamma\right)}{\left|z_{0}\right|^{2}} E_{z, \omega}^{\sigma}\left[\left\{\widehat{z}_{0} \cdot\left(Z_{\tau_{w}}-Z_{0}\right)\right\}^{2}\right]+c\left(\frac{D_{n(z)}^{*}}{\left|z_{0}\right|}\right)^{3}\right) . \tag{5.26}
\end{gather*}
$$

Comparing the law of $Z_{\tau_{w}}$ under $P_{z, \omega}^{\sigma}$ with $\widetilde{R}_{n(z), \sigma}^{*}(z, \cdot)$, cf. (4.138), with (4.139), and $\omega \in \bar{G}_{\sigma, n_{0}+1}$, we see that when $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, w, z \in 2 \mathcal{T}_{n_{0}+1}$, with (5.24):

$$
\begin{align*}
& \left|E_{z, \omega}^{\sigma}\left[Z_{\tau_{w}}-Z_{0}\right]-\widetilde{d}_{n(z), \sigma}^{*}(z, \omega)\right| \leq e^{-\kappa_{n_{0}}}  \tag{5.27}\\
& \left|E_{z, \omega}^{\sigma}\left[\left(Z_{\tau_{w}}-Z_{0}\right)_{i}\left(Z_{\tau_{\omega}}-Z_{0}\right)_{j}\right]-\alpha_{n(z)} \delta_{i j} L_{n(z)}^{2}-\left(\widetilde{\gamma}_{n(z), \sigma}^{*}\right)^{i, j}(z, \omega)\right| \leq e^{-\kappa_{n_{0}}},
\end{align*}
$$

for $1 \leq i, j \leq d$, with the notation of (4.14). Using (4.179), (1.49), and once again an analogous calculation as in Lemma 2.1, we see that under the same conditions as in (5.27)

$$
\begin{equation*}
\left|{\widetilde{d_{n}}(z), \sigma}_{*}(z, \omega)\right| \leq \kappa_{n_{0}} L_{n(z)} \nu_{n(z)},\left|\widetilde{\gamma}_{n(z), \sigma}^{*}(z, \omega)\right| \leq \kappa_{n_{0}} L_{n(z)}^{2} \nu_{n(z)} \tag{5.28}
\end{equation*}
$$

As a result we obtain, (recall $\left.\gamma+2-d=-\frac{1}{100}\right)$ :

$$
(\gamma+2) E_{z, \omega}^{\sigma}\left[\left\{\widehat{z}_{0} \cdot\left(Z_{\tau_{w}}-Z_{0}\right)\right\}^{2}\right]-E_{z, \omega}^{\sigma}\left[\left|Z_{\tau_{w}}-Z_{0}\right|^{2}\right] \leq-\frac{1}{100} \alpha_{n(z)} L_{n(z)}^{2}+\kappa_{n_{0}} L_{n(z)}^{2} \nu_{n(z)}
$$

Therefore for large $L_{0}, \sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, w, z \in 2 \mathcal{T}_{n_{0}+1}$, with (5.24), we find

$$
\begin{align*}
& E_{y, \omega}^{\sigma}\left[f_{w}\left(Z_{\tau_{w}}\right)\right] \leq \\
& f_{w}(z)\left[1+\frac{\kappa_{n_{0}}}{\left|z_{0}\right|} L_{n(z)} \nu_{n(z)}-\frac{\gamma}{2\left|z_{0}\right|^{2}} L_{n(z)}^{2}\left(\frac{\alpha_{n(z)}}{100}-\kappa_{n_{0}} \nu_{n(z)}\right)+c\left(\frac{D_{n(z)}^{*}}{\left|z_{0}\right|}\right)^{3}\right]  \tag{5.29}\\
& \stackrel{(5.24),(5.20)}{\leq} f_{w}(z)\left[1+\frac{L_{n(z)}}{\left|z_{0}\right|}\left(\kappa_{n 0} \nu_{n(z)}-\frac{c L_{n(z)}}{\left|z_{0}\right|}\right)\right] \leq f_{w}(z)
\end{align*}
$$

using (5.26), (5.28), and (4.17). The claim (5.23) now follows.
Coming back to (5.19), (5.7), we see that

$$
\begin{align*}
A_{2} & \leq 2 \widetilde{\ell}_{0} \sup _{1 \leq i \leq \widetilde{\ell}} \sup _{z \in \mathcal{T}_{n_{0}+1}:\left|z-\sigma_{i}\right| \geq \frac{L_{n}}{4 \tilde{0}_{0}}-\widetilde{D}_{n_{0}^{\prime}}} P_{z, \omega}^{\sigma}\left[H_{\bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)}<T \wedge T_{\mathcal{T}_{n_{0}+1}}\right]+e^{-\kappa_{n_{0}+1}} \\
& \leq 2 \widetilde{\ell}_{0} \sup _{1 \leq i \leq \widetilde{\ell}} \sup _{z \in \mathcal{T}_{n_{0}+1}: \frac{L_{n_{0}}}{4 \ell_{0}}-\widetilde{D}_{n_{0}^{\prime}} \leq\left|z-\sigma_{i}\right| \leq \frac{L_{0}}{4 \ell_{0}}} P_{z, \omega}^{\sigma}\left[H_{\bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)}<T \wedge T_{\mathcal{T}_{n_{0}+1}}\right]+e^{-\kappa_{n_{0}+1}} \tag{5.30}
\end{align*}
$$

using the strong Markov property in the last step.
With (4.139), $n=n_{0}$, and the Markov property, we observe that for large $L_{0}, \sigma \in$ $\Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, z \in \mathcal{T}_{n_{0}+1}$,

$$
\begin{equation*}
P_{z, \omega}^{\sigma}\left[\sup _{0 \leq k \leq T}\left|Z_{k}-Z_{0}\right|>\ell_{n_{0}}^{2} 30 \widetilde{\ell}_{0} \widetilde{D}_{n_{0}}\right] \leq e^{-\kappa_{n_{0}+1}} \tag{5.31}
\end{equation*}
$$

As a result when $z \in \mathcal{T}_{n_{0}+1}$ is such that for some $1 \leq i \leq \ell, \frac{L_{n_{0}}}{4 \widetilde{\ell}_{0}}-\widetilde{D}_{n_{0}^{\prime}} \leq\left|z-\sigma_{i}\right| \leq \frac{L_{n_{0}}}{4 \widetilde{\ell}_{0}}$, with (1.14), (1.40), we find

$$
\begin{equation*}
P_{z, \omega}^{\sigma}\left[H_{\bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)}<T \wedge T_{\mathcal{T}_{n_{0}+1}}\right] \leq P_{z, \omega}^{\sigma}\left[H_{\bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)}<T_{B\left(\sigma_{i}, L_{n_{0}}^{(1+\delta / 2)}\right)}\right]+e^{-\kappa_{n_{0}+1}} \tag{5.32}
\end{equation*}
$$

We can then introduce $\tau_{\sigma_{i}}^{k}, k \geq 0$, the iterates of the stopping time $\tau_{\sigma_{i}}$, cf. (5.21) with $w=\sigma_{i}$.

$$
\begin{equation*}
\tau_{\sigma_{i}}^{0}=0, \tau_{\sigma_{i}}^{1}=\tau_{\sigma_{i}}, \quad \tau_{\sigma_{i}}^{k+1}=\tau_{\sigma_{i}} \circ \theta_{\tau_{\sigma_{i}}^{k}}+\tau_{\sigma_{i}}^{k}, \text { for } k \geq 1 \tag{5.33}
\end{equation*}
$$

as well as

$$
\begin{equation*}
N=\inf \left\{k \geq 0 ; Z_{\tau_{\sigma_{i}}^{k}} \in \bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right) \cup B\left(\sigma_{i}, L_{n_{0}}^{(1+\delta / 2)}\right)^{c}\right\} \tag{5.34}
\end{equation*}
$$

Using induction over $k$, the strong Markov property and (5.23), we see that

$$
\begin{equation*}
E_{z, \omega}^{\sigma}\left[f_{\sigma_{i}}\left(Z_{\tau_{\sigma_{i}}^{N} \wedge k}\right)\right] \text { is a decreasing function of } k \geq 0 \tag{5.35}
\end{equation*}
$$

Further observe that for $z$ as above (5.32), $P_{z, \omega}^{\sigma}$-a.s., on the event $\left\{H_{\bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)}<\right.$ $\left.T_{B\left(\sigma_{i}, L_{n_{0}}^{(1+\delta / 2)}\right)}\right\}$, it holds that $Z_{\tau_{\sigma_{i}}^{N}} \in \bar{B}\left(\sigma_{i}, 10 \widetilde{D}_{n_{0}^{\prime}}\right)$, as follows from (5.21), (5.33), (5.34). Hence with Fatou's lemma, we find

$$
\begin{equation*}
P_{z, \omega}^{\sigma}\left[H_{\bar{B}\left(\sigma_{i}, 10 \tilde{D}_{n_{0}^{\prime}}\right)}<T_{B\left(\sigma_{i}, L_{n_{0}}^{(1+\delta / 2)}\right)}\right] \leq E_{z, \omega}^{\sigma}\left[f_{\sigma_{i}}\left(Z_{\tau_{\sigma_{i}}^{N}}\right), N<\infty\right] \leq f_{\sigma_{i}}(z) \tag{5.36}
\end{equation*}
$$

The above inequality together with (5.22), (5.30), shows that when $L_{0}$ is large,

$$
\begin{align*}
& A_{2} \leq \kappa_{n_{0}+1}\left(\frac{L_{n_{0}}}{L_{n_{0}^{\prime}+1}}\right)^{-\left(d-2-\frac{1}{100}\right)}+e^{-\kappa_{n_{0}+1}}  \tag{5.37}\\
& \quad \begin{array}{l}
(4.1) \\
\leq \kappa_{n_{0}+1} L_{n_{0}+1}^{-\frac{99}{100}\left((1+a)^{-1}-(1+a)^{-\left(m_{0}+1\right)}\right)^{(1.14),(1.17)}} \leq L_{n_{0}+1}^{-\frac{8}{10}} .
\end{array} .=\frac{1}{\leq} .
\end{align*}
$$

We now bound $A_{3}$. As in the case of $A_{2}$, we only need to consider the case $\sigma \neq \emptyset$, see below (5.8). We first introduce some notations. We consider the functions, (with $\omega \in \bar{G}_{\sigma, n_{0}+1}$, and $f$ as in (5.9)):

$$
\begin{equation*}
F^{e}(k, z)=E_{z, \omega}^{e}\left[f\left(Z_{T-k}\right)\right], F^{\sigma}(k, z)=E_{z, \omega}^{\sigma}\left[f\left(Z_{T-k}\right)\right], z \in \mathbb{R}^{d}, 0 \leq k \leq T \tag{5.38}
\end{equation*}
$$

We also introduce the probability kernels:

$$
\begin{align*}
& Q^{e} G(k, z)=E_{z, \omega}^{e}\left[G\left(\left(k+T_{U_{\sigma}} \wedge t_{0}\right) \wedge T, Z_{T_{U_{\sigma}} \wedge t_{0} \wedge(T-k)}\right)\right], 0 \leq k \leq T, z \in \mathbb{R}^{d}  \tag{5.39}\\
& Q^{\sigma} G(k, z)=E_{z, \omega}^{\sigma}\left[G\left(\left(k+T_{U_{\sigma}} \wedge t_{0}\right) \wedge T, Z_{T_{U_{\sigma}} \wedge t_{0} \wedge(T-k)}\right)\right], 0 \leq k \leq T, z \in \mathbb{R}^{d}
\end{align*}
$$

with $G$ bounded measurable on $\{0, \ldots, T\} \times \mathbb{R}^{d}, U_{\sigma}$ as in (5.7), and

$$
\begin{equation*}
t_{0}=k_{n_{0}^{\prime}+3} \stackrel{(4.138)}{=}\left(L_{n_{0}^{\prime}+3} / L_{n_{0}^{\prime}}\right)^{2} \tag{5.40}
\end{equation*}
$$

Loosely speaking, these kernels describe for the Markov chain in the true environment or in the environment after surgery how the process initiated at time $k \leq T$, and stopped at the deterministic time $T \wedge\left(k+t_{0}\right)$ quits $U_{\sigma}$. We also introduce sub-probability kernels describing returns to $K_{\sigma}$ prior to $T$ or exit from $\frac{3}{4} \mathcal{T}_{n_{0}+1}$ :

$$
\begin{align*}
& R^{e} G(k, z)=E_{z, \omega}^{e}\left[G\left(\left(k+H_{K_{\sigma}}\right) \wedge T, Z_{H_{K_{\sigma}} \wedge(T-k)}\right), H_{K_{\sigma}}<(T-k) \wedge T_{\frac{3}{4}} \mathcal{T}_{n_{0}+1}\right]  \tag{5.41}\\
& R^{\sigma} G(k, z)=E_{z, \omega}^{\sigma}\left[G\left(\left(k+H_{K_{\sigma}}\right) \wedge T, Z_{H_{K_{\sigma}} \wedge(T-k)}\right), H_{K_{\sigma}}<(T-k) \wedge T_{\frac{3}{4}} \mathcal{T}_{n_{0}+1}\right]
\end{align*}
$$

with $0 \leq k \leq T, z \in \mathbb{R}^{d}$, and $G$ as below (5.39).

Coming back to the definition of $A_{3}$ in (5.8), we see using the strong Markov property at time $H_{K_{\sigma}}$, analogous considerations as in the control of $A_{1}$ and (5.10), that for large $L_{0}, \sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, y \in B\left(0, \widetilde{D}_{n_{0}+1}\right):$

$$
\begin{align*}
& \left|A_{3}-A_{3}^{\prime}\right| \leq e^{-\kappa_{n_{0}+1}}, \text { with } \\
& A_{3}^{\prime} \stackrel{\text { def }}{=} E_{y, \omega}^{\sigma}\left[H_{K_{\sigma}} \leq \frac{T}{2} \wedge T_{\frac{1}{5} \mathcal{T}_{n_{0}+1}}, F^{e}\left(H_{K_{\sigma}}, Z_{H_{K_{\sigma}}}\right)-F^{\sigma}\left(H_{K_{\sigma}}, Z_{H_{K_{\sigma}}}\right)\right] . \tag{5.42}
\end{align*}
$$

Applying the strong Markov property, we see that for $0 \leq k \leq T, z \in \mathbb{R}^{d}$ :

$$
\begin{align*}
F^{e}(k, z)-F^{\sigma}(k, z) & =Q^{e} F^{e}(k, z)-Q^{\sigma} F^{\sigma}(k, z) \\
& =Q^{e}\left(F^{e}-F^{\sigma}\right)(k, z)+\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}(k, z) \tag{5.43}
\end{align*}
$$

The next lemma will provide an analogue of (4.139) for the Markov chain in the true environment (i.e. under $P_{z, \omega}^{e}$ ).
Lemma 5.3. When $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, z \in 3 \mathcal{T}_{n_{0}+1}, n_{0}^{\prime} \leq n \leq n_{0}$ :

$$
\begin{equation*}
P_{z, \omega}^{e}\left[\sup _{0 \leq k \leq k_{n}}\left|Z_{k}-Z_{0}\right| \geq 30 \widetilde{\ell}_{0} \widetilde{D}_{n}\right] \leq e^{-\kappa_{n_{0}+1}} \tag{5.44}
\end{equation*}
$$

with $k_{n} \stackrel{(4.138)}{=}\left(L_{n} / L_{n_{0}^{\prime}}\right)^{2}$, and $\tilde{\ell}_{0}$ as below (4.2).
Proof. The argument is similar to the proof of (4.139). The probability in (5.44) coincides with

$$
\begin{equation*}
P_{z, \omega}\left[\sup _{0 \leq k \leq k_{n}}\left|X_{k L_{n_{0}^{\prime}}^{2}}-X_{0}\right| \geq 30 \widetilde{\ell}_{0} \widetilde{D}_{n}\right] . \tag{5.45}
\end{equation*}
$$

On the event inside the above probability, $X$. exits the open set $U$ defined below (4.142):

$$
U=\left(\bigcup_{1 \leq i \leq \widetilde{\ell}_{0}} B\left(w_{i}, 6 \widetilde{D}_{n}\right)\right) \cup\left(\bigcup_{1 \leq i \leq \tilde{\ell}} B\left(\sigma_{i}, 6 \widetilde{D}_{n}\right)\right),
$$

where the $w_{i}$ are omitted when $n=n_{0}^{\prime}$. We denote with $S$ the stopping time on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)$ :

$$
S=\inf \left\{s \geq 0,\left|X_{s}-x\right| \geq 4 \widetilde{D}_{n}, \text { for all } x \in\left\{w_{1}, \ldots, w_{\widetilde{\ell}_{0}}, \sigma_{1}, \ldots, \sigma_{\widetilde{\ell}}\right\}\right\}
$$

where the $w_{i}$ are omitted when $n=n_{0}^{\prime}$. From the discussion above, with the notation (1.18), the expression in (5.45) is smaller than:

$$
\begin{align*}
& E_{z, \omega}\left[S<L_{n}^{2}, P_{X_{S}, \omega}\left[X_{L_{n}^{2}}^{*} \geq \widetilde{D}_{n}\right]\right] \stackrel{(2.10)}{\leq}  \tag{5.46}\\
& E_{z, \omega}\left[S<L_{n}^{2} \wedge T_{4 \tau_{n_{0}+1}}, P_{X_{S}, \omega}\left[X_{L_{n}^{2}}^{*} \geq \widetilde{D}_{n}\right]\right]+e^{-\kappa_{n_{0}+1}} \leq e^{-\kappa_{n_{0}+1}}
\end{align*}
$$

using the definition of $U$, and (2.2) in the last step. This proves the lemma.
We now work on the quantities that appear in the last line of (5.43). For $0 \leq k \leq$ $T, z \in \frac{1}{2} \mathcal{T}_{n_{0}+1}$, we can write:

$$
\begin{aligned}
& F^{e}(k, z)-F^{\sigma}(k, z)= \\
& E_{z, \omega}^{e}\left[H_{K_{\sigma}}<(T-k) \wedge T_{\frac{3}{4} \mathcal{T}_{n_{0}+1}}, f\left(Z_{T-k}\right)\right]-E_{z, \omega}^{\sigma}\left[H_{K_{\sigma}}<(T-k) \wedge T_{\frac{3}{4} \mathcal{I}_{n_{0}+1}}, f\left(Z_{T-k}\right)\right]+ \\
& E_{z, \omega}^{e}\left[H_{K_{\sigma}} \geq(T-k) \wedge T_{\frac{3}{4} \mathcal{T}_{n_{0}+1}}, f\left(Z_{T-k}\right)\right]-E_{z, \omega}^{\sigma}\left[H_{K_{\sigma}} \geq(T-k) \wedge T_{\frac{3}{4} \mathcal{I}_{n_{0}+1}}, f\left(Z_{T-k}\right)\right]
\end{aligned}
$$

Note that when $L_{0}$ is large $\ell_{n_{0}}^{2} \widetilde{D}_{n_{0}}<\frac{1}{8} L_{n_{0}+1}^{2}$, so that with (4.139) and (5.44) when $n=n_{0}$, difference of the last two terms of the above equality is bounded in absolute value by

$$
\begin{aligned}
& \left|E_{z, \omega}^{e}\left[H_{K_{\sigma}} \wedge T_{\frac{3}{4}}^{\mathcal{T}_{n_{0}+1}} \geq T-k, f\left(Z_{T-k}\right)\right]-E_{z, \omega}^{\sigma}\left[H_{K_{\sigma}} \wedge T_{\frac{3}{4}} \mathcal{I}_{n_{0}+1} \geq T-k, f\left(Z_{T-k}\right)\right]\right| \\
& +e^{-\kappa_{n_{0}+1}} \leq e^{-\kappa_{n_{0}+1}}
\end{aligned}
$$

using in the last step analogous estimates as for $A_{1}$, cf. (5.13). Further the terms in the first line of the right-hand side of the above equality are seen to coincide with $R^{e} F^{e}(k, z)-$ $R^{\sigma} F^{\sigma}(k, z)$, after application of the Markov property at time $H_{K_{\sigma}} \wedge(T-k)$. Using once again estimates as in the control of $A_{1}$, or in the derivation of (5.42), we see that $R^{e} F^{e}(k, z)-R^{\sigma} F^{\sigma}(k, z)$ differs at most by $e^{-\kappa_{n_{0}+1}}$ from $R^{e}\left(F^{e}-F^{\sigma}\right)(k, z)$. Collecting our bounds, we see that when $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, 0 \leq k \leq T, z \in \frac{1}{2} \mathcal{T}_{n_{0}+1}$ :

$$
\begin{equation*}
\left|\left(F^{e}-F^{\sigma}\right)(k, z)-R^{e}\left(F^{e}-F^{\sigma}\right)(k, z)\right| \leq e^{-\kappa_{n_{0}}+1} \tag{5.47}
\end{equation*}
$$

Letting $y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}$ play the role of $Z_{H_{\sigma}}$ in (5.42), and noting that in view of (5.39), (5.44), when $0 \leq k^{\prime} \leq T, Q^{e}\left(\left(k^{\prime}, y^{\prime}\right),\{0, \ldots, T\} \times\left(\frac{1}{2} \mathcal{T}_{n_{0}+1}\right)^{c}\right) \leq e^{-\kappa_{n_{0}+1}}$, we see with (5.43) and (5.47) that for $0 \leq k^{\prime} \leq T$ :

$$
\begin{equation*}
\left|\left(F^{e}-F^{\sigma}\right)\left(k^{\prime}, y^{\prime}\right)-Q^{e} R^{e}\left(F^{e}-F^{\sigma}\right)\left(k^{\prime}, y^{\prime}\right)-\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}\left(k^{\prime}, y^{\prime}\right)\right| \leq e^{-\kappa_{n_{0}+1}} \tag{5.48}
\end{equation*}
$$

Thanks to (5.43) the expression under the absolute value coincides with

$$
\begin{equation*}
\left[F^{e}-F^{\sigma}-Q^{e} R^{e} Q^{e}\left(F^{e}-F^{\sigma}\right)-\sum_{m=0}^{1}\left(Q^{e} R^{e}\right)^{m}\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}\right]\left(k^{\prime}, y^{\prime}\right) \tag{5.49}
\end{equation*}
$$

Using the strong Markov property, (5.39), (5.41), (2.1)

$$
\begin{align*}
& Q^{e} R^{e} Q^{e}\left(\left(k^{\prime}, y^{\prime}\right),\{0, \ldots, T\} \times\left(\frac{1}{2} \mathcal{T}_{n_{0}+1}\right)^{c}\right) \leq \\
& P_{y^{\prime}, \omega}^{e}\left[\sup _{k \leq T}\left|Z_{k}-Z_{0}\right| \geq \frac{1}{4} L_{n_{0}+1}^{2}\right] \stackrel{(5.44)}{\leq} e^{-\kappa_{n_{0}+1}} \tag{5.50}
\end{align*}
$$

Hence using (5.47) to transform (5.49), we deduce from (5.48), (5.50) that

$$
\begin{equation*}
\left|\left[F^{e}-F^{\sigma}-\left(Q^{e} R^{e}\right)^{2}\left(F^{e}-F^{\sigma}\right)-\sum_{m=0}^{1}\left(Q^{e} R^{e}\right)^{m}\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}\right]\left(k^{\prime}, y^{\prime}\right)\right| \leq e^{-\kappa_{n_{0}+1}} \tag{5.51}
\end{equation*}
$$

Note that (5.50) holds for $\left(Q^{e} R^{e}\right)^{m} Q^{e}, m \geq 0$, arbitrary in place of $\left(Q^{e} R^{e}\right) Q^{e}$, as follows from the strong Markov property. We can then repeat the above manipulation finitely many times and find that when $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}$, $0 \leq k^{\prime} \leq T$ :

$$
\begin{align*}
& \left|\left[F^{e}-F^{\sigma}-\left(Q^{e} R^{e}\right)^{m_{*}}\left(F^{e}-F^{\sigma}\right)-\sum_{0 \leq m<m_{*}}^{1}\left(Q^{e} R^{e}\right)^{m}\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}\right]\left(k^{\prime}, y^{\prime}\right)\right|  \tag{5.52}\\
& \leq e^{-\kappa_{n_{0}+1}}
\end{align*}
$$

with in the notations of (1.14), (1.17):

$$
\begin{equation*}
m_{*}=\left[a^{-1}(1+a)^{m_{0}+1}\right]+1 . \tag{5.53}
\end{equation*}
$$

Keeping in mind that $y^{\prime}$ plays the role of $Z_{H_{K \sigma}}$ and letting $k^{\prime}$ play the role of $H_{K_{\sigma}}$ in (5.42), we are now going to bound $\left[\left(Q^{e} R^{e}\right)^{m_{*}} 1\right]\left(k^{\prime}, y^{\prime}\right)$, for $0 \leq k^{\prime} \leq \frac{T}{2}, y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}$.

Lemma 5.4. When $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$, for $0 \leq k^{\prime} \leq \frac{T}{2}, y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}$,

$$
\begin{gather*}
\sup _{0 \leq m \leq m_{*}}\left(Q^{e} R^{e}\right)^{m}\left(\left(k^{\prime}, y^{\prime}\right),\left[\frac{3}{4} T, T\right] \times \mathbb{R}^{d}\right) \leq L_{n_{0}+1}^{-\frac{8}{10}},  \tag{5.54}\\
\left(Q^{e} R^{e}\right)^{m_{*}}\left(\left(k^{\prime}, y^{\prime}\right),[0, T] \times \mathbb{R}^{d}\right) \leq 2 L_{n_{0}+1}^{-\frac{8}{10}} \tag{5.55}
\end{gather*}
$$

Proof. We first prove (5.54). When $m=0$, the expression that appears in (5.54) vanishes, and we can restrict to the case $1 \leq m \leq m_{*}$. We can rewrite the quantity in (5.54) using the strong Markov property, (5.39), (5.41), as the $P_{y, \omega}^{e}$-probability of a certain event (loosely speaking expressing the occurrence of $m$ successive possibly truncated departures from $U_{\sigma}$ and returns to $K_{\sigma}$ prior to exit of $\frac{3}{4} \mathcal{T}_{n_{0}+1}$, with the $m$-th return taking place sometimes during $\left[\frac{3}{4} T-k^{\prime}, T-k^{\prime}\right)$ ). On this event since truncated departures have at most a duration of $t_{0}$, cf. (5.39), at least one of the return periods has a duration of at least

$$
\left(\frac{3}{4} T-k^{\prime}-m_{*} t_{0}\right) / m_{*} \geq \frac{T}{4 m_{*}}-t_{0}
$$

As a result we have:

$$
\begin{equation*}
\left(Q^{e} R^{e}\right)^{m}\left(\left(k^{\prime}, y^{\prime}\right),\left[\frac{3}{4} T, T\right] \times \mathbb{R}^{d}\right) \leq m \sup _{z \in \frac{3}{4} \mathcal{T}_{n_{0}+1}} \quad P_{z, \omega}^{e}\left[\frac{T}{4 m_{*}}-t_{0} \leq H_{K_{\sigma}}<T\right] \tag{5.56}
\end{equation*}
$$

The probability that appears in the right-hand side of (5.56) is similar to the first probability that appears in (5.14), $\left(y \in B\left(0, \widetilde{D}_{n_{0}+1}\right)\right.$ is now replaced with $z \in \frac{3}{4} \mathcal{T}_{n_{0}+1}$, and $\frac{T}{2}$ with $\frac{T}{4 m_{*}}-t_{0}$ ). The same estimates leading to (5.37) now yield for $L_{0}$ large:

$$
\begin{equation*}
m_{*} \sup _{z \in \frac{3}{4} \mathcal{T}_{n_{0}+1}} P_{z, \omega}^{e}\left[\frac{T}{4 m_{*}}-t_{0} \leq H_{K_{\sigma}}<T\right] \leq L_{n_{0}+1}^{-\frac{8}{10}} \tag{5.57}
\end{equation*}
$$

thus proving (5.54).
We now turn to the proof of (5.55). With (5.54) and using the strong Markov property in the second inequality, we find

$$
\begin{align*}
& \left(Q^{e} R^{e}\right)^{m_{*}}\left(\left(k^{\prime}, y^{\prime}\right),[0, T] \times \mathbb{R}^{d}\right) \leq L_{n_{0}+1}^{-\frac{8}{10}}+\left(Q^{e} R^{e}\right)^{m_{*}}\left(\left(k^{\prime}, y^{\prime}\right),\left(0, \frac{3 T}{4}\right) \times \mathbb{R}^{d}\right) \\
& \leq L_{n_{0}+1}^{-\frac{8}{10}}+\left(\sup _{z \in \frac{3}{4} \mathcal{T}_{n_{0}+1}} P_{z, \omega}\left[H_{K_{\sigma}} \circ \theta_{T_{U_{\sigma}} \wedge t_{0}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right]\right)^{m_{*}} . \tag{5.58}
\end{align*}
$$

The same argument employed in (5.16) - (5.18), shows that for $z \in \frac{3}{4} \mathcal{T}_{n_{0}+1}$, (recall $t_{0} \stackrel{(5.40)}{=}$ $\left.k_{n_{0}^{\prime}+3}\right)$ :

$$
\begin{equation*}
P_{z, \omega}^{e}\left[d\left(Z_{k}, K_{\sigma}\right) \geq \frac{L_{n_{0}^{\prime}+2}}{4 \widetilde{\ell}_{0}}-11 \widetilde{D}_{n_{0}^{\prime}}, \text { for some } 0 \leq k<t_{0}\right] \geq 1-e^{-\kappa_{n_{0}+1}} \tag{5.59}
\end{equation*}
$$

so that we find with (5.7)

$$
\begin{align*}
& P_{z, \omega}\left[H_{K_{\sigma}} \circ \theta_{T_{U_{\sigma}} \wedge t_{0}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right] \leq \\
& e^{-\kappa_{n_{0}+1}}+E_{z, \omega}^{e}\left[T_{U_{\sigma}}<t_{0}, P_{Z_{T_{U \sigma} \wedge t_{0}}, \omega}^{e}\left[H_{K_{\sigma}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right]\right] \tag{5.60}
\end{align*}
$$

But for $\bar{z} \in \mathcal{T}_{n_{0}+1} \backslash U_{\sigma}$ playing the role of $Z_{T_{U_{\sigma} \wedge t_{0}}, \omega}$ in the last term of (5.60), we find just as for (5.15):

$$
\begin{equation*}
P_{\bar{z}, \omega}^{e}\left[H_{K_{\sigma}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right] \leq P_{\bar{z}, \omega}^{\sigma}\left[H_{K_{\sigma}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right]+e^{-\kappa_{n_{0}+1}} \tag{5.61}
\end{equation*}
$$

The first term on the right-hand side of (5.61) can be bounded in the same way as in (5.30) - (5.37), to obtain with $L_{0}$ large:

$$
\begin{equation*}
P_{\bar{z}, \omega}^{\sigma}\left[H_{K_{\sigma}}<T_{\mathcal{T}_{n_{0}+1}} \wedge T\right] \leq \ell\left(\frac{c L_{n_{0}^{\prime}+2}}{D_{n_{0}^{\prime}+1}^{*}}\right)^{-\frac{99}{100}}+e^{-\kappa_{n_{0}+1}} \leq \ell_{n_{0}^{\prime}+1}^{-\frac{9}{10}} \tag{5.62}
\end{equation*}
$$

Coming back to (5.58), (5.60), we see that when $L_{0}$ is large, $\sigma \in \Sigma$, $\omega \in \bar{G}_{\sigma, n_{0}+1}, 0 \leq k^{\prime} \leq$ $\frac{T}{2}, y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}:$

$$
\begin{equation*}
\left.\left(Q^{e} R^{\sigma}\right)^{m_{*}}\left(\left(k^{\prime}, y^{\prime}\right),[0, T] \times \mathbb{R}^{d}\right)\right) \leq L_{n_{0}+1}^{-\frac{8}{10}}+\left(\ell_{n_{0}^{\prime}+1}^{-\frac{9}{10}}+e^{-\kappa_{n_{0}+1}}\right)^{m_{*}} \stackrel{(1.15),(5.53)}{\leq} 2 L_{n_{0}+1}^{-\frac{8}{10}} \tag{5.63}
\end{equation*}
$$

This proves the claim (5.55).
We return to (5.52), and observe with the help of the above lemma that when $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma_{n_{0}+1}}$, for $0 \leq k^{\prime} \leq \frac{T}{2}, y^{\prime} \in \frac{1}{4} \mathcal{T}_{n_{0}+1}$,

$$
\begin{equation*}
\left|\left(F^{e}-F^{\sigma}\right)\left(k^{\prime}, y^{\prime}\right)\right| \leq c\left(L_{n_{0}+1}^{-\frac{8}{10}}+\sup _{k \leq \frac{3}{4} T, z \in K_{\sigma} \cap\left(\frac{3}{4} \mathcal{I}_{n_{0}+1}\right)}\left|\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}(k, z)\right|\right) . \tag{5.64}
\end{equation*}
$$

We now bound the last term of (5.64). We consider $k \leq \frac{3}{4} T, z \in K_{\sigma} \cap\left(\frac{3}{4} \mathcal{T}_{n_{0}+1}\right)$, as above and introduce (recall $t_{0} \stackrel{(5.40)}{=} k_{n_{0}^{\prime}+3}$ )

$$
\begin{equation*}
\widetilde{k}=\inf \left\{m \in t_{0} \mathbb{Z}+T ; m \geq k+t_{0}\right\} \tag{5.65}
\end{equation*}
$$

With (5.39), and the Markov property in (5.38), we can write

$$
\begin{align*}
Q^{e} F^{\sigma}(k, z) & =E_{z, \omega}^{e}\left[F^{\sigma}\left(k+T_{U_{\sigma}} \wedge t_{0}, Z_{T_{U_{\sigma} \wedge t_{0}}}\right)\right] \\
& =E_{z, \omega}^{e}\left[E_{Z_{U_{\sigma} \wedge t_{0}}, \omega}^{\sigma}\left[F^{\sigma}\left(\widetilde{k}, Z_{\widetilde{k}-\bar{k}}\right)\right]\right], \tag{5.66}
\end{align*}
$$

where $\bar{k}=k+T_{U_{\sigma}} \wedge t_{0}$ is not part of the inner expectation. The same calculation for $Q^{\sigma} F^{\sigma}(k, z)$ and the strong Markov property yield:

$$
\begin{equation*}
Q^{\sigma} F^{\sigma}(k, z)=E_{z, \omega}^{\sigma}\left[F^{\sigma}\left(\widetilde{k}, Z_{\widetilde{k}-k}\right)\right] \tag{5.67}
\end{equation*}
$$

Using controls on the size of displacements of $Z$. in a time interval of length $t_{0}$ or $2 t_{0}$, under $P_{z, \omega}^{\sigma}$ or $P_{z, \omega}^{e}$, cf. (4.139), (5.44), we see that:

$$
\begin{align*}
& \sup _{k \leq \frac{3}{4} T, z \in K_{\sigma} \cap\left(\frac{3}{4} \mathcal{T}_{n_{0}+1}\right)}\left|\left(Q^{e}-Q^{\sigma}\right) F^{\sigma}(k, z)\right| \leq e^{-\kappa_{n_{0}+1}}+\operatorname{var} F^{\sigma} \text {, where }  \tag{5.68}\\
& \operatorname{var} F^{\sigma} \stackrel{\text { def }}{=} \sup \left\{\left|F^{\sigma}\left(\widetilde{k}, z_{1}\right)-F^{\sigma}\left(\widetilde{k}, z_{2}\right)\right|, z_{1}, z_{2} \in \mathcal{T}_{n_{0}+1},\left|z_{1}-z_{2}\right| \leq D_{n_{0}^{\prime}+3}^{*}\right.  \tag{5.69}\\
& \left.\widetilde{k} \in\left(t_{0} \mathbb{Z}+T\right) \cap\left[0, \frac{4}{5} T\right]\right\} .
\end{align*}
$$

We will bound var $F^{\sigma}$ with the help of the smoothness properties resulting from (4.179) and (5.38). We introduce a cut-off function $h$ with values in $[0,1]$ such that with (2.1):

$$
\begin{equation*}
h=1 \text { on } 2 \mathcal{T}_{n_{0}+1}, h=0 \text { on }\left(\frac{5}{2} \mathcal{T}_{n_{0}+1}\right)^{c}, \text { and }|h|_{\left(n_{0}+1\right)} \leq 1+\frac{c}{L_{n_{0}+1}^{\beta}} . \tag{5.70}
\end{equation*}
$$

Lemma 5.5. For large $L_{0}, \sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}, n_{0}^{\prime} \leq n \leq n_{0}$,

$$
\begin{equation*}
\left\|h R_{n, \sigma}^{*}\right\|_{n_{0}+1}=\left\|h\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}}\right\|_{n_{0}+1} \leq 1+\kappa_{n} \nu_{n} \tag{5.71}
\end{equation*}
$$

with $\nu_{n}$ defined in (4.17), and $k_{n} \stackrel{(4.138)}{=} L_{n}^{2} / L_{n_{0}^{\prime}}^{2}$.
Proof. The equality in (5.70) follows from (4.9), (5.69). Then with (4.9), (5.70), we can write

$$
\begin{equation*}
h\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n}}=h R_{n, \sigma}^{*} \stackrel{(4.8)}{=} h R_{n}^{0}+h S_{n, \sigma}^{*}=h R_{n}^{0}+h \widetilde{S}_{n, \sigma}^{*}+h\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right) . \tag{5.72}
\end{equation*}
$$

From (1.29), (1.55), (5.70) we have

$$
\begin{equation*}
\left\|h R_{n}^{0}\right\|_{n_{0}+1} \leq 1+\frac{c}{L_{n_{0}+1}^{\beta}} \tag{5.73}
\end{equation*}
$$

and from (4.140) we deduce

$$
\begin{equation*}
\left\|h\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right)\right\|_{n_{0}+1} \leq\left(\frac{L_{n_{0}+1}}{L_{n}}\right)^{\beta}\left\|h\left(S_{n, \sigma}^{*}-\widetilde{S}_{n, \sigma}^{*}\right)\right\|_{n} \leq\left(\frac{L_{n_{0}+1}}{L_{n}}\right)^{\beta} e^{-\kappa_{n_{0}}} \leq e^{-\kappa_{n_{0}+1}} \tag{5.74}
\end{equation*}
$$

If $g$ is such that $|g|_{\left(n_{0}+1\right)}=1$, and $x \in L_{n} \mathbb{Z}^{d}$ such that $\chi_{n, x} h \neq 0$, we can find $\widetilde{G}$ such that:

$$
\begin{equation*}
\text { Supp } \widetilde{G} \subseteq B\left(x, 4 D_{n}^{*}\right), \widetilde{G}=g-g(x) \text { on } B\left(x, 3 D_{n}^{*}\right),|\widetilde{G}|_{(n)} \leq \kappa_{n}\left(\frac{L_{n}}{L_{n_{0}+1}}\right)^{\beta} \tag{5.75}
\end{equation*}
$$

We thus see, cf. above (4.12), that with (1.49)

$$
\begin{align*}
& \left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*} g\right|_{(n)} \leq\left|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*} \widetilde{G}\right|_{(n)}+e^{-\kappa_{n}} \leq\left\|\chi_{n, x} \widetilde{S}_{n, \sigma}^{*}\right\|_{n} \kappa_{n}\left(\frac{L_{n}}{L_{n_{0}+1}}\right)^{\beta}+e^{-\kappa_{n}}  \tag{5.76}\\
& \stackrel{(4.179)}{\leq} \kappa_{n} \nu_{n}\left(\frac{L_{n}}{L_{n_{0}+1}}\right)^{\beta} .
\end{align*}
$$

As a consequence we see with (A.3) from the Appendix and (5.70) that

$$
\begin{equation*}
\left|h \widetilde{S}_{n, \sigma}^{*} g\right|_{\left(n_{0}+1\right)} \leq\left(\frac{L_{n_{0}+1}}{L_{n}}\right)^{\beta}\left|h \widetilde{S}_{n, \sigma}^{*} g\right|_{(n)} \leq \kappa_{n} \nu_{n}=\kappa_{n} \nu_{n}|g|_{n_{0}+1} \tag{5.77}
\end{equation*}
$$

Collecting (5.72), (5.73), (5.74), (5.77) we obtain (5.71).
We return to the task of bounding (5.69). With $\widetilde{k}$ as in (5.69) we have $T-\widetilde{k}-k_{n_{0}} \in t_{0} \mathbb{N}$, and hence we can write

$$
\begin{equation*}
T-\widetilde{k}-k_{n_{0}}=\sum_{n_{0}^{\prime}+3 \leq n \leq n_{0}} u_{n} k_{n}, \text { with } u_{n} \text { suitable integers in }\left[0, \ell_{n}^{2}-1\right) . \tag{5.78}
\end{equation*}
$$

Then for $z \in \mathcal{T}_{n_{0}+1}, f$ as in (5.6), (or (5.9)), we have:

$$
F^{\sigma}(\widetilde{k}, z) \stackrel{(5.38)}{=}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T-\widetilde{k}} f(z)=\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n_{0}}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{u_{n_{0}^{\prime}+3^{2}} k_{n_{0}^{\prime}+3}} \ldots\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{u_{n} k_{n}} f(z) .
$$

Using (4.9) and $(T-\widetilde{k}) \widetilde{D}_{n_{0}^{\prime}}<\frac{1}{10} L_{n_{0}+1}^{2}$, cf. below (5.10), we find

$$
\begin{align*}
F^{\sigma}(\widetilde{k}, z) & =R_{n_{0}, \sigma}^{*}\left(h R_{n_{0}^{\prime}+3, \sigma}^{*}\right)^{u_{n_{0}^{\prime}+3}} \ldots\left(h R_{n, \sigma}^{*}\right)^{u_{n}} \ldots\left(h R_{n_{0}, \sigma}^{*}\right)^{u_{n_{0}}} f(z)  \tag{5.79}\\
& =R_{n_{0}, \sigma}^{*} \widetilde{f}(z),
\end{align*}
$$

where in view of (5.71), (5.78)

$$
\begin{equation*}
|\widetilde{f}|_{\left(n_{0}+1\right)} \leq \prod_{n_{0}^{\prime}+3 \leq n \leq n_{0}}\left(1+\kappa_{n} \nu_{n}\right)^{\ell_{n}^{2}} \leq \exp \left\{\sum_{n_{0}^{\prime}+3 \leq n \leq n_{0}} \kappa_{n} \nu_{n} \ell_{n}^{2}\right\}^{(1.15),(4.17)} \leq{ }^{\leq} c . \tag{5.80}
\end{equation*}
$$

So we see that for $z_{1}, z_{2} \in \mathcal{T}_{n_{0}+1}$, with $\left|z_{1}-z_{2}\right| \leq D_{n_{0}^{\prime}+3}^{*}$,

$$
\left.\begin{array}{l}
\left|F^{\sigma}\left(\widetilde{k}, z_{1}\right)-F^{\sigma}\left(\widetilde{k}, z_{2}\right)\right| \stackrel{(5.79)}{\leq}\left|R_{n_{0}}^{0} \widetilde{f}\left(z_{1}\right)-R_{n_{0}}^{0} \widetilde{f}\left(z_{2}\right)\right|+\left|S_{n_{0}, \sigma}^{*} \widetilde{f}\left(z_{1}\right)-S_{n_{0}, \sigma}^{*} \widetilde{f}\left(z_{2}\right)\right| \\
\underset{(1.49),(1.56)}{\leq} \frac{c D_{n_{0}^{\prime}+3}^{*}}{L_{n_{0}}}+c\left(\frac{D_{n_{0}^{\prime}+3}^{*}}{L_{n_{0}}}\right)^{\beta} \nu_{n_{0}} \stackrel{(4.17),(4.1)}{\leq}  \tag{5.81}\\
\kappa_{n_{0}}\left(L_{n_{0}+1}^{-\left(\frac{1}{1+a}-(1+a)^{-\left(m_{0}-1\right)}\right)}+L_{n_{0}+1}^{-\left(\frac{\beta}{1+a}+\frac{\beta}{4(a+1)}-\beta(1+a)^{-\left(m_{0}-1\right)}-\left(\frac{\beta}{4}-\delta\right)(1+a)^{-\left(m_{0}+2\right)}\right)}\right) \\
\underset{(1.14),(1.17)}{\leq} L_{n_{0}+1}^{-(\beta+\delta)}\left(L_{n_{0}+1}^{-2 a}+L_{n_{0}+1}^{-\left(\frac{\beta}{4(1+a)}-\delta-a\right.} \frac{\beta}{1+a}-\frac{\beta}{100}-\frac{1}{100}\left(\frac{\beta}{4}-\delta\right)\right)
\end{array}\right) \leq c L_{n_{0}+1}^{-(\beta+\delta+2 a)} .
$$

So we have shown that when $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$,

$$
\begin{equation*}
\operatorname{var} F^{\sigma} \leq c L_{n_{0}+1}^{-(\beta+\delta+2 a)} \tag{5.82}
\end{equation*}
$$

Collecting (5.42), (5.64), (5.68), we obtain since $\beta+\delta+2 a<\frac{8}{10}$,

$$
\begin{equation*}
A_{3} \leq c L_{n_{0}+1}^{-(\beta+\delta+2 a)} \tag{5.83}
\end{equation*}
$$

Substituting in (5.9) the bounds (5.13), (5.37), (5.83) we now obtain (5.6) and this concludes the proof of Proposition 5.1.

As an application of Proposition 4.11 and 5.1, we have
Proposition 5.6. When $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$,

$$
\begin{equation*}
\left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1}-\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}}\right)\right\|_{n_{0}+1} \leq c L_{n_{0}+1}^{-(\delta+a)} \tag{5.84}
\end{equation*}
$$

Proof. We have

$$
\begin{align*}
& \left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1}-\left(R_{n_{0}}^{0}\right) \ell_{n_{0}}^{2}\right)\right\|_{n_{0}+1} \leq \\
& \left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1}-R_{n_{0}+1, \sigma}^{*}\right)\right\|_{n_{0}+1}+\left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1, \sigma}^{*}-\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}}\right)\right\|_{n_{0}+1} \stackrel{(4.180)}{\leq}  \tag{5.85}\\
& \left\|\chi_{n_{0}+1,0}\left(R_{n_{0}+1}-R_{n_{0}+1, \sigma}^{*}\right)\right\|_{n_{0}+1}+\nu_{n_{0}+1} .
\end{align*}
$$

With the notation of (5.6), and with (4.9), we also find that:

$$
\begin{align*}
& \chi_{n_{0}+1,0}\left(R_{n_{0}+1}-R_{n_{0}+1, \sigma}^{*}\right)=\chi_{n_{0}+1,0}\left(R_{n_{0}^{\prime}}\left(R_{n_{0}^{\prime}}\right)^{T}-R_{n_{0}^{\prime}, \sigma}^{*}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right)=  \tag{5.86}\\
& \chi_{n_{0}+1,0} R_{n_{0}^{\prime}}\left(\left(R_{n_{0}^{\prime}}\right)^{T}-\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right)+\chi_{n_{0}+1,0}\left(R_{n_{0}^{\prime}}-R_{n_{0}^{\prime}, \sigma}^{*}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T} .
\end{align*}
$$

With (1.60) and (5.6), we see that

$$
\begin{equation*}
\left\|R_{n_{0}^{\prime}} 1_{B\left(0, \widetilde{D}_{n_{0}+1}\right)}\left(\left(R_{n_{0}^{\prime}}\right)^{T}-\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right)\right\|_{n_{0}+1} \leq\left(\frac{L_{n_{0}+1}}{L_{n_{0}^{\prime}}}\right)^{\beta} c L_{n_{0}^{\prime}}^{\beta} L_{n_{0}+1}^{-(\beta+\delta+a)} \leq c L_{n_{0}+1}^{-(\delta+a)} \tag{5.87}
\end{equation*}
$$

Also note that when $|g|_{\infty} \leq 2$ and $g 1_{B\left(0, \widetilde{D}_{n_{0}+1}\right)}=0$, then with the notation (1.57), $\chi_{n_{0}+1,0} R_{n_{0}^{\prime}} g=\chi_{n_{0}+1,0} P_{1, \omega} P_{L_{n_{0}^{\prime}}^{2}-1, \omega} g$, and from inequalities such as in (2.10), and from (1.17), we see that $\left|1_{B\left(0, D_{\left.n_{0}+1\right)}\right.} P_{L_{n_{0}^{\prime}}^{2}-1, \omega} g\right|_{\infty} \leq e^{-c L_{n_{0}^{\prime}}}$, so that using (1.59) as in the proof of (1.60), we find that $\left|\chi_{n_{0}+1,0} R_{n_{0}^{\prime}} g\right|_{\left(n_{0}+1\right)} \leq e^{-c L_{n_{0}^{\prime}}}$. Coming back to (5.87), we hence obtain:

$$
\begin{equation*}
\left\|\chi_{n_{0}+1,0} R_{n_{0}^{\prime}}\left(\left(R_{n_{0}^{\prime}}\right)^{T}-\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right)\right\|_{n_{0}+1} \leq c L_{n_{0}+1}^{-(\delta+a)} . \tag{5.88}
\end{equation*}
$$

We now turn to the last term of (5.86) and observe that:

$$
R_{n_{0}^{\prime}}-R_{n_{0}^{\prime}, \sigma}^{*} \stackrel{(4.7)}{=}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right)+g_{\sigma}\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}\right) .
$$

With the same argument employed above (5.88), cf. (1.20), (1.37), for the notation, applied to the last expression of the following identity

$$
\begin{aligned}
& \chi_{n_{0}+1,0} g_{\sigma}\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}=\chi_{n_{0}+1,0} g_{\sigma}\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}\right) \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}+ \\
& \chi_{n_{0}+1,0} g_{\sigma} R_{n_{0}^{\prime}}\left(1-\chi_{D_{n_{0}+1}}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T} \stackrel{\text { def }}{=} A_{1}+A_{2},
\end{aligned}
$$

we see that $\left\|A_{2}\right\|_{n_{0}+1}$ is smaller than $e^{-c L_{n_{0}^{\prime}}}$. Further just as in (5.80) we see that:

$$
\left\|\chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1} \leq c
$$

and together with (4.6), (2.2), (2.46) we obtain:

$$
\begin{equation*}
\left\|A_{1}\right\|_{n_{0}+1} \leq\left\|\chi_{n_{0}+1,0} g_{\sigma}\left(S_{n_{0}^{\prime}}-\widetilde{S}_{n_{0}^{\prime}}\right) \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1}+e^{-\kappa_{n_{0}+1}} \leq e^{-\kappa_{n_{0}+1}} \tag{5.89}
\end{equation*}
$$

In view of the identity below (5.88), to control the rightmost expression in (5.86), it remains to bound $\left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1}$. To this end in analogy with (1.20) we define the probability kernel

$$
\begin{align*}
& R_{n_{0}^{\prime}}^{*}(x, d y)=P_{x, \omega}\left[X_{L_{n_{0}^{\prime}}^{2} \wedge \wedge_{n_{0}^{\prime}}^{*}} \in d y\right], x \in \mathbb{R}^{d}, \omega \in \Omega, \text { with }  \tag{5.90}\\
& T_{n_{0}^{\prime}}^{*}=\inf \left\{u \geq 0, X_{u}^{*} \geq D_{n_{0}^{\prime}}^{*}\right\}, \text { cf. }(4.10),(1.18) \text { for the notation } .
\end{align*}
$$

As in Lemma 5.3, see in particular (5.46), we see that when $L_{0}$ is large, $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$, for $y \in B\left(0, \widetilde{D}_{n_{0}+1}\right)$,

$$
P_{y, \omega}\left[X_{L_{n_{0}^{\prime}}^{2}}^{*} \geq \frac{D_{n_{0}^{\prime}}^{*}}{2}\right] \leq e^{-\kappa_{n_{0}^{\prime}}}
$$

Then with a slight variation on the proof of Proposition 2.5, for $x \in L_{n_{0}^{\prime}} \mathbb{Z}^{d} \cap B\left(0, D_{n_{0}+1}\right)$,

$$
\begin{equation*}
\left\|\chi_{n_{0}^{\prime}, x}\left(R_{n_{0}^{\prime}}^{*}-R_{n_{0}^{\prime}}\right)\right\|_{n_{0}^{\prime}} \leq e^{-\kappa_{n_{0}^{\prime}}} \tag{5.91}
\end{equation*}
$$

Employing a similar identity as above (5.89) in the first inequality, and (5.91) in the second, we find

$$
\begin{aligned}
& \left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1} \leq \\
& \left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}^{*}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right) \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1}+ \\
& \left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right) R_{n_{0}^{\prime}}\left(1-\chi_{D_{n_{0}+1}}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1}+e^{-\kappa_{n_{0}^{\prime}}} \leq \\
& \left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}^{*}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right) \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1}+e^{-\kappa_{n_{0}^{\prime}}},
\end{aligned}
$$

with the same argument as applied above (5.88). Note that thanks to (1.60), (4.6), (5.91), $\left\|\left(1-g_{\sigma}\right) \chi_{n_{0}^{\prime}, x}\left(R_{n_{0}^{\prime}}^{*}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right)\right\|_{n_{0}^{\prime}} \leq c L_{n_{0}^{\prime}}^{\beta}$, with $x$ as above (5.91). For $f$ with $|f|_{\left(n_{0}+1\right)} \leq 1$, and writing $Q=\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}^{*}-\widetilde{R}_{n_{0}^{\prime}}^{0}\right)$, we also find

$$
\begin{equation*}
\chi_{n_{0}+1,0} Q \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T} f=\chi_{n_{0}+1,0} Q \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n_{0}}} \widetilde{f} \tag{5.92}
\end{equation*}
$$

where $\tilde{f}$ just as in $(5.79),(5.80)$ satisfies

$$
\begin{equation*}
|\widetilde{f}|_{\left(n_{0}+1\right)} \leq c \tag{5.93}
\end{equation*}
$$

Further if $x \in L_{n_{0}^{\prime}} \mathbb{Z}^{d}$ is such that $d\left(x\right.$, Supp $\left.\chi_{n_{0}+1,0}\right) \leq 30 \sqrt{d} L_{n_{0}^{\prime}}$, we can use a cut-off function and construct $\widetilde{H}_{1}, \widetilde{H}_{2}$ supported in $B\left(x, 3 D_{n_{0}^{\prime}}^{*}\right)$ (where $\chi_{D_{n_{0}+1}}(\cdot)=1$ ), such that in $\bar{B}\left(x, 2 D_{n_{0}^{\prime}}^{*}\right)$

$$
\begin{align*}
& \widetilde{H}_{1} \text { coincides with } R_{n_{0}}^{0} \widetilde{f}(\cdot)-R_{n_{0}}^{0} \widetilde{f}(x), \\
& \begin{aligned}
& \widetilde{H}_{2} \text { coincides with } S_{n_{0}, \sigma}^{*} \widetilde{f}(\cdot)-S_{n_{0}, \sigma}^{*} \widetilde{f}(x) \stackrel{(4.8),(4.9)}{=}\left(R_{n_{n}^{\prime}, \sigma}^{*}\right)^{k_{n_{0}}} \widetilde{f}(\cdot) \\
& \quad-\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{k_{n_{0}}} \widetilde{f}(x)-R_{n_{0}}^{0} \widetilde{f}(\cdot)+R_{n_{0}}^{0} \widetilde{f}(x),
\end{aligned} \tag{5.94}
\end{align*}
$$

and so that they satisfy the bounds

$$
\begin{equation*}
\left|\widetilde{H}_{1}\right|_{\left(n_{0}^{\prime}\right)} \stackrel{(1.56)}{\leq} \kappa_{n_{0}^{\prime}} \frac{L_{n_{0}^{\prime}}}{L_{n_{0}}},\left|\widetilde{H}_{2}\right|_{\left(n_{0}^{\prime}\right)} \stackrel{(4.179)}{\leq} \kappa_{n_{0}^{\prime}}\left(\frac{L_{n_{0}^{\prime}}}{L_{n_{0}}}\right)^{\beta} \nu_{n_{0}} . \tag{5.95}
\end{equation*}
$$

As a result we obtain

$$
\begin{aligned}
&\left|\chi_{n_{0}^{\prime}, x} Q \chi_{D_{n_{0}+1}}\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T} f\right|_{\left(n_{0}^{\prime}\right)} \leq\left|\chi_{n_{0}^{\prime}, x} Q \widetilde{H}_{1}\right|_{\left(n_{0}^{\prime}\right)}+\left|\chi_{n_{0}^{\prime}, x} Q \widetilde{H}_{2}\right|_{\left(n_{0}^{\prime}\right)} \\
& \underset{(5.2),(4.6)}{\leq} \kappa_{n_{0}^{\prime}} L_{n_{0}^{\prime}}^{\beta}\left(\frac{L_{n_{0}^{\prime}}}{L_{n_{0}}}+\left(\frac{L_{n_{0}^{\prime}}}{L_{n_{0}}}\right)^{\beta} \nu_{n_{0}}\right) .
\end{aligned}
$$

We thus find

$$
\begin{align*}
& \left\|\chi_{n_{0}+1,0}\left(1-g_{\sigma}\right)\left(R_{n_{0}^{\prime}}-\widetilde{R}_{n_{0}^{\prime}}\right)\left(R_{n_{0}^{\prime}, \sigma}^{*}\right)^{T}\right\|_{n_{0}+1} \leq \\
& \kappa_{n_{0}^{\prime}} L_{n_{0}^{\prime}}^{\beta}\left(\frac{L_{n_{0}+1}}{L_{n_{0}}}\right)^{\beta}\left(\left(\frac{L_{n_{0}^{\prime}}}{L_{n_{0}}}\right)^{1-\beta}+\nu_{n_{0}}\right)+e^{-\kappa_{n_{0}^{\prime}} \leq L_{n_{0}+1}^{-(\delta+a)}} \tag{5.96}
\end{align*}
$$

using similar calculations as in the bottom lines of (5.81). Collecting (5.88), (5.89), (5.96), we obtain (5.84).

Before concluding the proof of Theorem 1.1, we yet have to control the difference $\alpha_{n_{0}+1}-\alpha_{n_{0}}$.

Proposition 5.7. Under the assumptions of Theorem 1.1, when $L_{0}$ is large,

$$
\begin{equation*}
\left|\alpha_{n_{0}+1}-\alpha_{n_{0}}\right| \leq L_{n_{0}}^{-\left(1+\frac{9}{10}\right) \delta} \tag{5.97}
\end{equation*}
$$

Proof. Recall the definition of $\alpha_{n}$ in (1.22). In analogy with (2.5) we consider the function, cf. (1.37) for the notation:

$$
\begin{equation*}
f(z)=\chi_{2 \widetilde{D}_{n_{0}+1}}(z) \frac{|z|^{2}}{L_{n_{0}+1}^{2}}, z \in \mathbb{R}^{d} \tag{5.98}
\end{equation*}
$$

so that $|f|_{\left(n_{0}+1\right)} \leq \kappa_{n_{0}+1}$, and:

$$
\begin{equation*}
\alpha_{n_{0}+1} \stackrel{(1.22)}{=} \mathbb{E}\left[\widetilde{R}_{n_{0}+1} f(0)\right] . \tag{5.99}
\end{equation*}
$$

We denote with $\widetilde{\Omega}$ the event

$$
\begin{align*}
\widetilde{\Omega}= & \left\{\omega \in \Omega ; \text { for }|y| \leq 30 \sqrt{d} L_{n_{0}+1}, P_{y, \omega}\left[X_{L_{n_{0}+1}^{2}}^{*} \geq v\right] \leq \exp \left\{-\frac{v}{D_{n_{0}+1}}\right\}\right.  \tag{5.100}\\
& \text { for all } \left.v \geq D_{n_{0}+1}\right\} \cap\left\{\omega \in \Omega ; \text { for all } x \in L_{n_{0}} \mathbb{Z}^{d} \cap\left(5 \mathcal{T}_{n_{0}+1}\right), x \in \widetilde{B}_{n_{0}}(\omega)\right\}
\end{align*}
$$

With (2.9) and (1.47), we see that when $L_{0}$ is large,

$$
\begin{equation*}
\mathbb{P}\left[\widetilde{\Omega}^{c}\right] \leq \frac{1}{10} L_{n_{0}+1}^{-M_{0}}+c\left(\frac{L_{n_{0}+1}^{2}}{L_{n_{0}}}\right)^{d} L_{n_{0}}^{-M_{0}} \underset{(1.46)}{(1.14),(1.15)} L_{n_{0}+1}^{-10} . \tag{5.101}
\end{equation*}
$$

Then for $\omega \in \widetilde{\Omega}$, we see that (cf. (1.37) for the notation):

$$
\begin{aligned}
& \left|\widetilde{R}_{n_{0}+1} f(0)-\left(R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right)^{\ell_{n_{0}}^{2}} f(0)\right| \leq\left|\widetilde{R}_{n_{0}+1} f(0)-R_{n_{0}+1} f(0)\right|+ \\
& \left|\left(R_{n_{0}}^{0}+S_{n_{0}}\right)^{\ell_{n_{0}}^{2}} f(0)-\left(R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right)^{\ell_{n_{0}}^{2}} f(0)\right| \leq e^{-\kappa_{n_{0}+1}}+ \\
& \left|\sum_{0 \leq k<\ell_{n_{0}}^{2}}\left(R_{n_{0}}\right)^{k}\left(1-\chi_{\widetilde{D}_{n_{0}+1}}\right) S_{n_{0}}\left(R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right)^{\ell_{n_{0}}^{2}-k-1} f(0)\right|
\end{aligned}
$$

using (2.46) with $n=n_{0}+1$, and perturbation expansion in the last step. Since $R_{n_{0}}^{0}+$ $\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}=\left(1-\chi_{\widetilde{D}_{n_{0}+1}}\right) R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}}} R_{n_{0}}$ contracts the sup-norm, we see with (5.100), that when $L_{0}$ is large, for $\omega \in \widetilde{\Omega}$ :

$$
\begin{equation*}
\left|\widetilde{R}_{n_{0}+1} f(0)-\left(R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right)^{\ell_{n_{0}}^{2}} f(0)\right| \leq e^{-\kappa_{n_{0}+1}} \tag{5.102}
\end{equation*}
$$

Using perturbation expansion as in (4.15) we find that for all $\omega \in \Omega$ :

$$
\begin{align*}
& \left(R_{n_{0}}^{0}+\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right)^{\ell_{n_{0}}^{2}} f(0)-\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}} f(0)= \\
& \sum_{0 \leq k<\ell_{n_{0}}^{2}}\left(R_{n_{0}}^{0}\right)^{k} \chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}-k-1} f(0)+  \tag{5.103}\\
& \sum_{\substack{k 0+\cdots+k_{m}+m=\ell_{n_{0}}^{2} \\
k_{i} \geq 0, m \geq 2}} \chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\left(R_{n_{0}}^{0}\right)^{k_{1}} \ldots \chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\left(R_{n_{0}}^{0}\right)^{k_{m}} f(0) .
\end{align*}
$$

Further for $\omega \in \widetilde{\Omega},\left\|\chi_{\widetilde{D}_{n_{0}+1}} S_{n_{0}}\right\| \stackrel{(2.2),(2.46)}{\leq} c L_{n_{0}}^{-\delta}$, so that the term in the last line of (5.103) is smaller in absolute value than:

$$
\begin{align*}
\sum_{\substack{k_{0}+\cdots+k_{m}+m=\ell_{n_{0}}^{2} \\
k_{i} \geq 0, m \geq 2}}\left(c L_{n_{0}}^{-\delta}\right)^{m} \kappa_{n_{0}+1} & =\kappa_{n_{0}+1}\left[\left(1+c L_{n_{0}}^{-\delta}\right)^{\ell_{n_{0}}^{2}}-1-c \ell_{n_{0}}^{2} L_{n_{0}}^{-\delta}\right]  \tag{5.104}\\
& \leq \kappa_{n_{0}+1} L_{n_{0}}^{-2 \delta+4 a},
\end{align*}
$$

with $c$ denoting the same constant in both members of the equality, and using a similar argument as in (4.172).

Coming back to (5.102), (5.103), noting that $\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}} f(0)=P_{\alpha_{n_{0}} L_{n_{0}+1}^{2}} f(0)$, cf. (1.21), (1.54), and that in view of (1.49) i) and (5.98) this quantity differs at most by $e^{-\kappa_{n_{0}+1}}$ from $d \alpha_{n_{0}}$, we see that for $\omega \in \widetilde{\Omega}$ :

$$
\begin{align*}
& \left|\widetilde{R}_{n_{0}+1} f(0)-d \alpha_{n_{0}}-\sum_{0 \leq k<\ell_{n_{0}}^{2}}\left(R_{n_{0}}^{0}\right)^{k} \chi_{\widetilde{D}_{n_{0}+1}} \widetilde{S}_{n_{0}}\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}-k-1} f(0)\right| \leq  \tag{5.105}\\
& \kappa_{n_{0}+1} L_{n_{0}}^{-2 \delta+4 a},
\end{align*}
$$

where we used (2.46) with $n=n_{0}$.
Observe that for $z \in B\left(0, \frac{3}{2} \widetilde{D}_{n_{0}+1}\right)$, with (1.49) i) and (5.98),

$$
\sup _{0 \leq k<\ell_{n_{0}}^{2}}\left|\left(R_{n_{0}}^{0}\right)^{\ell_{n_{0}}^{2}-k-1}(f-g)(z)\right| \leq e^{-\kappa_{n_{0}+1}} \text {, with } g(\cdot)=\frac{|\cdot|^{2}}{L_{n_{0}+1}^{2}} \text {. }
$$

Hence with (5.105) we see that when $L_{0}$ is large, for $\omega \in \widetilde{\Omega}$ :

$$
\begin{align*}
& \mid \widetilde{R}_{n_{0}+1} f(0)-d \alpha_{n_{0}}- \\
& \left.\quad \sum_{0 \leq k<\ell_{n_{0}}^{2}} \int P_{\alpha_{n_{0}} k L_{n_{0}}^{2}}(0, d z) \chi_{\widetilde{D}_{n_{0}+1}}(z)\left(\frac{2 \widetilde{d}_{n_{0}}(z, \omega)}{L_{n_{0}+1}^{2}} \cdot z+\sum_{i=1}^{d} \frac{\widetilde{\gamma}_{\frac{1}{i}, i}^{i, i}(z, \omega)}{L_{n_{0}+1}^{2}}\right) \right\rvert\, \leq  \tag{5.106}\\
& \kappa_{n_{0}+1} L_{n_{0}}^{-2 \delta+4 a} .
\end{align*}
$$

In view of (1.24), (1.25), the $\mathbb{P}$-expectation of the sum in (5.106) vanishes. Hence with (5.101) we see that for large $L_{0}$ :

$$
\begin{align*}
& \left|\mathbb{E}\left[\widetilde{\Omega}, \sum_{0 \leq k<\ell_{n_{0}}^{2}} \int P_{\alpha_{n_{0}} k L_{n_{0}}^{2}}(0, d z) \chi_{\widetilde{D}_{n_{0}+1}}(z)\left(\frac{2 \widetilde{d}_{n_{0}}(z, \omega)}{L_{n_{0}+1}^{2}} \cdot z+\sum_{i=1}^{d} \frac{\widetilde{\gamma}_{n_{0}}^{i, i}(z, \omega)}{L_{n_{0}+1}^{2}}\right)\right]\right| \leq  \tag{5.107}\\
& \kappa_{n_{0}} \ell_{n_{0}}^{2} L_{n_{0}+1}^{-10} \leq L_{n_{0}+1}^{-9} .
\end{align*}
$$

So using (5.101), (5.105), (5.107), we see that when $L_{0}$ is large

$$
\begin{aligned}
& d\left|\alpha_{n_{0}+1}-\alpha_{n_{0}}\right| \leq\left|\mathbb{E}\left[\widetilde{R}_{n_{0}+1} f(0)-d \alpha_{n_{0}}, \widetilde{\Omega}^{c}\right]\right|+\mathbb{E}\left[\widetilde{R}_{n_{0}+1} f(0)-d \alpha_{n_{0}}, \widetilde{\Omega}\right] \mid \leq \\
& \kappa_{n_{0}+1} L_{n_{0}}^{-2 \delta+4 a} \stackrel{(1.14),(1.40)}{\leq} L_{n_{0}}^{-\left(1+\frac{9}{10}\right) \delta},
\end{aligned}
$$

and (5.97) is proved.

We can now conclude the proof of Theorem 1.1. We have just shown (1.50) and there remains to complete the proof of (1.47) with $n=n_{0}+1$. With (5.84), we see that when $L_{0}$ is large, for $\sigma \in \Sigma, \omega \in \bar{G}_{\sigma, n_{0}+1}$,

$$
\begin{aligned}
\left\|\chi_{n_{0}+1,0} S_{n_{0}+1}\right\|_{n_{0}+1} & \leq c L_{n_{0}+1}^{-(\delta+a)}+\left\|P_{\alpha_{n_{0}} L_{n_{0}+1}^{2}}-P_{\alpha_{n_{0}+1} L_{n_{0}+1}^{2}}\right\|_{n_{0}+1} \\
& \leq c L_{n_{0}+1}^{-(\delta+a)}+c\left|\alpha_{n_{0}+1}-\alpha_{n_{0}}\right| \stackrel{(5.97)}{\leq} c L_{n_{0}+1}^{-(\delta+a)}
\end{aligned}
$$

using in the second inequality a similar bound as in (4.173). Further with (5.5) we find $\mathbb{P}\left[\left(\bigcup_{\sigma \in \Sigma} \bar{G}_{\sigma, n_{0}+1}\right)^{c}\right] \leq \frac{1}{10} L_{n_{0}+1}^{-M_{0}}$. These bounds together with (2.9) and (2.46) show that

$$
\mathbb{P}\left[0 \notin \mathcal{B}_{n_{0}+1}(\omega)\right] \leq\left(\frac{1}{10}+\frac{1}{10}\right) L_{n_{0}+1}^{-M_{0}} \leq L_{n_{0}+1}^{-M_{0}} .
$$

This concludes the proof of (1.47) for $n=n_{0}+1$, and hence of Theorem 1.1.

## 6 Invariance principle, transience and homogenization

In this section as mentioned in the Introduction, we apply Theorem 1.1 and prove an invariance principle and transience for isotropic diffusions in random environment that are small perturbations of Brownian motion, cf. Theorem 6.3. We also provide an application to homogenization, cf. Theorem 6.4. We begin with a lemma that is helpful when applying Theorem 1.1.

Lemma 6.1. When $L_{0}$ is large, for $\omega \in \Omega, 0 \leq n \leq m_{0}+1$,

$$
\begin{equation*}
\left\|\chi_{n, 0}\left(P_{1, \omega}-P_{1}\right) P_{L_{n}^{2}-1}\right\|_{n} \leq \frac{1}{10} L_{n}^{-\delta} \tag{6.1}
\end{equation*}
$$

$c f .(1.17),(1.38),(1.40),(1.54),(1.57)$ for the notations.
Proof. We recall the convention $L_{-1}=1$, see below (1.15), and extend using this convention the definitions $|\cdot|_{(n)},\|\cdot\|_{n}, \chi_{n, x}$, to the case $n=-1$, cf. (1.28), (1.30), (1.38). We also introduce the probability kernels, see above (1.21) for the notation

$$
\begin{gather*}
\widetilde{P}_{1, \omega}(z, d y)=P_{x, \omega}\left[X_{1 \wedge T_{-1}} \in d y\right], \widetilde{P}_{1}(x, d y)=W_{x}\left[X_{1 \wedge T_{-1}} \in \cdot\right], x \in \mathbb{R}^{d}, \text { where }  \tag{6.2}\\
T_{-1}=\inf \left\{u \geq 0, X_{u}^{*} \geq L_{0}^{\frac{1}{10}}\right\} \tag{6.3}
\end{gather*}
$$

With the same proof as in Proposition 2.5, using exponential inequalities, cf. [23, p. 145], in place of (2.45), we see that for large $L_{0}$, for $\omega \in \Omega, x \in \mathbb{Z}^{d}$,

$$
\begin{equation*}
\left\|\chi_{-1, x}\left(P_{1, \omega}-\widetilde{P}_{1, \omega}\right)\right\|_{-1} \vee\left\|\chi_{-1, x}\left(P_{1}-\widetilde{P}_{1}\right)\right\|_{-1} \leq e^{-c L_{0}^{1 / 10}} \tag{6.4}
\end{equation*}
$$

Hence it follows that for $0 \leq n \leq m_{0}+1$,

$$
\begin{align*}
& \left\|\chi_{-1, x}\left(P_{1, \omega}-P_{1}\right) P_{L_{n}^{2}-1}\right\|_{-1} \leq\left\|\chi_{-1, x}\left(P_{1, \omega}-\widetilde{P}_{1, \omega}\right) P_{L_{n}^{2}-1}\right\|_{-1}+ \\
& \left\|\chi_{-1, x}\left(\widetilde{P}_{1, \omega}-\widetilde{P}_{1}\right) P_{L_{n}^{2}-1}\right\|_{-1}+\left\|\chi_{-1, x}\left(\widetilde{P}_{1}-P_{1}\right) P_{L_{n}^{2}-1}\right\|_{-1} \leq c e^{-c L_{0}^{1 / 10}}+  \tag{6.5}\\
& \left\|\chi_{-1, x}\left(\widetilde{P}_{1, \omega}-\widetilde{P}_{1}\right) P_{L_{n}^{2}-1}\right\|_{-1} .
\end{align*}
$$

With a similar argument as in (5.94), for $0 \leq n \leq m_{0}+1$, and $f$ with $|f|_{(n)} \leq 1$, we can construct with a cut-off function, a function $\widetilde{H}$ supported in $B\left(x, 3 L_{0}^{1 / 10}\right)$, such that:

$$
\begin{equation*}
\widetilde{H} \text { agrees with } P_{L_{n}^{2}-1} f-P_{L_{n}^{2}-1} f(x) \text { in } B\left(x, 2 L_{0}^{1 / 10}\right) \text { and }|\widetilde{H}|_{(-1)} \leq c \frac{L_{0}^{\frac{1}{10}}}{L_{n}} \tag{6.6}
\end{equation*}
$$

where (1.56) has been used for the last inequality. We hence find that with large $L_{0}$

$$
\left|\chi_{-1, x}\left(\widetilde{P}_{1, \omega}-\widetilde{P}_{1}\right) P_{L_{n}^{2}-1} f\right|_{(-1)}=\left|\chi_{-1, x}\left(\widetilde{P}_{1, \omega}-\widetilde{P}_{1}\right) H\right|_{(-1)} \stackrel{(1.62),(6.4),(6.6)}{\leq} c L_{0}^{\frac{1}{10}} L_{n}^{-1},
$$

and hence with (6.5), (6.6):

$$
\begin{align*}
& \left\|\chi_{n, 0}\left(P_{1, \omega}-P_{1}\right) P_{L_{n}^{2}-1}\right\|_{n} \leq L_{n}^{\beta}\left\|\chi_{n, 0}\left(P_{1, \omega}-P_{1}\right) P_{L_{n}^{2}-1}\right\|_{-1} \\
& \leq L_{n}^{\beta}\left(c e^{-c L_{0}^{\frac{1}{10}}}+c L_{0}^{\frac{1}{10}} L_{n}^{-1}\right) \stackrel{(1.17)}{\leq} \frac{1}{10} L_{n}^{-\delta} . \tag{6.7}
\end{align*}
$$

This proves our claim.
The next proposition is instrumental and enables to construct good couplings of the diffusion in random environment with Brownian motion. From now on we specify the choices of $\nu=2, \beta=\frac{1}{2}, a, c_{0}, \varphi, \psi, \zeta, M_{0}, M$, cf. (1.5), (1.13), (1.14), (1.32), (1.43), (1.46). In accordance with the convention concerning constants started above Theorem 1.1, constants will solely depend on $d, K, R$ in view of the choices we just made. We denote with $\widetilde{X}_{t}, t \geq 0$, and $\widetilde{X}_{t}^{0}, t \geq 0$, the canonical processes on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)^{2}$, the space on which we will construct the coupling measures.

Proposition 6.2. ( $d \geq 3$ )
Given $K>1, R>0$, there exists $\eta_{0}>0$, depending only on $d, K, R$, such that for $a(x, \omega), b(x, \omega)$ as in (1.2), satisfying (1.4), (1.7), (0.4), and

$$
\begin{equation*}
|a(x, \omega)-I| \leq \eta_{0},|b(x, \omega)| \leq \eta_{0}, \text { for } x \in \mathbb{R}^{d}, \omega \in \Omega \tag{6.8}
\end{equation*}
$$

then there is an event $\bar{\Omega}$ with full $\mathbb{P}$-measure and a finite $N(\cdot)$ on $\bar{\Omega}$, such that for $\omega \in \bar{\Omega}$, when $n \geq N(\omega)$ :

$$
\begin{equation*}
\text { for all } x \in L_{n} \mathbb{Z}^{d} \cap\left(4 \mathcal{T}_{n+3}\right), x \in \mathcal{B}_{n}(\omega) \text {, (cf. (1.32), (2.1) for notations), } \tag{6.9}
\end{equation*}
$$

and for any $y \in \mathbb{R}^{d}$ there is a coupling measure $\widetilde{Q}_{n, y, \omega}$ on $C\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right)^{2}$ such that under $\widetilde{Q}_{n, y, \omega}$,

$$
\begin{align*}
& \widetilde{X}_{.}^{0} \text { is distributed as } X_{\alpha_{n}} . \text { under } W_{y} \text {, }  \tag{6.10}\\
& \widetilde{X}_{. \wedge T_{2 \tau_{n+3}}(\widetilde{X})} \text { is distributed as } X_{. \wedge T_{2 \tau_{n+3}} \text { under } P_{y, \omega} \text {, and, }}^{\widetilde{Q}_{n, y, \omega}\left[\sup _{u \leq L_{n+3}^{2}}\left|\widetilde{X}_{u}-\widetilde{X}_{u}^{0}\right| \geq 3 \widetilde{D}_{n}\right] \leq L_{n}^{-\delta / 2} .} \tag{6.11}
\end{align*}
$$

Proof. In the sequel we use the expression "small enough $\eta_{0}$ ", in place $\eta_{0} \leq c$, with $c$ a constant, with the meaning explained above Proposition 6.2. From now on we assume $\eta_{0}<1$ small enough so that (1.3), (1.5) are satisfied. We now choose constants $L_{0}$ and $c_{2}$ according to Theorem 1.1, Lemma 6.1, and such that for all $n \geq 0$, (recall $W_{0}$ denotes the Wiener measure)
i) if in (2.45), $\kappa_{n}^{0}=\frac{1}{2}\left(\widetilde{D}_{n} / D_{n}\right)$, then $e^{-\kappa_{n}}$ in (2.46) is smaller than $\frac{1}{10} L_{n}^{-\delta}$,
ii) $W_{0}\left[X_{L_{n}^{2}}^{*} \geq v\right] \leq \frac{1}{10} \exp \left\{-\frac{4 v}{D_{n}}\right\}$, for $v \geq \frac{1}{4} D_{n}$,
iii) $\quad\left(E^{W_{0}}\left[\left|X_{L_{n}^{2}}\right|^{4}\right]^{\frac{1}{2}}+\widetilde{D}_{n}^{2}\right) W_{0}\left[X_{L_{n}^{2}}^{*}>\frac{\widetilde{D}_{n}}{4}\right]^{\frac{1}{2}} \leq \frac{1}{100}$,
iv) $\left|\chi_{n, 0}\right|_{(n)} \sup _{\frac{1}{2} \leq \alpha \neq \alpha^{\prime} \leq 4} \frac{\left\|P_{\alpha L_{n}^{2}}-P_{\alpha^{\prime} L_{n}^{2}}\right\|_{n}}{\left|\alpha-\alpha^{\prime}\right|} \leq L_{0}$, cf. (4.173),
and

$$
\begin{equation*}
\sum_{n \geq 0} L_{n}^{-\left(1+\frac{9}{10}\right) \delta}<\frac{1}{10} . \tag{6.14}
\end{equation*}
$$

We have now specified $L_{0}$, and we will first see that:

$$
\begin{align*}
& \text { for } \eta_{0} \text { small enough, }(1.47),(1.48),(1.49) \text { hold for all } \\
& n_{0} \geq m_{0}+1, \text { and }\left|\alpha_{0}-1\right|<\frac{1}{10} \tag{6.15}
\end{align*}
$$

To this end, first recall from (1.9) that for $\omega \in \Omega, x \in \mathbb{R}^{d}$, there is an $\left(\mathcal{F}_{t}\right)$-Brownian motion $\beta$. such that $P_{x, \omega}$-a.s., for all $t \geq 0$,

$$
\begin{equation*}
X_{t}=x+\int_{0}^{t} \sigma\left(X_{s}, \omega\right) d \beta_{s}+\int_{0}^{t} b\left(X_{s}, \omega\right) d s, \text { with } \sigma(\cdot, \omega)=a(\cdot, \omega)^{\frac{1}{2}} \tag{6.16}
\end{equation*}
$$

Note that for $y \in \mathbb{R}^{d}, \omega \in \Omega, \sigma(y, \omega)-I=(a(y, \omega)-I)(\sigma(y, \omega)+I)^{-1}$, so for small $\eta_{0}$, $y \in \mathbb{R}^{d}, \omega \in \Omega$, with (6.8),

$$
\begin{equation*}
|\sigma(y, \omega)-I| \leq c \eta_{0} \tag{6.17}
\end{equation*}
$$

Further from the exponential martingale inequalities, cf. [23], p. 145,

$$
\begin{align*}
& P_{x, \omega}\left[\sup _{v \leq t}\left|\int_{0}^{v} \sigma\left(X_{s}, \omega\right) d \beta_{s}-\beta_{v}\right| \geq u\right] \leq c \exp \left\{-\frac{c u^{2}}{\eta_{0}^{2} t}\right\}  \tag{6.18}\\
& \text { for } u, t>0, x \in \mathbb{R}^{d}, \omega \in \Omega
\end{align*}
$$

Choosing $\eta_{0}$ small, with (6.8), (6.16), (6.18), we see that for $\omega \in \Omega, 0 \leq n \leq m_{0}+1$, $x \in L_{n} Z^{d}, A \subseteq C_{n}(x), \gamma \in\left\{1, \ldots, 2 d 5^{d-1}\right\}$, and the notation (1.44),

$$
\begin{equation*}
J_{n, x, A, \gamma}(\omega)=0 \tag{6.19}
\end{equation*}
$$

so that (1.48) holds for $0 \leq n \leq m_{0}+1$. Likewise with (6.13) ii), we see that choosing $\eta_{0}$ small we can make sure that for $\omega \in \Omega, 0 \leq n \leq m_{0}+1, y \in \mathbb{R}^{d}$,

$$
\begin{equation*}
P_{y, \omega}\left[X_{L_{n}^{2}}^{*} \geq v\right] \leq \exp \left\{-\frac{v}{D_{n}}\right\}, \text { for all } v \geq D_{n} \tag{6.20}
\end{equation*}
$$

Further we have

$$
\chi_{n, 0}\left(R_{n}-P_{L_{n}^{2}}\right)=\chi_{n, 0} P_{1, \omega}\left(P_{L_{n}^{2}-1, \omega}-P_{L_{n}^{2}-1}\right)+\chi_{n, 0}\left(P_{1, \omega}-P_{1}\right) P_{L_{n}^{2}-1},
$$

and with (1.60), (6.1), (6.18), it follows that choosing $\eta_{0}$ small, for $\omega \in \Omega$, and $0 \leq n \leq$ $m_{0}+1$,

$$
\begin{equation*}
\left\|\chi_{n, 0}\left(R_{n}-P_{L_{n}^{2}}\right)\right\|_{n} \leq \frac{1}{5} L_{n}^{-\delta} \tag{6.21}
\end{equation*}
$$

Recall that, cf. (1.22)

$$
\alpha_{n}=\frac{1}{d L_{n}^{2}} E_{0}\left[\left|X_{L_{n}^{2} \wedge T_{n}}\right|^{2}\right],
$$

and note that for small $\eta_{0}$, with (6.13) iii), (6.18), for $0 \leq n \leq m_{0}+1$,

$$
\begin{aligned}
& \left|E_{0}\left[\left|X_{L_{n}^{2}}\right|^{2}\right]-E_{0}\left[\left|X_{L_{n}^{2} \wedge T_{n}}\right|^{2}\right]\right| \leq E_{0}\left[\left(\left|X_{L_{n}^{2}}\right|^{2}+\widetilde{D}_{n}^{2}\right), T_{n}<L_{n}^{2}\right] \\
& \leq\left(E_{0}\left[\left|X_{L_{n}^{2}}\right|^{4}\right]^{\frac{1}{2}}+\widetilde{D}_{n}^{2}\right)\left(P_{0}\left[\sup _{s \leq L_{n}^{2}}\left|\beta_{s}\right| \geq \frac{\widetilde{D}_{n}}{4}\right]^{\frac{1}{2}}+P_{0}\left[\sup _{0 \leq s \leq L_{n}^{2}}\left|\int_{0}^{s}\left(\sigma\left(X_{s}, \omega\right)-I\right) d \beta_{s}\right| \geq \frac{\widetilde{D}_{n}}{4}\right]^{\frac{1}{2}}\right) \\
& \leq \frac{1}{20}
\end{aligned}
$$

So when $\eta_{0}$ is small enough, for $0 \leq n \leq m_{0}+1$,

$$
\begin{equation*}
\left|\alpha_{n}-1\right| \leq \frac{1}{20 d L_{n}^{2}}+\frac{1}{d L_{n}^{2}} \left\lvert\, E_{0}\left[\left|X_{L_{n}^{2}}\right|^{2}\right]-E_{0}\left[\left|\beta_{L_{n}^{2}}\right|^{2}\right] \stackrel{(6.16),(6.18)}{\leq} \frac{1}{10 L_{n}^{2}}\right. \tag{6.22}
\end{equation*}
$$

and hence

$$
\begin{array}{ll}
\text { i) } & \left|\alpha_{n}-\alpha_{n+1}\right| \leq L_{n}^{-\left(1+\frac{9}{10}\right) \delta}, 0 \leq n \leq m_{0} \text {, and } \\
\text { ii) } & \alpha_{n} \in\left[\frac{1}{4}, 4\right]\left(=\left[\frac{1}{2 \nu}, 2 \nu\right]\right), \text { for } 0 \leq n \leq m_{0}+1 \tag{6.23}
\end{array}
$$

This proves that (1.49) holds for $0 \leq n \leq m_{0}+1$. Then observe that for $0 \leq n \leq m_{0}+1$, $\omega \in \Omega$,

$$
\left\|\chi_{n, 0} \widetilde{S}_{n}\right\|_{n} \leq\left\|\chi_{n, 0}\left(\widetilde{S}_{n}-S_{n}\right)\right\|_{n}+\left\|\chi_{n, 0}\left(R^{n}-P_{L_{n}^{2}}\right)\right\|_{n}+\left\|\chi_{n, 0}\left(P_{\alpha_{n} L_{n}^{2}}-P_{L_{n}^{2}}\right)\right\|_{n}
$$

so that using (6.20), (2.46), (6.13) i) to bound the first term in the right-hand side, (6.21) to bound the second term, (6.13) iv), (6.22), (6.23) ii) to bound the last term, we see that when $\eta_{0}$ is small, for $\omega \in \Omega, 0 \leq n \leq m_{0}+1$,

$$
\begin{equation*}
\left\|\chi_{n, 0} \widetilde{S}_{n}\right\|_{n} \leq \frac{1}{10} L_{n}^{-\delta}+\frac{1}{5} L_{n}^{-\delta}+\frac{1}{5} L_{0} L_{n}^{-2} \leq L_{n}^{-\delta} \tag{6.24}
\end{equation*}
$$

Hence with (6.20), we see that for small $\eta_{0}$, when $\omega \in \Omega, 0 \leq n \leq m_{0}+1$,

$$
\begin{equation*}
0 \in \mathcal{B}_{n}(\omega) \tag{6.25}
\end{equation*}
$$

We can now apply Theorem 1.1, and with (6.14) note that $\left|\alpha_{0}-1\right|<\frac{1}{10}$ implies that (1.49) remains also satisfied by induction, so that (6.15) is proved.

As a next step observe that for $n \geq m_{0}+1$,

$$
\begin{aligned}
& \mathbb{P}\left[\text { for some } x \in L_{n} \mathbb{Z}^{d} \cap\left(4 \mathcal{T}_{n+3}\right), x \notin \mathcal{B}_{n}(\omega)\right] \leq \\
& c\left(\frac{L_{n+3}^{2}}{L_{n}}\right)^{d} L_{n}^{-M_{0}} \stackrel{(1.46)}{\leq} c L_{n}^{2 d(1+a)^{3}-100 d(1+a)^{m_{0}+2}} \leq c L_{n}^{-98 d},
\end{aligned}
$$

and this last quantity is the general term of a convergent series. With Borel-Cantelli's lemma, we see that there is an event $\bar{\Omega}$ with full $\mathbb{P}$-measure, and a finite $N(\cdot)$ on $\bar{\Omega}$, such that when $n \geq N(\omega)$, (6.9) holds.

Let us now fix $\omega \in \bar{\Omega}$. Given $n \geq N(\omega)$, we denote with $h$ some $[0,1]$-valued continuous function with value 1 on $2 \mathcal{T}_{n+3}$, and 0 on $\left(3 \mathcal{T}_{n+3}\right)^{c}$. With Proposition 3.1, we have for $y \in \mathbb{R}^{d}$ a coupling measure $Q_{n, y}$ on $\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right)^{\mathrm{N}}$, under which the canonical processes $\bar{X}_{k}, k \geq 0$, and $\bar{X}_{k}^{0}, k \geqq 0$, have the laws of the Markov chains on $\mathbb{R}^{d}$ starting at $y$ with respective transitions $\widetilde{R}_{n, h}$, cf. (3.4), and $R_{n}^{0}$. Using product of bridge measures in time $L_{n}^{2}$ between $z, z^{\prime} \in \mathbb{R}^{d}$

$$
\begin{aligned}
Q_{z, z^{\prime}}^{L_{n}^{2}}= & \left(h(x) p_{L_{n}^{2}, \omega}\left(z, z^{\prime}\right)+(1-h(x)) p_{\alpha_{n} L_{n}^{2}}\left(z, z^{\prime}\right)\right)^{-1}\left(h(x) p_{L_{n}^{2}}\left(z, z^{\prime}\right) P_{z, z^{\prime}, \omega}^{L_{n}^{2}}\right. \\
& \left.+(1-h(x)) p_{\alpha_{n} L_{n}^{2}}\left(z, z^{\prime}\right) P_{z, z^{\prime}}^{L_{n}^{2}}\right)
\end{aligned}
$$

with $P_{z, z^{\prime}, \omega}^{L_{n}^{2}}$ and $P_{z, z^{\prime}}^{L_{n}^{2}}$ respectively denoting the bridge measures in time $L_{n}^{2}$ between $z, z^{\prime}$ for the diffusion in the environment $\omega$, and for Brownian motion with covariance $\alpha_{n} I$, to interpolate between times $k L_{n}^{2}$ and $(k+1) L_{n}^{2}$, for the first canonical process and product of bridge measures $P_{z, z^{\prime}}^{L_{n}^{2}}$ to interpolate between times $k L_{n}^{2}$ and $(k+1) L_{n}^{2}$ for the second canonical process (so that conditional on $\bar{X}_{k}, k \geq 0, \bar{X}_{k}^{0}, k \geq 0$, all these bridges are independent) we can construct a coupling measure $\widetilde{Q}_{n, y, \omega}$ on $C^{2}\left(\mathbb{R}_{+}, \mathbb{R}^{d}\right), y \in \mathbb{R}^{d}$, so that (6.10), (6.11) holds. Then using (3.6), (6.9), (1.39), we find:

$$
\begin{align*}
& \widetilde{Q}_{n, y, \omega}\left[\sup _{u \leq L_{n+3}^{2}}\left|\widetilde{X}_{u}-\widetilde{X}_{u}^{0}\right| \geq 3 \widetilde{D}_{n}\right] \leq \\
& \left(\frac{L_{n+3}}{L_{n}}\right)^{4}\left(\kappa_{n} L_{n}^{-\delta}+e^{-\kappa_{n}}\right)+2\left(\frac{L_{n+3}}{L_{n}}\right)^{2} e^{-\kappa_{n}} \leq L_{n}^{-\delta / 2} \tag{6.26}
\end{align*}
$$

when $n$ is large enough. Hence increasing $N(\cdot)$ if necessary, we see that for $\omega \in \bar{\Omega},(6.10)$, (6.11), (6.12) holds, and this finishes the proof of Proposition 6.2.

We are now ready to state and prove our main applications.
Theorem 6.3. $(d \geq 3)$
With $\eta_{0}(d, K, R)>0$, as in Proposition 6.2, when $a(x, \omega), b(x, \omega)$, as in (1.2), satisfy (1.4), (1.7), (0.4) as well as (6.8), i.e.

$$
|a(x, \omega)-I| \leq \eta_{0},|b(x, \omega)| \leq \eta_{0}, \text { for } x \in \mathbb{R}^{d}, \omega \in \Omega
$$

then $\mathbb{P}$-a.s.,

$$
\begin{equation*}
\text { for all } x \in \mathbb{R}^{d}, P_{x, \omega} \text {-a.s., } \lim _{t \rightarrow \infty}\left|X_{t}\right|=\infty \tag{6.28}
\end{equation*}
$$

Proof. We keep the notations of Proposition 6.2. We first prove (6.27). From (6.23) we know that $\alpha_{n}$ converges and we write

$$
\begin{equation*}
\sigma^{2} \stackrel{\text { def }}{=} \lim _{n} \alpha_{n}\left(\in\left[\frac{1}{4}, 4\right]\right) . \tag{6.29}
\end{equation*}
$$

The claim (6.27) will follow once we prove that for any $\omega$ in $\bar{\Omega}$, in the notations of Proposition 6.2,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} E_{0, \omega}\left[F\left(\frac{1}{\sqrt{t}} X_{. t}\right)\right]=E^{W_{0}}\left[F\left(X_{\sigma^{2}} .\right)\right] \tag{6.30}
\end{equation*}
$$

for any $F$ on $C\left([0, T], \mathbb{R}^{d}\right), T>0$, bounded by 1 , Lipschitz relative to the distance function

$$
\begin{equation*}
D_{T}\left(w, w^{\prime}\right)=\sup _{s \leq T}\left|w(s)-w^{\prime}(s)\right| \wedge 1, w, w^{\prime} \in C\left([0, T], \mathbb{R}^{d}\right) \tag{6.31}
\end{equation*}
$$

with Lipschitz constant 1 , with a slight abuse of notations in (6.30). For $t$ large we define the integer $n(t) \geq 0$, such that

$$
\begin{equation*}
L_{n(t)+1}^{2} \leq t<L_{n(t)+2}^{2} \tag{6.32}
\end{equation*}
$$

and observe that for $\omega \in \bar{\Omega}, F$ as above and large $t$

$$
\begin{align*}
& \left|E_{0, \omega}\left[F\left(\frac{1}{\sqrt{t}} X_{\cdot t}\right)\right]-E^{W_{0}}\left[F\left(X_{\sigma^{2} \cdot}\right)\right]\right| \leq a_{1}+a_{2}+a_{3}, \text { where } \\
& a_{1}(t)=\left|E_{0, \omega}\left[F\left(\frac{1}{\sqrt{t}} X_{\cdot t}\right)\right]-E_{0, \omega}\left[F\left(\frac{1}{\sqrt{t}} X_{(\cdot t) \wedge T_{2 \tau_{n(t)+3}}}\right)\right]\right|  \tag{6.33}\\
& a_{2}(t)=\left|E_{0, \omega}\left[F\left(\frac{1}{\sqrt{t}} X_{(\cdot t)} \wedge T_{2 \tau_{n(t)+3}}\right)\right]-E^{W_{0}}\left[F\left(\frac{1}{\sqrt{t}} X_{\alpha_{n(t)}} \cdot\right)\right]\right|, \\
& a_{3}(t)=\left|E^{W_{0}}\left[F\left(\sqrt{\alpha}_{n(\cdot)} X .\right)\right]-E^{W_{0}}[F(\sigma X .)]\right|
\end{align*}
$$

and we have used Brownian scaling for $a_{3}(\cdot)$. From (6.29) and dominated convergence, we see that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} a_{3}(t)=0 \tag{6.34}
\end{equation*}
$$

Further when $t$ is large,

$$
\begin{gather*}
a_{1}(t) \leq 2 P_{0, \omega}\left[T_{2} \mathcal{T}_{n(t)+3}<T t\right] \stackrel{(6.32)}{\leq} 2 P_{0, \omega}\left[T_{2 \mathcal{T}_{n(t)+3}}<T L_{n(t)+2}^{2}\right] \\
\stackrel{(2.10)}{\leq} c \exp \left\{-c L_{n(t)+3}^{2}\right\}, \text { so that } \\
\lim _{t \rightarrow \infty} a_{1}(t)=0 . \tag{6.35}
\end{gather*}
$$

As for $a_{2}(t)$, using the coupling measure $\widetilde{Q}_{n(t), 0, \omega}$ from Proposition 6.2 , we find with (6.10), (6.11), that for large $t$

$$
\begin{aligned}
& a_{2}(t)=\left\lvert\, E^{\widetilde{Q}_{n(t), 0, \omega}}\left[\left.F\left(\frac{1}{\sqrt{t}} \widetilde{X}_{\left.(\cdot t) \wedge T_{2 T_{n(t)+3}(\tilde{X})}\right)}-F\left(\frac{1}{\sqrt{t}} \widetilde{X}_{\cdot t}^{0}\right)\right] \right\rvert\, \leq\right.\right. \\
& \quad E^{\widetilde{Q}_{n(t), 0, \omega}}\left[\sup _{u \leq T t} \frac{\left|\widetilde{X}_{u \wedge T_{2} \tau_{n(t)+3}(\tilde{X})}-\widetilde{X}_{u}^{0}\right|}{\sqrt{t}} \wedge 1\right] \underset{(2.10)}{\leq} \frac{(6.12),(6.32)}{\leq} \frac{3 \widetilde{D}_{n(t)}}{\sqrt{t}}+L_{n(t)}^{-\delta / 2}+c e^{-c L_{n(t)+3}^{2}},
\end{aligned}
$$

so that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} a_{2}(t)=0 \tag{6.36}
\end{equation*}
$$

Combining (6.34) - (6.36), the claim (6.30) follows. This proves (6.27).
We now prove (6.28). When $n$ is large, it follows from standard estimates on Brownian motion and (1.49) that for $|z|=L_{n+1}$,

$$
\begin{equation*}
W_{z}\left[X_{\alpha_{n}} . \text { exits } B\left(0,2 L_{n+2}\right) \text { before time } L_{n+3}^{2} \text { or entering } \bar{B}\left(0,4 \widetilde{D}_{n}\right)\right] \geq 1-\frac{\kappa_{n}}{\ell_{n}} \tag{6.37}
\end{equation*}
$$

Then for $\omega \in \bar{\Omega}$, with Proposition 6.2 and (6.37) we see that for large $n$ and $|z|=L_{n+1}$,

$$
\begin{equation*}
\widetilde{Q}_{n, z, \omega}\left[\tilde{X} . \text { enters } \bar{B}\left(0, L_{n}\right) \text { before exiting } B\left(0, L_{n+2}\right)\right] \leq L_{n}^{-\delta / 2}+\frac{\kappa_{n}}{\ell_{n}} \leq \frac{\kappa_{n}}{\ell_{n}} \tag{6.38}
\end{equation*}
$$

With (6.11), we thus see that for large $n$ and $|z|=L_{n+1}$,

$$
P_{z, \omega}\left[H_{\bar{B}\left(0, L_{n}\right)}<T_{B\left(0, L_{n+2}\right)}\right] \leq \frac{\kappa_{n}}{\ell_{n}} \leq \ell_{n}^{-1 / 2}
$$

so that with the strong Markov property we find:

$$
\begin{equation*}
P_{z, \omega}\left[H_{\bar{B}\left(0, L_{n}\right)}=\infty\right] \geq \prod_{k \geq 0}\left(1-\ell_{n+k}^{-1 / 2}\right) \underset{n \rightarrow \infty}{\longrightarrow} 1 \tag{6.39}
\end{equation*}
$$

It now follows in a standard way that when $\omega \in \bar{\Omega}$,

$$
\begin{equation*}
\text { for } x \in \mathbb{R}^{d}, P_{x, \omega}\left[\lim _{t \rightarrow \infty}\left|X_{t}\right|=\infty\right]=1 \tag{6.40}
\end{equation*}
$$

and this proves (6.28).
We conclude this section with an application to homogenization in random media. Given $f, g$ bounded functions on $\mathbb{R}^{d}$ respectively continuous and Hölder continuous, under the assumptions of Theorem 6.3, for $\omega \in \Omega$ and $\epsilon>0$, there is a unique bounded solution of the Cauchy problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{\epsilon}=L_{\epsilon} u_{\epsilon}+g \text { in }(0, \infty) \times \mathbb{R}^{d}  \tag{6.41}\\
\left.u_{\epsilon}\right|_{t=0}=f
\end{array}\right.
$$

where

$$
\begin{equation*}
L_{\epsilon}=\frac{1}{2} \sum_{i, j=1}^{d} a_{i j}\left(\frac{x}{\epsilon}, \omega\right) \partial_{i j}^{2}+\sum_{i=1}^{d} \frac{1}{\epsilon} b_{i}\left(\frac{x}{\epsilon}, \omega\right) \partial_{i} \tag{6.42}
\end{equation*}
$$

see for instance [9, Theorem 12, p. 25], and [10, Theorem 5.3]. The asymptotic behavior of $u_{\epsilon}$, as $\epsilon \rightarrow 0$, is intimately related to the invariance principle proved in Theorem 6.3.

Theorem 6.4. ( $d \geq 3$ )
Under the same assumptions as in Theorem 6.3, on a set of full $\mathbb{P}$-measure, for any $f, g$ as above, the solution $u_{\epsilon}$ of (6.41) converges uniformly on compact subsets of $\mathbb{R}_{+} \times \mathbb{R}^{d}$ to the solution of the Cauchy problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{0}=\sigma^{2} \Delta u_{0}+g \text { in }(0, \infty) \times \mathbb{R}^{d}  \tag{6.43}\\
\left.u_{\epsilon}\right|_{t=0}=f
\end{array}\right.
$$

with $\sigma^{2}$ as in (6.27).

Proof. Consider $\omega \in \bar{\Omega}$, (cf. Proposition 6.2), and $\epsilon>0$, with [10, Theorem 5.3], we can write

$$
\begin{equation*}
u_{\epsilon}(s, x)=E_{x / \epsilon, \omega}\left[f\left(\epsilon X_{s / \epsilon^{2}}\right)-\int_{0}^{s} g\left(\epsilon X_{v / \epsilon^{2}}\right) d v\right] \text {, for } s \geq 0, x \in \mathbb{R}^{d} \tag{6.44}
\end{equation*}
$$

Letting $\epsilon^{-1}$ play the role of $t$ in (6.32), very similar bounds as in (6.33) - (6.36), with some obvious modifications for the bound above (6.36) yield that as $\epsilon \rightarrow 0$,

$$
\begin{align*}
& u_{\epsilon} \text { converges uniformly on compact subsets of } \mathbb{R}_{+} \times \mathbb{R}^{d} \text { to } \\
& u_{0}(s, x)=E^{W_{x}}\left[f\left(X_{\sigma^{2} s}\right)-\int_{0}^{s} g\left(X_{\sigma^{2} v}\right) d v\right] \tag{6.45}
\end{align*}
$$

and our claim now follows.
The proofs of the last two theorems illustrate the fact that the measures constructed in Proposition 6.2 offer a very quantitative and handy comparison of the isotropic diffusion in random environment with Brownian motion.

## A Appendix

This appendix collects several results concerning the Hölder-norms $|\cdot|_{(n)},\|\cdot\|_{n}$, cf. (1.28), (1.30). In particular the effective control of these norms with the help of wavelets is discussed in Proposition A.2. We begin with the convenient

Lemma A.1. ( $n \geq 0, L_{n}$ as in (1.15), $\beta \in(0,1)$ )
Consider a non-empty index set $I, f,\left(g_{i}\right)_{i \in I}$, scalar functions on $\mathbb{R}^{d}$, $\left(x_{i}\right)_{i \in I}$, points of $\mathbb{R}^{d}$, such that

$$
\begin{align*}
& f=g_{i}, \text { on } B\left(x_{i}, 2 L_{n}\right), i \in I, \text { and }  \tag{A.1}\\
& \operatorname{Supp} f \subseteq \bigcup_{i \in I} \bar{B}\left(x_{i}, L_{n}\right), \text { then }  \tag{A.2}\\
& |f|_{(n)} \leq 3 \sup _{i \in I}\left|g_{i}\right|_{(n)} . \tag{A.3}
\end{align*}
$$

Moreover if $f$ is a scalar function, $\Gamma>0$, and

$$
\begin{align*}
& \sup _{x \in \mathbb{R}^{d}}|f(x)| \leq \Gamma  \tag{A.4}\\
& |f(x)-f(y)| \leq \Gamma\left|\frac{x-y}{L_{n}}\right|^{\beta}, \text { for } x, y \text { in the open } L_{n} \text {-neighborhood }  \tag{A.5}\\
& \text { of the support of } f \text { and }|x-y|<L_{n} \text {, }
\end{align*}
$$

then

$$
\begin{equation*}
|f|_{(n)} \leq 3 \Gamma \tag{A.6}
\end{equation*}
$$

Proof. We first prove (A.3). Note that

$$
|f|_{\infty} \leq \sup _{i \in I}\left|g_{i}\right|_{\infty}
$$

and for $x, y$ in $\mathbb{R}^{d}$ with $|x-y| \geq L_{n}$,

$$
L_{n}^{\beta} \frac{|f(x)-f(y)|}{|x-y|^{\beta}} \leq 2 \sup _{i}\left|g_{i}\right|_{\infty} .
$$

On the other hand, when $x, y$ are distinct points of $\mathbb{R}^{d}$, with $|x-y|<L_{n}$ and say $x \in \operatorname{Supp} f$, then $x \in \bar{B}\left(x_{i_{0}}, L_{n}\right)$, for some $i_{0} \in I$. One then has

$$
L_{n}^{\beta} \frac{|f(x)-f(y)|}{|x-y|^{\beta}} \stackrel{(A .1)}{=} L_{n}^{\beta} \frac{\left|g_{i_{0}}(x)-g_{i_{0}}(y)\right|}{|x-y|^{\beta}},
$$

whereas when none of $x, y$ belongs to Supp $f$, the left member vanishes. The claim (A.3) now follows.

We now prove (A.6). Note that when $x, y$ are such that $|x-y| \geq L_{n}$, then

$$
L_{n}^{\beta} \frac{|f(x)-f(y)|}{|x-y|^{\beta}} \leq 2|f|_{\infty} \stackrel{(A .4)}{\leq} 2 \Gamma .
$$

On the other hand when $x, y$ are distinct points of $\mathbb{R}^{d}$ with $|x-y|<L_{n}$, and either some or none of them belongs to $\operatorname{Supp} f$, we find with (A.5)

$$
L_{n}^{\beta} \frac{|f(x)-f(y)|}{|x-y|^{\beta}} \leq \Gamma,
$$

and the claim (A.6) now follows.
The next result will provide an effective control of the Hölder-norms (1.28), (1.30), with the help of the expansion in an orthonormal basis of wavelets. The fact that such bases give rise to a handy control of the Hölder-property is well known, cf. Daubechies [6, p. 199-203], Mallat [16, p. 169-173]. The proposition we will now prove, gives a version of these results useful for the calculations of Section 4. We introduce the sequence of non-negative integers $J_{n}, n \geq 0$, such that

$$
\begin{equation*}
2^{J_{n}} \leq L_{n}<2^{J_{n+1}} \tag{A.7}
\end{equation*}
$$

and recall the $L^{2}\left(\mathbb{R}^{d}\right)$-orthogonal expansion in (1.35).
Proposition A.2. $(d \geq 1,0<\beta<1, \varphi, \psi)$
There is a constant $\Gamma>1$, depending on $d, \beta, \varphi, \psi$, such that for $n \geq 0$, and $f$ compactly supported bounded measurable function, one has, cf. (1.35) for the notations,

$$
\begin{equation*}
\frac{1}{\Gamma}|f|_{(n)} \leq \sup _{\substack{\alpha, \ell \leq J_{n}, p \in \mathbb{Z}^{d} \\ \alpha \neq 0, \text { for } \ell<J_{n}}} 2^{\beta\left(J_{n}-\ell\right)}\left|c_{\alpha, \ell, p}^{J_{n}}\right| \leq \Gamma|f|_{(n)} . \tag{A.8}
\end{equation*}
$$

Moreover, when $A$ is a bounded linear operator mapping bounded measurable functions on $\mathbb{R}^{d}$ into bounded measurable compactly supported functions on $\mathbb{R}^{d}$, and $A$ vanishes for functions supported in the complement of some compact subset of $\mathbb{R}^{d}$, then

$$
\begin{equation*}
\frac{1}{\Gamma}\|A\|_{n} \leq \sup _{\substack{\alpha, \ell \leq J_{n}, p \in \mathbb{Z}^{d} \\ \alpha \neq 0, \text { when } \ell<J_{n}}} \sum_{\substack{\alpha^{\prime} \ell^{\prime} \leq J_{n}, p^{\prime} \in \mathbb{Z}^{d} \\ \alpha^{\prime} \neq 0, \text { when } \ell^{\prime}<J_{n}}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \frac{1}{2^{d \ell}}\left|\left\langle\theta_{\alpha, \ell, p}, A \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right\rangle\right| \leq \Gamma\|A\|_{n}, \tag{A.9}
\end{equation*}
$$

with the notation $\langle h, g\rangle=\int h(x) g(x) d x$.

Proof. We begin with the proof of (A.8). For $f$ as in the statement, $\alpha \in\{0,1\}^{d}, \ell \leq J_{n}, p \in \mathbb{Z}^{d}$, with $\alpha \neq 0$, when $\ell<J_{n}$, the coefficients $c_{\alpha, \ell, p}^{J_{n}}$ of (A.8), are expressed in view of (1.35), as

$$
\begin{equation*}
c_{\alpha, \ell, p}^{J_{n}}=\frac{1}{2^{\ell d}} \int_{\mathbb{R}^{d}} f(x) \theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right) d x, \tag{A.10}
\end{equation*}
$$

(note incidentally that for $\left.n \geq 0, \ell \leq J_{n}, \alpha \neq 0, c_{\alpha, \ell, p}^{J_{n}}=c_{\alpha, \ell, p}^{J_{n+1}}\right)$. Denoting throughout the proof with $c$ a positive constant changing from place to place and solely depending on $d, \beta, \varphi, \psi$, we find that for $\ell \leq J_{n}, p \in \mathbb{Z}^{d}, \alpha \in\{0,1\}^{d}$, with $\alpha \neq 0$ if $\ell<J_{n}$ :

$$
\begin{equation*}
\left|\left.\right|_{\alpha, \ell, p} ^{J_{n}}\right| \leq c|f|_{\infty} \leq c|f|_{(n)} . \tag{A.11}
\end{equation*}
$$

Note that when $\alpha \neq 0, \theta_{\alpha_{i}}=\psi$, for some $1 \leq i \leq d$, in (1.33), hence

$$
\begin{equation*}
\int \theta_{\alpha}(x) d x=0, \text { for } \alpha \neq 0 . \tag{A.12}
\end{equation*}
$$

We see that for $\ell<J_{n}, p \in \mathbb{Z}^{d}, \alpha \neq 0$ :

$$
\begin{equation*}
\left|c_{\alpha, \ell, p}^{J_{n}}\right|=2^{-\ell d} \int_{2^{\ell}\left(p+\operatorname{Supp} \theta_{\alpha}\right)}\left(f(x)-f\left(2^{\ell} p\right)\right) \theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right) d x, \tag{A.13}
\end{equation*}
$$

and hence

$$
\begin{equation*}
\left|c_{\alpha, \ell, p}^{J_{n}}\right| \leq c\left(\frac{2^{\ell}}{L_{n}}\right)^{\beta}|f|_{(n)} \leq c 2^{\beta\left(\ell-J_{n}\right)}|f|_{(n)} . \tag{A.14}
\end{equation*}
$$

The right inequality in (A.8) now follows from (A.11), (A.14).
Conversely, expanding $f$ as in (1.35), assume that

$$
\begin{equation*}
\rho_{f} \stackrel{\text { def }}{=} \sup \left\{\left|c_{\alpha, \ell, p}^{J_{n}}\right| 2^{\beta\left(J_{n}-\ell\right)} ; \alpha \in\{0,1\}^{d}, \ell \leq J_{n}, p \in \mathbb{Z}^{d}, \alpha \neq 0 \text { when } \ell<J_{n}\right\}<\infty . \tag{A.15}
\end{equation*}
$$

Observe that for $\bar{\ell}_{1} \leq \bar{\ell}_{0} \leq J_{n}$ and $x \in \mathbb{R}^{d}$,

$$
\begin{align*}
& \left|\sum_{\substack{\alpha, p \\
\bar{\ell}_{1} \leq \ell \leq \bar{\ell}_{0}}} c_{\alpha, \ell, p}^{J_{n}} \theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right)\right| \leq \rho_{f} \sum_{\substack{\alpha, p \\
\bar{\ell}_{1} \leq \ell \leq \bar{\ell}_{0}}} 2^{\beta\left(\ell-J_{n}\right)}\left|\theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right)\right|  \tag{A.16}\\
& \leq c \rho_{f} \sum_{\bar{\ell}_{1} \leq \ell \leq \bar{\ell}_{0}} 2^{\beta\left(\ell-J_{n}\right)} \leq c \rho_{f} 2^{\beta\left(\bar{\ell}_{0}-J_{n}\right)},
\end{align*}
$$

since for each $\ell \leq J_{n}$, at most $c$ of the summands in the expression after the first inequality do not vanish. In particular $\sum_{\bar{\tau}_{1} \leq \ell \leq J_{n}}^{\alpha, p} c_{\alpha, \ell, p}^{J_{n}} \theta_{\alpha, \ell, p}$ converges uniformly (and of course in $L^{2}$ ) towards $f$, which is continuous and satisfies:

$$
\begin{equation*}
|f|_{\infty} \leq c \rho_{f} \tag{A.17}
\end{equation*}
$$

Note that when $|x-y| \geq 2^{J_{n}}$, one has

$$
\begin{equation*}
|f(x)-f(y)| \leq 2|f|_{\infty} \leq 2 c \rho_{f} \leq c \rho_{f}\left|\frac{x-y}{L_{n}}\right|^{\beta} \tag{A.18}
\end{equation*}
$$

On the other hand, when $|x-y|<2^{J_{n}}$, so that

$$
\begin{equation*}
2^{\bar{\ell}_{0}}<|x-y| \leq 2^{\bar{\ell}_{0}+1}, \text { with } \bar{\ell}_{0}<J_{n}, \tag{A.19}
\end{equation*}
$$

we introduce $\widetilde{f}=\sum_{\substack{\alpha, p \\ \bar{\ell}_{0} \leq \ell \leq J_{n}}} c_{\alpha, \ell, p}^{J_{n}} \theta_{\alpha, \ell, p}$, and find

$$
\begin{align*}
& |f(x)-f(y)| \leq 2|f-\widetilde{f}|_{\infty}+|\widetilde{f}(x)-\widetilde{f}(y)| \stackrel{(A .16)}{\leq} \\
& c \rho_{f} 2^{\beta\left(\bar{\ell}_{0}-J_{n}\right)}+\left|\sum_{\bar{\ell}_{0, p} \leq \ell \leq J_{n}} c_{\alpha, \ell, p}^{J_{n}}\left(\theta_{\alpha}\left(\frac{x}{2^{\ell}}-p\right)-\theta_{\alpha}\left(\frac{y}{2^{\ell}}-p\right)\right)\right| \leq \\
& c \rho_{f} 2^{\beta\left(\bar{\ell}_{0}-J_{n}\right)}+c \rho_{f} \sum_{\bar{\ell}_{0} \leq \ell \leq J_{n}} 2^{\beta\left(\ell-J_{n}\right)}\left|\frac{x-y}{2^{\ell}}\right| \stackrel{(A .19)}{\leq}  \tag{A.20}\\
& c \rho_{f}\left(\left|\frac{x-y}{L_{n}}\right|^{\beta}+|x-y| \sum_{\ell_{0} \leq \ell \leq J_{n}} 2^{-(1-\beta) \ell-\beta J_{n}}\right) \leq \\
& c \rho_{f}\left(\left|\frac{x-y}{L_{n}}\right|^{\beta}+|x-y| 2^{-(1-\beta) \bar{\ell}_{0}-\beta J_{n}}\right) \stackrel{(A .19)}{\leq} c \rho_{f}\left|\frac{x-y}{L_{n}}\right|^{\beta} .
\end{align*}
$$

Combining (A.17), (A.18), (A.20), the proof of (A.8) is completed.
We now turn to the proof of (A.9). We begin with the proof of the left-hand inequality. We denote with $\Phi_{A}$ the middle expression of (A.9), which we assume finite. We pick a $[0,1]$-valued function $h$, compactly supported such that

$$
\begin{align*}
& |h|_{(n)} \leq 3, \text { and }  \tag{A.21}\\
& A(h g)=A(g) \text { for any bounded measurable } g . \tag{A.22}
\end{align*}
$$

Indeed given our assumptions on $A$, we can for instance pick $h$ of the form (1.37), with $u$ large, and use (A.6). For $g$ with $|g|_{(n)} \leq 1$, we define

$$
\begin{equation*}
f=h g \tag{A.23}
\end{equation*}
$$

so that expanding $f$ as in (1.35) with ( $J_{n}$ in place of $j_{0}$ ), and keeping the notation (A.15) for $\rho_{f}$, we find:

$$
\begin{equation*}
\rho_{f} \stackrel{(A .8)}{\leq} c|f|_{(n)} \stackrel{(1.29),(A .21)}{\leq} c|g|_{(n)} \leq c . \tag{A.24}
\end{equation*}
$$

Since $A(g)=A(f)$ is bounded measurable and compactly supported, we find:

$$
\begin{equation*}
A(g)=A(f) \stackrel{(1.35),(A .10)}{=} \sum_{\substack{\alpha, \ell \leq J_{n, p} \\ \alpha \neq 0, \text { for } \ell<J_{n}}} \frac{1}{2^{\ell d}}\left\langle\theta_{\alpha, \ell, p}, A(f)\right\rangle \theta_{\alpha, \ell, p} . \tag{A.25}
\end{equation*}
$$

We also know that the partial sums $\widetilde{f}$, cf. above (A.20), converge uniformly to $f$, as $\ell_{0}$ tends to $-\infty$, and only finitely many terms in the sum defining $f$ do not identically vanish on the support of $h$. Using the continuity of $A$ for the sup-norm, we find that for $\alpha \in\{0,1\}^{d}, \ell \leq J_{n}, p \in \mathbb{Z}^{d}$, with $\alpha \neq 0$, for $\ell<J_{n}$, with hopefully obvious notations:

$$
\begin{align*}
& 2^{\beta\left(J_{n}-\ell\right)} \frac{1}{2^{\ell d}}\left|\left\langle\theta_{\alpha, \ell, p}, A(f)\right\rangle\right| \leq 2^{\beta\left(J_{n}-\ell\right)-\ell d} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\left|c_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}^{J_{n}}(f)\right|\left|\left\langle\theta_{\alpha, \ell, p}, A\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)\right\rangle\right| \\
& \stackrel{(A .15)}{\leq} \rho_{f} \sum_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} \frac{2^{\beta \ell^{\prime}}}{2^{\beta \ell}} \frac{1}{2^{\ell d}}\left|\left\langle\theta_{\alpha, \ell, p}, A\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)\right\rangle\right| . \tag{A.26}
\end{align*}
$$

Keeping in mind (A.24), we see coming back to (A.25) with the help of (A.8) that $A(g)$ is a $\beta$-Hölder continuous function and:

$$
\begin{equation*}
|A(g)|_{(n)} \leq c \Phi_{A},(\text { cf. above (A.21) for the notation) } \tag{A.27}
\end{equation*}
$$

This proves the left inequality of (A.9).
We now prove the right inequality of (A.9). Without loss of generality we assume $\|A\|_{n}$ finite, i.e. $A$ maps boundedly the set of bounded $\beta$-Hölder continuous functions endowed with $|\cdot|_{(n)}$, into itself. Consider $\alpha_{0} \in\{0,1\}^{d}, \ell_{0} \leq J_{n}, p_{0} \in \mathbb{Z}^{d}$, with $\alpha_{0} \neq 0$, if $\ell_{0}<J_{n}$, and $\mathcal{J}^{\prime}$ a finite set of $\left(\alpha^{\prime}, \ell^{\prime}, p^{\prime}\right)$ satisfying analogous constraints. Using the convention sign $(0)=1$, we define

$$
\begin{equation*}
f=\sum_{\mathcal{J}^{\prime}} \operatorname{sign}\left(\left\langle\theta_{\alpha_{0}, \ell_{0}, p_{0}}, A\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)\right\rangle\right) 2^{\beta \ell^{\prime}} \theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}} . \tag{A.28}
\end{equation*}
$$

From (A.8), we deduce that

$$
\begin{gather*}
|f|_{(n)} \leq c 2^{\beta J_{n}} \text {, and that }  \tag{A.29}\\
|A(f)|_{(n)} \stackrel{(A .8),(A .10)}{\geq} c 2^{\beta\left(J_{n}-\ell_{0}\right)} \frac{1}{2^{\ell_{0} d}}\left|\left\langle\theta_{\alpha_{0}, \ell_{0}, p_{0}}, A(f)\right\rangle\right| \\
\stackrel{(A .28)}{=} c 2^{\beta\left(J_{n}-\ell_{0}\right)} \sum_{\mathcal{J}^{\prime}} \frac{2^{\beta \ell^{\prime}}}{2^{\ell_{0} d}}\left|\left\langle\theta_{\alpha_{0}, \ell_{0}, p_{0}}, A\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)\right\rangle\right| \\
\stackrel{(A .29)}{\geq} c|f|_{(n)} \sum_{\mathcal{J}^{\prime}} \frac{2^{\beta^{\prime}}}{2^{(d+\beta) \ell_{0}}}\left|\left\langle\theta_{\alpha_{0}, \ell_{0}, p_{0}}, A\left(\theta_{\alpha^{\prime}, \ell^{\prime}, p^{\prime}}\right)\right\rangle\right| .
\end{gather*}
$$

Since $f$ in (A.28) is not identically zero and $\alpha_{0}, \ell_{0}, p_{0}$, and $\mathcal{J}^{\prime}$ are arbitrary, we find that

$$
\begin{equation*}
\|A\|_{n} \geq c \Phi_{A}, \text { (cf. above (A.21) for the notation). } \tag{A.31}
\end{equation*}
$$

This finishes the proof of (A.9), and of Proposition A.2.

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